



Assessment of Inertial Confinement Fusion Targets

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ASSESSMENT OF INERTIAL CONFINEMENT FUSION TARGETS

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Panel on the Assessment of Inertial Confinement Fusion Targets
Board on Physics and Astronomy
Board on Energy and Environmental Systems
Division on Engineering and Physical Sciences
NATIONAL RESEARCH COUNCIL
OF THE NATIONAL ACADEMIES

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35 responsible for the report were chosen for their special competences and with regard for appropriate
36 balance.

37

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Preface and Acknowledgments

In the fall of 2010, the Office of the U.S. Department of Energy’s (DOE’s) Under Secretary for Science asked for a National Research Council (NRC) committee to investigate the prospects for generating power using inertial confinement fusion (ICF) concepts, acknowledging that a key test of viability for this concept—ignition¹—could be demonstrated at the National Ignition Facility (NIF) at Lawrence Livermore National Laboratory (LLNL) in the relatively near term. The committee was asked to provide an unclassified report. However, DOE indicated that to fully assess this topic, the committee’s deliberations would have to be informed by the results of some classified experiments and information, particularly in the area of ICF targets and nonproliferation. Thus, the Panel on the Assessment of Inertial Confinement Fusion Targets (“the panel”) was assembled, composed of experts able to access the needed information (for member biographies, see Appendix A). The panel was charged with advising the Committee on the Prospects for Inertial Confinement Fusion Energy Systems on these issues, both by internal discussion and by this unclassified report. The statement of task for the panel is given in Box P.1.

Box P.1 Statement of Task for the Panel on the Assessment of Inertial Confinement Fusion Targets

A Panel on Fusion Target Physics (“the panel”) will serve as a technical resource to the Committee on Inertial Confinement Energy Systems (“the Committee”) and will prepare a report that describes the R&D challenges to providing suitable targets, on the basis of parameters established and provided to the Panel by the Committee.

The Panel on Fusion Target Physics will prepare a report that will assess the current performance of fusion targets associated with various ICF concepts in order to understand:

1. The spectrum output;
2. The illumination geometry;
3. The high-gain geometry; and
4. The robustness of the target design.

The panel will also address the potential impacts of the use and development of current concepts for Inertial Fusion Energy on the proliferation of nuclear weapons information and technology, as appropriate. The Panel will examine technology options, but will not provide recommendations specific to any currently operating or proposed ICF facility.

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The panel interpreted the terms used in its statement of task in the following way. “Illumination geometry” not only is interpreted to mean the physical arrangement and timing of laser or particle beams incident on the target but also is generalized to mean “delivering driver energy to the target.” In this way, the magnetic forces in pulsed-power schemes are also included. “High-gain geometry” is interpreted as designs that enable the energy incident on the target to be converted efficiently into fuel burn and high yield.² “Spectrum output” is interpreted to include all of the types of emissions (photons, ions, neutrons, and debris) from the fusion target and their energy spectra. Depending on the type of reaction chamber used (solid wall,

¹ The operative definition of ignition adopted by the panel, “gain greater than unity,” is the same as that used in the earlier National Research Council NRC report: *Review of the Department of Energy's Inertial Confinement Fusion Program*, Washington, D.C.: National Academy Press (1997).

² High yield is defined broadly as much more than 10 times the fusion energy produced as driver energy delivered to the target.

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203 wetted wall, liquid wall, gas-filled, evacuated, and so on) these emissions may or may not reach
204 the chamber wall; however, a detailed discussion of the effects on the wall is beyond the scope of
205 this report. “Robustness of the target design” is interpreted in two ways: (1) the inherent
206 “physics robustness,” which relates to the performance margins of the design being large enough
207 compared to the physics uncertainties that reliable performance can be assured under ideal
208 conditions, and (2) “engineering robustness,” which relates to the target’s ability to deliver
209 reliable performance even under nonideal conditions such as variations in driver energy, target
210 manufacturing defects, errors in target positioning, or driver beam misalignment.

211 This unclassified report contains all of the panel’s conclusions and recommendations. In
212 some cases, additional support and documentation required the discussion of classified material,
213 which appears in classified appendixes in a separate version of this report. ICF is an active
214 research field, and scientific understanding continues to evolve. The information discussed here
215 is accurate as of the date presented to the panel (see Appendix B), though in some cases more
216 recent updates are included; if so, this is noted in the text.

217 This report was reviewed in draft form by individuals chosen for their diverse
218 perspectives and technical expertise in accordance with procedures approved by the National
219 Research Council’s Report Review Committee. The purpose of this independent review is to
220 provide candid and critical comments that will assist the institution in making its published
221 report as sound as possible and to ensure that the report meets institutional standards for
222 objectivity, evidence, and responsiveness to the study charge. The review comments and draft
223 manuscript remain confidential to protect the integrity of the process.

224 We wish to thank the following individuals for their review of this report:
225

226 Bedros Afeyan, Polymath Research Inc.,
227 Roger Bangerter, E.O. Lawrence Berkeley National Laboratory (retired),
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233 Arjun Makhijani, Institute for Energy and Environmental Research,
234 David Overskei, Decision Factors Inc.,
235 Robert Rosner, University of Chicago, and
236 Douglas Wilson, Los Alamos National Laboratory.

237
238 Although the reviewers listed above have provided many constructive comments and
239 suggestions, they were not asked to endorse the conclusions or recommendations, nor did they
240 see the final draft of the report before its release. The review of this report was overseen by
241 Louis J. Lanzerotti, New Jersey Institute of Technology. Appointed by the NRC, he was
242 responsible for making certain that an independent examination of this report was carried out in
243 accordance with institutional procedures and that all review comments were carefully
244 considered. Responsibility for the final content of this report rests entirely with the authoring
245 committee and the institution.

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247 The panel also thanks the NRC staff for its dedicated work, in particular Sarah Case, who
248 got the panel started off on the correct path, and Greg Eyring, who persevered in getting both the
249 classified and the unclassified reports over many hurdles.

250

251 John F. Ahearne, Chair

252 Panel on Assessment of Inertial Confinement Fusion Targets

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Summary

In the fall of 2010, the Office of the U.S. Department of Energy’s (DOE’s) Under Secretary for Science asked for a National Research Council (NRC) committee to investigate the prospects for generating power using inertial fusion energy (IFE), noting that a key test of viability for this concept—ignition³—could be demonstrated at the National Ignition Facility (NIF) at Lawrence Livermore National Laboratory (LLNL) in the relatively near term. In response, the NRC formed both the Committee on the Assessment of the Prospects for Inertial Fusion Energy (“the committee”) to investigate the overall prospects for IFE in an unclassified report and the separate Panel on Fusion Target Physics (“the panel”) to focus on issues specific to fusion targets, including the results of relevant classified experiments and classified information on the implications of IFE targets for the proliferation of nuclear weapons.

This is the report of the Panel on Fusion Target Physics, which is intended to feed into the broader assessment of IFE being done by the NRC committee. It consists of an unclassified body, which contains all of the panel’s conclusions and recommendations, as well as three classified appendices, which provide additional support and documentation.

BACKGROUND

Fusion is the process by which energy is produced in the sun, and, on a more human scale, is the one of the key processes involved in the detonation of a thermonuclear bomb. If this process could be “tamed” to provide a controllable source of energy that can be converted to electricity—as nuclear fission has been in currently operating nuclear reactors—it is possible that nuclear fusion could provide a new method for producing low-carbon electricity to meet the U. S. and world growing energy needs.

For inertial fusion to occur in a laboratory, fuel material (typically deuterium and tritium) must be confined for an adequate length of time at an appropriate density and temperature to overcome the Coulomb repulsion of the nuclei and allow them to fuse. In inertial confinement fusion (ICF)—the concept investigated in this report⁴—a driver (e.g., a laser, particle beam, or pulsed magnetic field) delivers energy to the fuel target, heating and compressing it to the conditions required for ignition. Most ICF concepts compress a small amount of fuel directly to thermonuclear burn conditions (a hot spot) and propagate the burn via alpha particle deposition through adjacent high-density fuel regions, thereby generating a significant energy output.

There are two major concepts for inertial confinement fusion target design: direct-drive targets, in which the driver energy strikes directly on the fuel capsule, and indirect-drive targets, in which the driver energy first strikes the inside surface of a hollow chamber (a hohlraum) surrounding the fuel capsule, producing energetic X-rays that compress the fuel capsule. Conventional direct and indirect drive share many key physics issues (e.g., energy coupling, the need for driver uniformity, and hydrodynamic instabilities); however, there are also issues that are unique to each concept.

³ The operative definition of ignition adopted by the panel, “gain greater than unity,” is the same as that used in the earlier National Research Council NRC report: *Review of the Department of Energy's Inertial Confinement Fusion Program*, Washington, D.C.: National Academy Press (1997).

⁴ Inertial confinement fusion (ICF) is the process by which the target is heated and compressed by the driver to reach fusion conditions. Inertial fusion energy (IFE) is the process by which useful energy is extracted from ignition and burn of ICF fuel targets.

333 The only facility in the world that was designed to conduct ICF experiments that address
 334 the ignition scale is the NIF at LLNL. The NIF driver is a solid-state laser. For the first ignition
 335 experiments, the NIF team has chosen indirect-drive targets. The NIF can also be configured for
 336 direct drive. In addition, important work on laser-driven, direct-drive targets (albeit at less than
 337 ignition scale) is also under way in the United States at the Naval Research Laboratory and the
 338 OMEGA laser at the University of Rochester. Heavy-ion-beam drivers are being investigated at
 339 the Lawrence Berkeley National Laboratory (LBNL), LLNL, and the Princeton Plasma Physics
 340 Laboratory (PPPL), and magnetic implosion techniques are being explored on the Z machine at
 341 Sandia National Laboratory (SNL) and at Los Alamos National Laboratory (LANL). Important
 342 ICF research is also under way in other countries, as discussed later in this report.

343 **SPECIFIC CONCLUSIONS AND RECOMMENDATIONS**

344 The panel's key conclusions and recommendations, all of them specific to various aspects
 345 of inertial confinement fusion, are presented below. They are labeled according to the chapter
 346 and number order in which they appear in the text, to provide the reader with an indicator of
 347 where to find a more complete discussion. This summary ends with two overarching conclusions
 348 and an overarching recommendation derived from viewing all of the information presented to the
 349 panel as a whole.

350 **Targets for Indirect Laser Drive**

351 **CONCLUSION 4-1: The national program to achieve ignition using indirect laser drive**
 352 **has several physics issues that must be resolved if it is to achieve ignition.** At the time of this
 353 writing, the capsule/hohlraum performance in the experimental program, which is carried out at
 354 the NIF, has not achieved the compressions and neutron yields expected based on computer
 355 simulations. At present, these disparities are not well understood. While a number of hypotheses
 356 concerning the origins of the disparities have been put forth, it is apparent to the panel that the
 357 treatments of the detrimental effects of laser-plasma interactions (LPI) in the target performance
 358 predictions are poorly validated and may be very inadequate. A much better understanding of
 359 laser-plasma interactions will be required of the ICF community.

360 **CONCLUSION 4-2: Based on its analysis of the gaps in current understanding of target**
 361 **physics and the remaining disparities between simulations and experimental results, the**
 362 **panel assesses that ignition using laser indirect drive is not likely in the next several years.**
 363 The National Ignition Campaign (NIC) plan—as the panel understands it—suggests that ignition
 364 is planned after the completion of a tuning program lasting 1-2 years that is presently under way
 365 and scheduled to conclude at the end of FY2012. While this success-oriented schedule remains
 366 possible, resolving the present issues and addressing any new challenges that might arise are
 367 likely to push the timetable for ignition to 2013-2014 or beyond.

370 **Targets for Indirect-Drive Laser Inertial Fusion Energy**

371 **CONCLUSION 4-4: The target design for a proposed indirect-drive inertial fusion energy**
 372 **system (the laser inertial fusion energy or LIFE program developed by LLNL)**

379 **incorporates plausible solutions to many technical problems, but the panel assesses that the**
 380 **robustness of the physics design for the LIFE target concept is low.**

- 381 • The proposed LIFE target presented to the panel has several modifications relative to
 382 the target currently used in the NIC (for example, rugby hohlraums, shine shields, and
 383 high-density carbon ablaters) and the effects of these modifications may not be
 384 trivial. For this reason, R&D and validation steps would still be needed.
- 385 • There is no evidence to indicate that the margin in the calculated target gain ensures
 386 either its ignition or sufficient gain for the LIFE target. If ignition is assumed, the
 387 gain margin briefed to the panel, which ranged from 25 percent to almost 60 percent
 388 when based on a calculation that used hohlraum and fuel materials characteristic of
 389 the NIC rather than the LIFE target, is unlikely to compensate for the phenomena
 390 relegated to it—for example, the effects of mix—under any but the most extremely
 391 favorable eventuality. In addition, the tight coupling of LIFE to what can be tested on
 392 the NIF constrains the potential design space for laser-driven, indirect-drive IFE.

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Targets for Direct-Drive Laser Inertial Fusion Energy

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397 **CONCLUSION 4-6: The prospects for ignition using laser direct drive have improved**
 398 **enough that it is now a plausible alternative to laser indirect drive for achieving ignition**
 399 **and for generating energy.**

400

- 401 • The major concern with laser direct drive has been the difficulty of achieving the
 402 symmetry required to drive such targets. Advances in beam-smoothing and pulse-
 403 shaping appear to have lessened the risks of asymmetries. This assessment is
 404 supported by data from capsule implosions (performed at the University of
 405 Rochester's OMEGA laser), but it is limited by the relatively low drive energy of the
 406 implosion experiments that have thus far been possible. Because of this, the panel's
 407 assessment of laser-driven, direct-drive targets is not qualitatively equivalent to that
 408 of laser-driven, indirect-drive targets.
- 409 • Further evaluation of the potential of laser direct-drive targets for IFE will require
 410 experiments at drive energies much closer to the ignition scale.
- 411 • Capsule implosions on OMEGA have established an initial scaling point that
 412 indicates the potential of direct-drive laser targets for ignition and high yield.
- 413 • Polar direct-drive targets⁵ will require testing on the NIF.
- 414 • Demonstration of polar-drive ignition on the NIF will be an important step toward an
 415 IFE program.
- 416 • If a program existed to reconfigure NIF for polar drive, direct-drive experiments that
 417 address the ignition scale could be performed as early as 2017.

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⁵ In polar direct drive, the driver beams are clustered in one or two rings at opposing poles. To increase the uniformity of the drive, polar drive beams strike the capsule obliquely, and the driver energy is biased in favor of the more equatorial beams.

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Fast Ignition

Fast ignition (FI) requires a combination of long-pulse (implosion) and short-pulse (ignition) lasers. Aspects of fast ignition by both electrons and protons were briefed to the panel. Continued fundamental research into fast ignition theory and experiments, the acceleration of electrons and ions by ultrashort-pulse lasers, and related high-intensity laser science is justified. However, issues surrounding low laser-target energy coupling, a complicated target design, and the existence of more promising concepts (such as shock ignition) led the panel to the next conclusion regarding the relative priority of fast ignition for fusion energy.

CONCLUSION 4-5: At this time, fast ignition appears to be a less promising approach for IFE than other ignition concepts.

Laser-Plasma Interactions

A variety of LPI take place when an intense laser pulse hits the target capsule or surrounding hohlraum. Undesirable effects include backscattering of laser light, which can result in loss of energy; cross-beam energy transfer among intersecting laser beams, which can cause loss of energy or affect implosion symmetry; acceleration of suprathermal “hot electrons,” which then can penetrate and preheat the capsule’s interior and limit later implosion; and filamentation, a self-focusing instability that can exacerbate other LPI. LPI have been a key limiting factor in laser inertial confinement fusion, including the NIC indirect-drive targets, and are still incompletely understood.

CONCLUSION 4-11: The lack of understanding surrounding laser-plasma interactions remains a substantial but as yet unquantified consideration in ICF and IFE target design.

RECOMMENDATION 4-1: DOE should foster collaboration among different research groups on the modeling and simulation of laser-plasma interactions.

Heavy-Ion Targets

A wide variety of heavy-ion target designs has been investigated, including indirect-drive, hohlraum/capsule targets that resemble NIC targets. Recently, the emphasis has shifted to direct-drive targets, but to date the analysis of how these targets perform has been based on computation rather than experiment, and the codes have not been benchmarked with experiments in relevant regimes.

CONCLUSION 4-12: The U.S. heavy-ion-driven fusion program is considering direct-drive and indirect-drive target concepts. There is also significant current work on advanced target designs.⁶ This work is at a very early stage, but if successful, may provide very high gain.

⁶ Advanced designs include direct-drive, conical X-target configurations, see Chapter 2.

- 465 • The work in the heavy-ion fusion (HIF) program involves solid and promising
466 science.
467 • Work on heavy-ion drivers is complementary to the laser approaches to IFE and
468 offers a long-term driver option for beam-driven targets.
469 • The HIF program relating to advanced target designs is in a very early stage and is
470 unlikely to be ready for technical assessment in the near term.
471 • The development of driver technology will take several years and the cost to build a
472 significant accelerator driver facility for any target is likely to be very high.

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Z-Pinch Targets

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477 Current Z-pinch direct-drive concepts utilize the pressure of a pulsed, high magnetic field
478 to implode deuterium-tritium fuel to fusion conditions. Simulations predict that directly using the
479 pressure of the magnetic field to implode and compress the target can greatly increase the
480 efficiency with which the electrical energy is coupled to the fuel as compared with the efficiency
481 of indirect drive from Z-pinch X-ray sources. There is work under way on both classified and
482 unclassified target designs.

483

484 **CONCLUSION 4-13: Sandia National Laboratory is leading a research effort on a Z-pinch**
485 **scheme that has the potential to produce high gain with good energy efficiency, but**
486 **concepts for an energy delivery system based on this driver are too immature to be**
487 **evaluated at this time.**

488 It is not yet clear that the work at SNL will ultimately result in the high gain predicted by
489 computer simulations, but initial results are promising and it is the panel's opinion that
490 significant progress in the physics may be made in a year's time. The pulsed power approach is
491 unique in that its goal is to deliver a large amount of energy (~10 MJ) to targets with good
492 efficiency (≥ 10 percent) and to generate large fusion yields at low repetition rates.

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Target Fabrication

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497 Current targets for inertial confinement fusion experiments tend to be one-off designs,
498 with specifications that change according to the experiments being run. In contrast, targets for
499 future IFE power plants will have to have standard, low-cost designs that are mass-produced in
500 numbers as high as a million targets per day per power plant. The panel examined the technical
501 feasibility of producing targets for various drivers, including limited aspects of fabrication for
502 IFE. However, a full examination of the issues of mass production and low cost is the province
503 of the NRC IFE committee study.

504

505 **CONCLUSION 4-7: In general, the science and engineering of manufacturing fusion**
506 **targets for laser-based ICF are well advanced and meet the needs of those experiments,**
507 **although additional technologies may be needed for IFE.** Extrapolating this status to predict
508 the success of manufacturing IFE targets is reasonable if the target is only slightly larger than the
509 ICF target and the process is scalable. However, subtle additions to the design of the ICF target

510 to improve its performance (greater yield) and survivability in an IFE power plant may
511 significantly affect the manufacturing paradigm.

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Proliferation Risks of IFE

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516 Many modern nuclear weapons rely on a fusion stage as well as a fission stage, and there
517 has been discussion of the potential for host state proliferation—particularly vertical
518 proliferation⁷—associated with the siting of an IFE power plant. The panel was asked to evaluate
519 the proliferation risks associated with IFE, particularly with regard to IFE targets.

520

521 **CONCLUSION 3-1: At present, there are more proliferation concerns associated with**
522 **indirect-drive targets than with direct-drive targets.** However, the spread of technology
523 around the world may eventually render these concerns moot. Remaining concerns are likely to
524 focus on the use of classified codes for target design.

525

526 **CONCLUSION 3-2: The nuclear weapons proliferation risks associated with fusion power**
527 **plants are real but are likely to be controllable. These risks fall into three categories:**

528

- Knowledge transfer,
- Special Nuclear Material (SNM) production, and
- Tritium diversion.

530

531

532

OVERARCHING CONCLUSIONS AND RECOMMENDATION

533 While the focus of this panel was on ICF target physics, the need to evaluate driver-target
534 interactions required considering driver characteristics as well. This broader analysis led the
535 panel to the following overarching conclusions and a recommendation.

536

537 **OVERARCHING CONCLUSION 1: NIF has the potential to support the development and**
538 **further validation of physics and engineering models relevant to several IFE concepts, from**
539 **indirect-drive hohlraum designs to polar direct-drive ICF and shock ignition.**

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- **In the near to intermediate term, NIF is the only platform that can provide information relevant to a wide range of IFE concepts at ignition scale. Insofar as target physics is concerned, it is a modest step from NIF scale to IFE scale.**
- **Targets for all laser-driven IFE concepts (both direct-drive and indirect-drive) can be tested on NIF. In particular, reliable target performance would need to be demonstrated before investments could confidently be made in development of laser-driven IFE target designs.**

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OVERARCHING CONCLUSION 2: It would be advantageous to continue research on a range of IFE concepts, for two reasons:

⁷ Vertical proliferation refers to the enhancement of a country's capability to move from simple weapons to more sophisticated weapons.

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- 559
- **The challenges involved in the current laser indirect-drive approach in the single-pulse National Nuclear Security Administration program at the NIF have not yet been resolved and**
 - **The alternatives to laser indirect drive have technical promise to produce high gain.**

560 In particular, the panel concludes that laser direct drive is a viable concept to be pursued
561 on the NIF. SNL's work on Z-pinch can serve to mitigate risk should the NIF not operate as
562 expected. This work is at a very early stage but is highly complementary to the NIF approach,
563 because none of the work being done at SNL relies on successful ignition at the NIF, and key
564 aspects of the target physics can be investigated on the existing Z-machine. Finally, emerging
565 heavy-ion designs could be fruitful in the long term.

566

567 **OVERARCHING RECOMMENDATION: The panel recommends against pursuing a**
568 **down-select decision for IFE at this time, either for a specific concept such as LIFE or for a**
569 **specific target type/driver combination.**

570

571 Further R&D will be needed on indirect drive and other ICF concepts, even following
572 successful ignition at the NIF, to determine the best path for IFE in the coming decades.

573

1

Introduction

Inertial fusion energy (IFE) has been a concept since the 1970s, and the National Research Council (NRC) has performed several reviews of the Department of Energy's (DOE's) programs for inertial confinement fusion (ICF)—the essential concept underlying IFE—since that time (NRC 1986, 1990, and 1997). This report of the Panel on Fusion Target Physics supports and informs a broader study on the prospects for IFE being undertaken by a separate NRC committee.⁸ The broader study is motivated by a desire on the part of DOE, the sponsor, to determine a clearer path forward for the IFE concept, in view of the prospect that a key test of viability for this concept—ignition—can be demonstrated at the National Ignition Facility (NIF) at Lawrence Livermore National Laboratory (LLNL) in the relatively near term.

To address its Statement of Task (see the Preface), the panel heard from many sources, listed in Appendix B, and visited several laboratories involved in U.S. efforts in ICF and IFE—LLNL, Sandia National Laboratory, Lawrence Berkeley National Laboratory, the University of Rochester Laboratory for Laser Energetics, and the Naval Research Laboratory—and heard from representatives of additional programs at the Los Alamos National Laboratory.

The panel's focus in this study is IFE targets, including both direct-drive and indirect-drive targets. To distinguish its role as clearly as possible from that of the main study committee, the panel drew a conceptual sphere around the outside of the target and considered anything crossing the surface of the sphere (energy coming in, reaction products going out) as well as physics processes taking place inside the sphere, to be within its purview. In addition, the panel considered the technical feasibility of fabricating various target concepts to be within its charge, but deemed the mass manufacturing of high-performance, cost-effective targets for future power plants to be part of the main committee's responsibility. Inevitably, there were certain topics at the interface between the charges of the panel and the main committee, such as the survivability of the injected target in the extreme environment of the reaction chamber. In such cases, the panel felt that it was preferable that the panel and committee reports should overlap rather than risk the possibility that important topics might be left out.

Chapter 2 provides a brief technical background on IFE and a discussion of key concepts related to ICF targets and their role in IFE. In Chapter 3, the proliferation risks of specific target designs are discussed, as well as the broader proliferation risks associated with IFE plants and research facilities. Chapter 4 evaluates the current status of various targets, considering the results of actual experiments on their performance as well as the analytical and predictive capabilities of available codes and simulations. This analysis is used to characterize the state of our current understanding of fusion target physics and to identify the major issues that remain to be resolved. The classified version of this report contains additional appendixes discussing classified material that the panel considers relevant to its conclusions and recommendations.

⁸ The Committee on the Prospects for Inertial Confinement Fusion Energy Systems.

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Technical Background

This chapter briefly introduces the key concepts necessary to understand inertial confinement fusion (ICF), inertial fusion energy (IFE), and target physics.

INERTIAL CONFINEMENT FUSION AND INERTIAL FUSION ENERGY

Nuclear fusion—the process by which the nuclei of atoms such as deuterium or tritium combine to form a heavier nucleus, such as that of helium—can release a significant amount of energy. Fusion is the process by which energy is produced in the sun and, on a more human scale, is the one of the key processes involved in the detonation of a thermonuclear bomb.

If this process can be tamed to provide a controllable source of energy that can be converted to electricity—as the nuclear fission process is used in nuclear reactors—it is possible that nuclear fusion could be a new way to produce low-carbon electricity to meet the growing energy needs of the United States and the world. However, this possibility is far from imminent, and a great deal of scientific and engineering work remains to be done before a commercial nuclear fusion plant can be demonstrated.

For inertial fusion to occur in a laboratory, heating of the fuel material (typically deuterium and tritium) must be confined to a small enough hot spot to overcome the Coulomb repulsion of the nuclei and allow fusion to initiate in a small region of the fuel (“ignition”). If successful, this process will release sufficient energy to sustain the fusion “burn” that will propagate through the fuel, generating a significant energy output. Two concepts are typically discussed for accomplishing this confinement: (1) magnetic confinement fusion (MCF), in which magnetic fields are used to confine the plasma, and (2) ICF, the topic of the current report, in which a driver delivers energy to the surface of a pellet of fuel, heating and compressing it. Potential drivers include lasers, particle beams, and X-rays, among other concepts.

In ICF, energy supplied by the driver is applied, either directly or indirectly, to the outer layer of a fuel pellet that is typically made up of an ablator material (e.g., beryllium, doped plastic, or high-density carbon) that explodes outward as it heats. This outward explosion of the surface layer forces the remainder of the fuel (typically light elements such as deuterium and tritium) to accelerate inward to conserve momentum. The timing of the inward fuel acceleration is controlled carefully in order to compress the fuel using a minimum of energy. At the same time, sudden increases in the driver power profile both accelerate the implosion and send shock waves into the center of the fuel, heating it sufficiently that fusion reactions begin to occur.⁹

The goal of ICF is to initiate a self-sustaining process in which the energetic alpha particles emitted by the ongoing fusion reactions heat the surrounding fuel to the point where it also begins to undergo fusion reactions. The percentage of fuel that undergoes fusion is referred to as the “burn-up fraction.” The fuel gain G (defined as the ratio of the total energy released by the target to the driving beam energy impinging upon it) depends on the burn-up fraction, and gains greater than about 10 will need to be demonstrated to validate the target physics of any approach to a practical IFE power plant.

⁹ What is described here is known as hot-spot ignition; other potential concepts for ignition are being considered, and are introduced briefly later in this chapter.

659 Important target physics includes processes that deflect or absorb driver energy within the
660 target; the transport of energy within the target; capsule preheat; conversion of energy to the
661 inward-directed implosion by ablation; fuel compression and heating; thermonuclear reactions;
662 transport and deposition of neutron and alpha-particle energy resulting in bootstrapping
663 thermonuclear reactions; and hydrodynamic disassembly and output. Models exist for all of these
664 processes, but some are more predictive than others. Some processes are difficult to simulate,
665 such as laser-plasma interactions, the generation and transport of hot electrons in self-consistent
666 magnetic fields, nonlocal-thermal-equilibrium atomic physics, hydrodynamic instabilities, mix,
667 and debris generation. These models continue to evolve to keep pace with experiments. Other
668 processes, such as large-scale hydrodynamics, thermonuclear reactions, and X-ray-, neutron- and
669 alpha-particle transport appear to be simulated adequately using standard numerical models.

670 The Department of Energy (DOE) is funding multiple efforts to investigate the physics of
671 ICF; many of these efforts have the potential to inform current understanding of the prospects for
672 IFE. Over the next several years, experiments will be ongoing at the National Ignition Facility
673 (NIF) at Lawrence Livermore National Laboratory (LLNL) that are aimed at achieving ICF
674 ignition. At the same time, experiments such as those at the University of Rochester's Laboratory
675 for Laser Energetics, the Naval Research Laboratory, Lawrence Berkeley Laboratory, and Sandia
676 National Laboratory continue to advance our understanding and control of ICF using different
677 technology and physics approaches. However, it should be recognized that up to this point, the
678 majority of the funding and efforts related to ICF target physics are provided by—and related
679 to—the U.S. nuclear weapons program and its stockpile stewardship efforts and are not directly
680 aimed at energy applications.

681 The DOE's Centurion-Halite program revolved around a series of underground
682 experiments conducted in the 1980s in which target capsules were driven by the energy from
683 nuclear explosions. Additional discussion of the program is provided in classified Appendix D.

684

685 **BASICS OF ICF TARGET PHYSICS AND DESIGN**

686

687 **Target Design: Direct and Indirect Drive, Z-pinch**

688

689 There are two major concepts for ICF target design: direct-drive targets, in which the
690 driver energy (e.g., in the form of laser beams, particle beams, or magnetic field pressure)
691 directly strikes the fuel capsule (see Figure 2-1); and indirect-drive targets, in which the driver
692 energy first strikes a hollow chamber (a "hohlraum") surrounding the fuel capsule, producing
693 energetic X-rays that compress the fuel capsule (see Figure 2-2). Conventional direct and indirect
694 drive share many key physics issues, such as energy coupling, the need for driver uniformity, and
695 hydrodynamic instabilities; however, there are issues that are unique to each concept.

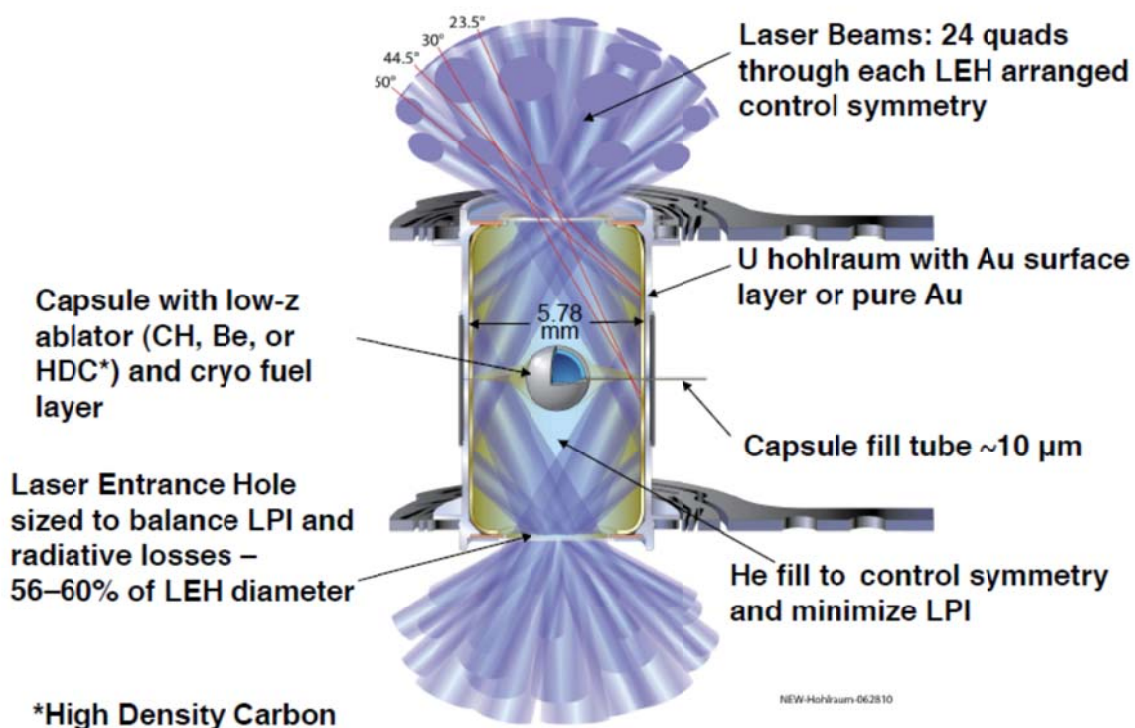
696 Generally, the elements of the fuel capsule are similar for direct drive and indirect drive,
697 at least with respect to laser drivers. Fuel capsules are typically spherical, with several layers: an
698 outer ablator layer; a layer of cryogenic frozen fuel; and a center of gaseous fuel, typically
699 deuterium-tritium (D-T). A sample fuel capsule is shown in Figure 2-3.

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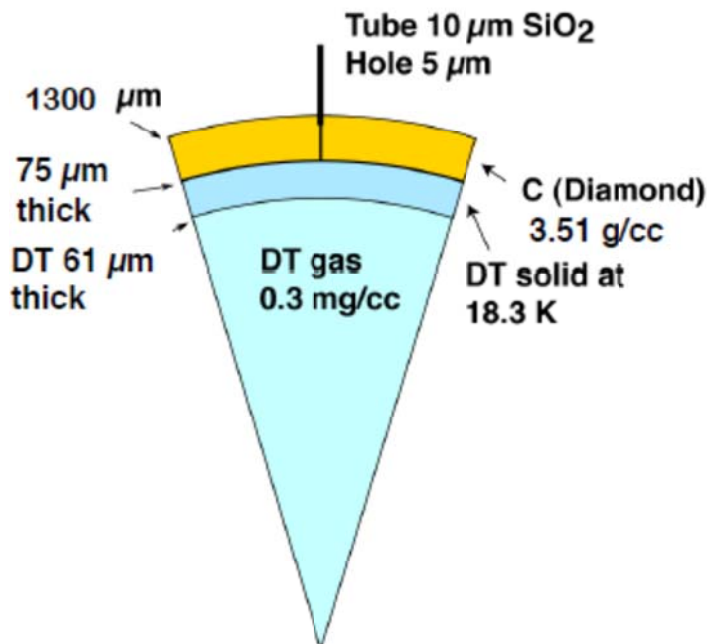
Implosion are driven by the rocket effect from the blow-off plasma.

701
 702 **FIGURE 2-1** In the case of direct drive, the fuel pellet is illuminated symmetrically by the
 703 driver energy, resulting in implosion. SOURCE: R. Betti, University of Rochester, presentation
 704 to the NRC IFE committee titled “Tutorial on the Physics of Inertial Confinement Fusion,” on
 705 April 22, 2011.
 706



707
 708 **FIGURE 2-2** In the case of indirect drive, driver energy incident on a hohlraum is converted to
 709 X-rays, which then impinge symmetrically on the fuel capsule, causing it to implode. This figure
 710

711 shows the laser beam geometry used in the National Ignition Campaign (NIC) at the Lawrence
 712 Livermore National Laboratory. LEH, laser entrance hole; LPI, laser-plasma interactions; HDC,
 713 high-density carbon. SOURCE: J. Lindl, LLNL, presentation to the panel titled “The National
 714 Ignition Campaign on NIF and Its Extension to Targets for IFE,” on February 16, 2011.
 715



716
 717 **FIGURE 2-3** Section of a spherical fuel capsule design showing the ablator layer (in this case
 718 pure carbon), a layer of DT ice, and an inner core of DT gas. Source: J. Lindl, LLNL,
 719 presentation to the panel titled “The National Ignition Campaign on NIF and Its Extension to
 720 Targets for IFE,” on February 16, 2011.

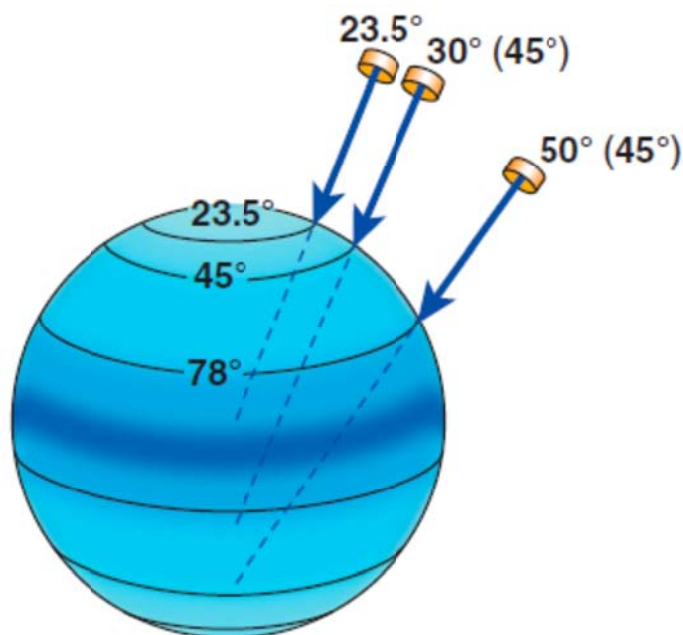
721
 722 Several of the key differences between direct drive and indirect drive for ICF are
 723 discussed briefly in the sections that follow.

724 725 **Direct Drive**

726
 727 Direct-drive concepts for ICF using laser drivers are currently being researched at the
 728 University of Rochester’s Laboratory for Laser Energetics (LLE) and the Naval Research
 729 Laboratory (NRL). Concepts using heavy-ion beam drivers are being studied at Lawrence
 730 Berkeley National Laboratory (LBNL), and Sandia National Laboratories (SNL) is developing
 731 direct-drive concepts for pulsed-power drivers.

732 The major benefit of direct-drive target design is the calculated potential for higher
 733 energy gain than to indirect drive. This relatively large gain is in large part due to avoiding the
 734 losses that occur during the conversion of laser beams or particle beams to X-rays in the
 735 hohlraum, discussed in detail in the next section. Avoiding these losses results in a higher
 736 percentage of driver energy absorbed by the capsule in direct drive, thus increasing the efficiency
 737 and potentially decreasing the size of the driver required.

738 Polar direct drive is a variant of the spherically symmetric, direct-drive illumination
 739 geometry shown in Figure 2-1. As shown in Figure 2-4, the driver beams are clustered in one or
 740 two rings at opposing poles. To increase the uniformity of the drive, polar drive beams strike the
 741 capsule obliquely, and the driver energy is biased in favor of the more equatorial beams.
 742 Although the polar illumination geometry is consequently less efficient than the spherically
 743 symmetric geometry, it is more compatible with the current NIF configuration.
 744



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 746
 747 **FIGURE 2-4** In the polar direct-drive illumination geometry, the driver beams are incident from
 748 directions above and below the fuel capsule but not near the equator. SOURCE: R. L. McCrory,
 749 University of Rochester, presentation to the panel titled “Laser-Driven Inertial Fusion Energy:
 750 Direct-Drive Targets Overview,” on February 16, 2011.

751
 752 Since the 1980s, there has been an ongoing effort in laser science that has been focused
 753 on improving the performance of direct-drive laser systems for both solid-state and KrF lasers.
 754 For solid-state lasers, these advances include frequency tripling (for improved energy coupling
 755 and lower instability growth rates), smoothing by spectral dispersion (SSD), and polarization
 756 smoothing, to reduce imprinting of beam nonuniformities on the target. Recently LLE developed
 757 SSD with multiple phase-modulation frequencies (Multi-FM) and proposed using this technique
 758 to modify NIF for polar direct drive.

759 High-energy KrF lasers were developed to utilize the deep ultraviolet (248 nm)
 760 wavelength of the system. Induced spatial incoherence (ISI) was developed to smooth the beams,
 761 and recently focal zooming¹⁰ was demonstrated to improve the efficiency of coupling the laser
 762 with imploding targets. Direct-drive target experiments on the OMEGA laser have shown steady
 763 improvement towards theoretical yield limits by combining a large number (60) of laser beams,
 764 better laser beam smoothing techniques, and improved beam pointing and target placement at the

¹⁰ Zooming involves reducing the driver spot size to match the diameter of the imploding capsule, thereby increasing the efficiency of energy coupling between driver and target.

765 target chamber center. Although historically much of the discussion of direct-drive fusion has
 766 involved laser drivers (e.g., LLE's work at the OMEGA laser facility and the Nike KrF laser
 767 experiments at NRL), direct-drive ICF has potential for use with other drivers. In particular, the
 768 panel was briefed on direct-drive targets by members of the LBNL heavy-ion driver program.

769 However, there are difficulties involved in using direct-drive fusion. A direct-drive
 770 capsule must tolerate four major sources of perturbations to ignite and burn: drive asymmetry,
 771 inhomogeneous capsule surface finish, ice roughness in the layer between the cryogenic DT and
 772 the DT gas; and driver imprint.¹¹ The effects of the driver imprint and drive asymmetry are
 773 reduced for indirect drive. In addition, without a hohlraum to protect the capsule from the high
 774 temperatures in the chamber, and if there is no buffer gas to protect the chamber walls from
 775 emitted alpha particles, alternative methods must be found to address these threats.

776

777 **Indirect Drive**

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779 As shown in Figure 2-2, indirect drive (whether using laser drivers or an alternative
 780 driver, such as heavy-ion beams) consists of driver beams entering a hohlraum, which is
 781 essentially a hollow cylinder, typically made of gold, or oblong capsule with (in the case of laser
 782 drivers) openings on either end. LLNL is currently leading research into indirect-drive concepts
 783 for laser-driven ICF at the NIF. The driver beams are directed to enter the openings on either end
 784 of the hohlraum, and strike the interior of the hohlraum in four circular arrays, two near the
 785 center, and two nearer the ends (see Figure 2-2). The energy deposited by the laser beams on the
 786 interior of the hohlraum produces a hot plasma that radiates primarily in X-rays at a temperature
 787 of about 300 eV or 3.3 million K. These X-rays are then absorbed by the capsule, resulting in
 788 implosion.

789 A virtue of the hohlraum in an actual IFE target is that it functions as a thermal shroud to
 790 protect the integrity of the cryogenic fuel capsule inside the target. This allows the target
 791 chamber to contain an inert gas (xenon) at low pressure to help protect the walls of the target
 792 chamber from X-rays emitted by high-Z materials in the exploding target.

793

794 *Benefits of Indirect Drive for Smoothing*

795

796 Spatial nonuniformities at any scale can significantly increase the deviation of the actual
 797 implosion of an inertial fusion capsule from the conditions it was designed to achieve, with the
 798 result that the conditions inside the imploded capsule lie in a less favorable location in
 799 thermodynamic phase space than intended. Indirect drive of laser targets was conceived and
 800 developed to eliminate the effects of nonuniformities within each laser beam delivered to the
 801 target chamber.

802 The smoothing obtained through the use of indirect drive is a consequence of
 803 transforming the energy of each laser from a focused beam into thermal radiation. Any
 804 nonuniformity in a laser beam entering an indirect-drive target chamber transfers to the wall of
 805 the hohlraum enclosing the target, heating its material to a heterogeneous plasma. This
 806 heterogeneity is somewhat smoothed by energy transport processes within the radiating plasma
 807 itself, but a stronger smoothing effect occurs because the X-rays originating in each localized

¹¹ For laser drivers, driver imprint occurs early in time when the target ablator is cold and dense. It is related to the asymmetries from modulations in individual laser beams (short wavelength) and perturbations from overlapping drive beams or by beams with slightly differing arrival times and angles of incidence (longer wavelength).

808 mass of plasma affect the entire portion of the target capsule surface to which it has a direct line
 809 of sight. The result is that localized variations in X-ray emission are averaged over the capsule
 810 surface, and rapid changes of drive conditions over the surface of the capsule are eliminated.

811 The development and use of indirect drive was the primary focus of LLNL on the 10-
 812 beam NOVA laser. This experience led to the development of the NIF indirect-drive
 813 configuration, which is much more sophisticated, using 192 laser beams in inner and outer
 814 clusters to control symmetry and pulse shape (see Figure 2-2).

815 Although the capsule absorption of X-rays is more efficient than the direct absorption of
 816 laser light in direct-drive fusion, enough energy is lost in the heating of the hohlraum to
 817 significantly reduce the efficiency of indirect-drive fusion relative to direct-drive fusion. This
 818 results in lower calculated potential gains for indirect-drive fusion targets.

819 As with direct drive, although its primary development historically has been with laser
 820 drivers, indirect drive has been used in IFE system designs with other drivers (e.g., heavy ions
 821 and early Z-pinch schemes). The key is to deposit enough energy on the inner surface of the
 822 hohlraum to produce a hot plasma that radiates thermal X-rays.

823 One of the key reasons that indirect-drive targets were developed is that ICF can model
 824 on a laboratory scale some aspects of a thermonuclear explosion. This is highly useful for the
 825 applications of ICF at the NIF at LLNL that are related to the long-term stewardship of the U.S.
 826 nuclear stockpile. This motivation has been a key aspect in the development of the indirect-drive
 827 approach for IFE, since one could leverage insights from better-funded weapons programs for
 828 the less well funded energy programs. However, there remains debate about whether this
 829 provides significant benefits for energy generation using ICF, and some argue that the indirect-
 830 drive approach—if commercialized and distributed overseas—could increase the risk that
 831 nuclear weapons knowledge and information will proliferate. This topic is analyzed in more
 832 detail in the classified Appendix E and in Chapter 3.

833

834 **Z-pinch Target**

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836 In recent ICF and IFE studies, Z-pinch targets are imploded by the pressure of ultrahigh
 837 magnetic fields generated by high currents (e.g., 20-60 MA for ~100 ns) provided by pulsed-
 838 power generators rather than by the ablation pressure generated by illuminating a capsule with a
 839 high-power laser. While laser fusion capsules are typically spherical shells, Z-pinch targets are
 840 typically conducting cylindrical shells containing DT fuel. Since magnetic field strength
 841 increases inversely with the radius of the conductor in which the current flows (I/r), as long as
 842 the driver has the appropriate electrical characteristics to deliver current to the increasingly high-
 843 inductance target, the magnetic pressure (proportional to B^2) continues to grow, accelerating the
 844 cylindrical implosion and compressing the fuel. For appropriate design conditions, the DT fuel
 845 can be heated to sufficient temperature to initiate fusion reactions and compressed to sufficient
 846 areal density (bulk density ρ times fuel radius r) to trap emitted alpha particles and initiate
 847 bootstrap heating.

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849

850 **Physics of Different Types of Ignition**

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852 **Hot-Spot Ignition**

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854 Hot-spot ignition, described briefly earlier in this chapter, is the most commonly
855 discussed and best understood method for achieving ignition. Hot-spot ignition refers to the
856 creation of a small central mass of fuel that is heated to temperatures sufficient to begin efficient
857 thermonuclear burn (~ 10 keV), surrounded by a larger mass of dense but colder fuel that has
858 sufficient areal density (>300 mg/cm²) to trap alpha particles and initiate bootstrap heating.¹²

859 The primary reason for utilizing hot-spot ignition is to minimize the driver energy
860 requirements. Heating fuel to 10 keV is energy-intensive, so the goal is to use the driver energy
861 to launch a series of shocks that simultaneously coalesce and heat only a small central mass to
862 fusion temperatures, while quasi-isentropically compressing the main fuel mass as close to the
863 Fermi-degenerate limit (the minimum energy state for high-density matter) as possible. The
864 energy deposited by fusion alpha particles rapidly heats the cold, dense main fuel, causing it to
865 reach thermonuclear burn conditions. The fusion burn terminates when the rapidly heated fuel
866 mass overcomes the inertia of implosion and explodes to lower densities and temperatures where
867 fusion reaction rates rapidly decrease (hence the term “inertial confinement”).

868 In order to use minimum driver energy, it is important to compress most of the fuel near
869 the Fermi-degenerate adiabat. At least four laser pulses are required to provide the compression
870 energy in a time-dependent fashion that is consistent with this goal. More, smaller pulses—or
871 even a continuous power profile—could also be used, but the four-pulse system is the easiest to
872 control and observe experimentally.

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875 **Fast Ignition**

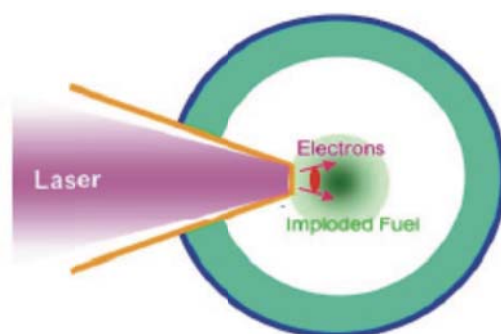
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877 In FI, ignition is separated from the compression phase. The fuel is compressed (using
878 lasers or another driver) at a lower velocity than in hot-spot ignition. The goal is to create a fuel
879 mass that has at least the 300 mg/cm² areal density required to capture alpha particles, but not the
880 DT temperature to initiate fusion burn. The energy to ignite a small portion of this compressed
881 fuel is provided by a high-intensity, ultrashort-pulse laser. For the correct conditions, the
882 thermonuclear burn propagates from this heated fuel volume into the rest of the cold, imploded
883 fuel.

884 The leading approach to fast ignition uses a hollow cone of high-density material inserted
885 into the fuel capsule so as to allow clean entry of this second laser beam to the compressed fuel
886 assembly (see Figure 2-5). The principle of fast ignition was first demonstrated at the Institute of
887 Laser Engineering in Osaka, Japan, in experiments performed on the Gekko-XII laser (Kodama
888 et al., 2002).

889

¹² R.L. McCrory, University of Rochester, presentation to the panel titled “Laser-Driven Inertial Fusion Energy: Direct-Drive Targets Overview,” on February 16, 2011.



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891

892 **FIGURE 2-5** In this version of fast ignition, a short, high-intensity laser pulse enters the cone of
893 a cone-and-capsule assembly after the fuel capsule has been compressed by an earlier pulse,
894 producing a pulse of hot electrons that initiate fusion. SOURCE: Juan Fernandez, LANL,
895 presentation to the panel titled “Inertial Confinement Fusion Targets at Los Alamos National
896 Laboratory,” May, 2011.

897
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899 Shock Ignition

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901 Shock ignition is yet another variant on the theme of slowing the main fuel implosion to
902 minimize driver energy requirements, adding one more drive element to locally heat a limited
903 quantity of fuel to thermonuclear burn conditions, and then using alpha-particle deposition to
904 propagate the burn wave into the assembled fuel mass. In shock ignition, rather than using a
905 separate, high-intensity, ultrashort-pulse laser to heat the ignited volume, a short, high-intensity
906 “spike” is added to the end of the main drive pulse shape to launch a very strong shock into the
907 fuel. This inward-propagating shock collides with the outward-propagating shock constituted by
908 the growing region of high-density fuel at the center, producing a spherical shell of fuel at a
909 much higher temperature. The principle of shock ignition has been demonstrated in experiments
910 on the OMEGA laser at LLE (Betti et al., 2007). Since the target has a smaller radius at the time
911 that the high-intensity spike is required to launch the final shock, it is energetically advantageous
912 if the laser optics can accommodate focal zooming or, alternatively, if the high-intensity spike
913 can come from a separate set of lasers with smaller intrinsic spot size. An issue that arises with
914 shock ignition is that the final, high-intensity spike exceeds the threshold for laser-plasma
915 interactions, which can interfere with the desired effect (see further discussion in Chapter 4).

916

917 Z-Pinch Ignition

918

919 Z-pinch targets need to achieve the same overall fuel parameters—that is, sufficient
920 temperature to initiate thermonuclear burn and area mass density to initiate alpha-particle
921 bootstrap heating of the remaining fuel mass. Since the targets are typically cylindrical, the
922 convergence is only two-dimensional and it is more difficult to meet the ρr criterion. Some target
923 designs work on the hot-spot ignition principle, in which a small central mass is shock-heated to
924 thermonuclear temperatures.

925 Alternatively, in magnetized-target fusion (MTF), the fuel mass is preheated by an energy
926 source (e.g. a laser beam) to place it on a higher adiabat. Field coils are placed around the target
927 to provide a seed magnetic field throughout the fuel volume. The magnetized, preheated fuel is

928 then imploded at a lower implosion velocity than is used in hot-spot ignition to minimize driver
 929 energy requirements. The magnetic field is applied to inhibit fuel cooling during the slow
 930 implosion process (i.e., inhibit cross-field transport). The higher initial adiabat allows the
 931 magnetically insulated fuel to reach thermonuclear conditions at smaller convergence ratios. The
 932 principle of MTF has not yet been successfully demonstrated. MTF is normally considered more
 933 as an attempt to find an easier path to ignition rather than as a path to high yield and high gain,
 934 but recent numerical simulations indicate that high-gain MTF is possible using cylindrical
 935 implosions with a cryogenic DT layer (Slutz and Vesey, 2012).

936 937 938 **What Determines the Degree of Fuel Burn and Gain** 939

940 Fusion yield Y scales strongly with capsule absorbed energy ($Y \sim E^{5/3}$), which implies
 941 there is a strong premium on efficiently delivering energy from the driver to the capsule. Energy
 942 must be absorbed symmetrically into the fuel to avoid instabilities. Each target design has
 943 different transport and deposition issues:

- 944 • Indirect drive (e.g., in the NIC at the NIF) requires transport of lasers through a
 945 background gas and delivery through laser entrance holes (LEH) in the hohlraum (see
 946 Chapter 4). Most of the driver energy goes to heating the hohlraum wall and the dense
 947 plasma blown off the wall, so the process is inherently inefficient.
- 948 • Direct drive simplifies transport and focusing issues, but it is critical to avoid the
 949 generation of hot electrons (which cause fuel preheat) from laser-plasma interactions.
 950 This method is more efficient because it is direct, but symmetry and deposition
 951 physics are very important.
- 952 • Z-pinchs require a direct electrical connection between driver and target through a
 953 recyclable transmission line (RTL). As the target implodes and the Z-pinch
 954 inductance increases, there may be potential loss regions. Because of the RTL, each
 955 shot requires the replacement of substantial structure.
- 956 • Heavy ions are charged particles that are susceptible to plasma instabilities when they
 957 are focused to the intensities required for ICF (>500 TW). Accelerators work best at
 958 low currents, so achieving a high power requires high particle energies, which makes
 959 their energy deposition range long. This complicates target design.

960
961 As noted above, fusion yield is calculated to scale as absorbed energy $E^{5/3}$, so delivering
 962 more energy to the target results in significantly higher yield. For the same driver energy, direct
 963 drive delivers more energy to the fuel than does indirect drive. Implicit in this yield-scaling is the
 964 fact that the increasing fusion energy output comes from burning more fuel. Burning more fuel
 965 requires compressing more fuel to near Fermi-degenerate conditions, which requires more
 966 energy to be absorbed by the target. Since most of the fuel mass is in DT at solid (ice) density,
 967 more fuel mass means targets of larger radius. Larger target radius has the additional benefit that
 968 it increases the inertial confinement time of the fuel mass (determined by the imploded fuel
 969 radius divided by the sound speed) and increases the burn-up fraction of the DT fuel
 970 disassembly. The burn-up fraction depends on the areal density of the fuel capsule:

$$971 \quad f_b = \rho r / (\rho r + \beta(T))$$

974 where $\beta(T) = 5.5\text{--}6.5 \text{ g/cm}^2$ for optimal burn conditions. For a burn-up fraction greater than
 975 about $1/3$, ρr must be greater than about 3 g/cm^2 .

976 All designs try to use driver energy efficiently; thus, they implode a cold mass of fuel
 977 isentropically and a small amount of fuel to high temperature—either by hot-spot ignition, fast
 978 ignition, or shock ignition. Instabilities can limit the propagation of burn from the ignition region
 979 to the remaining fuel. “Yield over clean” (YOC) is a measure of the deviation of experiments
 980 from ideal simulations.

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Spectrum Output

985 The fusion reaction determines the initial partitioning of energy into alpha particles, X-
 986 rays, and neutrons. The spectrum of particles hitting the IFE target chamber wall is a function of
 987 the intervening materials, whether from the hohlraum, support structures (e.g., RTLs), or
 988 chamber fill gas.

989 Indirect-drive targets have high-Z materials in the hohlraum that emit copious X-ray
 990 radiation. Xenon gas can be used to absorb these X-rays and mitigate chamber wall damage (see
 991 Chapter 4). The xenon gas will get hot, but the hohlraum is believed capable of protecting the
 992 cryogenic fuel as it transits the chamber.

993 Direct drive usually assumes a vacuum in the target chamber, because the fuel pellet
 994 cannot be thermally insulated from a hot background gas. A shroud containing helium gas at low
 995 pressure and temperature has been considered, although it presents many difficulties. Even
 996 though the target is made of low-Z materials, there are still X-rays and ions that strike the wall
 997 and deposit their energy very locally. Magnetic diversion of ions is being considered in some
 998 designs to protect the chamber wall.

999 Z-pinch reactors would have yields above 1 GJ and RTL structures in the chamber.¹³ This
 1000 can lead to debris and shrapnel. The RTLs also can contain substantial residual magnetic field
 1001 energy, which needs to be accounted for in determining which particles hit the wall. Thick, Li-
 1002 containing liquid walls can be used to protect the chamber surface from short-range ions,
 1003 neutrons, and X-rays.

1004 Heavy-ion driver concepts are tending to use liquid walls and perhaps background gases.
 1005 There do not appear to be any unique or particularly challenging aspects to the heavy-ion output
 1006 spectrum as compared with laser direct-drive or indirect-drive systems.

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 1008

Target Injection and Fabrication

1011 For energy to be produced in a fusion reactor, the target (which is the fuel source) will be
 1012 obliterated. Thus, for IFE to produce a steady flow of energy, a steady supply of new targets
 1013 must be introduced into the system. The more frequently the targets are introduced and converted
 1014 into energy, the more power is produced; and similarly, the more energy that is available in each
 1015 target, the more power is produced. It is the details of these targets, and how efficiently the

¹³ M. Cuneo et al., Sandia National Laboratories, presentation to the NRC IFE committee titled “Pulsed Power IFE: Background, Phased R&D, and Roadmap, April 1, 2011.

1016 energy is released, that distinguish the different concepts for IFE. These differences and
1017 technical challenges are discussed in detail in Chapter 4.

1018 How frequently targets can be introduced into the fusion reactor (the repetition rate) is
1019 determined by engineering practicalities of each fusion concept. The repetition rate for the
1020 concepts discussed here varies from 0.1 to 20 Hz. These values are calculated estimates; the
1021 technical challenges of delivering targets into the fusion chamber at these rates with the required
1022 precision, while preserving the integrity of the target, has been—in the absence of a
1023 comprehensive IFE program—only superficially addressed. Specific engineering concepts will
1024 require comprehensive testing to determine whether the proposed repetition rates, and
1025 subsequent power production, are feasible. Equally important is to understand whether any
1026 degradation to the configuration of the target during this injection process could reduce fusion
1027 performance below the calculated performance.

1028 Operating a fusion reactor at a repetition rate of 20 Hz will consume 1.728 million targets
1029 per day. No credible process for cost-effectively producing this number of targets has been
1030 developed. Current ICF experiments show that there is a technical path for manufacturing targets
1031 that meet critical specifications; whether this technical path is a viable method for mass-
1032 producing targets remains to be established. These considerations are discussed next.

1033

1034 **Target Injection**

1035

1036 For laser-driven IFE, the target injection process poses four challenges: accuracy and
1037 repeatability (both spatially and temporally) of target placement; ability to track the target, target
1038 survival, and clearing of the chamber. These challenges are discussed in the following
1039 paragraphs.

1040 A necessary condition for achieving the optimal energy output from each target is that the
1041 target be uniformly compressed by the laser beams. This requires the target to arrive at the same
1042 point in space and at the same instant as the multiple laser beams. For the direct-drive target, the
1043 target must be within 20 μm (rms between the centerline of laser beamlets to the centerline of the
1044 target). Concepts developed and tested as part of the High Average Power Laser (HAPL)
1045 program¹⁴ (see Box 4-2) showed that a surrogate target could be repeatedly placed within 10 mm
1046 of target chamber center, where a final engagement system does the final pointing. For the
1047 indirect-drive targets currently under development, the target is required to be within 100 μm of
1048 the focus of the laser beam,¹⁵ which appears to be within the capabilities of the system developed
1049 by the HAPL program; however, one difference between the direct- and indirect-drive
1050 approaches to fusion is that the indirect-drive approach has a higher gas pressure in the reactor
1051 chamber that may affect the repeatability of the injection process (Norimatsu et al., 2003). These
1052 are issues to be resolved in a technology development program.

1053 The second challenge is the ability to track the target to make real-time, minor
1054 corrections to the pointing of the laser beams at the target. Here technical progress was achieved
1055 during the HAPL program by demonstrating the ability to track a target moving at 5 m/s and to
1056 steer beams in real time so as to engage it with $\pm 28 \mu\text{m}$ accuracy (Carlson et al., 2007). The
1057 system has been designed assuming an injection velocity of 50 m/s.

¹⁴ J. Sethian, Naval Research Laboratory, presentation to the panel titled “The HAPL Program to Develop the Science and Technologies for Direct-Drive Laser Fusion Energy,” September 20, 2011.

¹⁵ M. Dunne, LLNL, “LIFE Target System Performance,” presentation to the panel on July 7, 2011.

1058 The third technical challenge is to preserve the target's critical specifications until the
1059 moment of the implosion. The problems are significantly different in this case for direct- and
1060 indirect-drive targets. For indirect-drive targets, the surrounding hohlraum will provide thermal
1061 protection. However, laser access to the target is through thin membranes ($<0.1 \mu\text{m}$ thick) at each
1062 end of the hohlraum, and these holes will allow a sizeable heat load (both radiative and
1063 conductive) to be delivered to the target. The radiation portion of this heat load is reduced by the
1064 presence of internal shields within the hohlraum, which will also disrupt convective cells, but the
1065 conductive heat load is unaffected and the target's temperature is calculated to rise $\sim 85 \text{ mK}$,
1066 which is less than the 100 mK ceiling specified in one system design.¹⁶ The benefit of these
1067 structures to the target's preservation is appreciable; however, this benefit comes at the cost of a
1068 complex structure that needs to be built to high precision, and this precision must be maintained
1069 during the acceleration loads that the target experiences when it is injected into the reactor. These
1070 loads to the target assembly have been calculated and are stated to be acceptable.¹⁷

1071 For direct-drive targets, target survival is the major challenge. The exact heat load to the
1072 target is strongly dependent on engineering parameters such as the gas pressure in the reactor
1073 chamber, the time the target is inside and exposed to the environment, and the temperature of the
1074 reactor; heat fluxes in excess of 1 W/cm^2 to the target will compromise the target's performance
1075 (Tillack et al., 2010; Bobeica, Ph.D. thesis, Bobeica et al., 2005).

1076 Multiple strategies are envisioned for minimizing the heat load; two possibilities are to
1077 add protective layers to the outer surface of the target and to minimize the gas pressure in the
1078 reactor (Petzoldt et al., 2002). Testing such strategies is a critical step in determining the
1079 engineering feasibility of the laser direct-drive fusion energy option.

1080 Finally, it is necessary to clear the chamber of debris between shots. In the past, there has
1081 been a tendency to minimize this problem because the other issues appear so much more
1082 daunting. However, new concepts, higher repetition rates (with incrementally more mass injected
1083 into the chamber per unit time), and the possibility of increasing the gas pressure in the reactor to
1084 improve the durability of the reactor structure (high gas pressure will reduce the X-ray and ion-
1085 induced damage to the chamber wall) complicate the process of clearing the chamber.

1086 Concepts for injecting targets for pulsed-power fusion energy are radically different and
1087 less fully developed than their laser-driven fusion energy counterparts. The signature difference
1088 is that targets are consumed at a rate of 0.1 Hz and that the target is a more massive structure (up
1089 to 50 kg) that includes transmission lines that couple the power to the target.¹⁸ Removing spent
1090 targets and installing new targets will be done using automated machinery.¹⁹ While this process
1091 is conceptually feasible, there remain substantial engineering considerations that need to be
1092 resolved to determine whether this process can be completed within 10 seconds.

1093 The heavy-ion fusion energy concepts originated as a variation of laser-driven concepts
1094 in which the driver energy is supplied by heavy ions accelerated by a linear accelerator.
1095 Subsequently, a variety of target-design concepts have been proposed: an indirect-drive design
1096 ($3\text{-}4 \text{ GeV Bi}^{+1}$); polar direct-drive design (3 GeV Hg^{+1}); and a single-sided direct-drive

¹⁶ Ibid.

¹⁷ Ibid.

¹⁸ M. Herrmann, Sandia National Laboratories, "Z-pinch Target Physics," presentation to the panel on February 17, 2011.

¹⁹ M. Cuneo et al., Sandia National Laboratories, "The Potential for a Z-pinch Fusion System for IFE," presentation to the panel on May 10, 2011.

1097 configuration (90 GeV U^{+4}).²⁰ The target-design concepts use indirect-drive, direct-drive, and
1098 single-sided direct-drive configurations. The target injection challenges are similar for heavy-ion
1099 and laser-driven fusion: the indirect-drive target benefits from the thermal shielding provided by
1100 the hohlraum, while the direct-drive target remains vulnerable to the hostile environment of the
1101 reactor chamber. Beyond these commonalities with laser-driven fusion, no target injection
1102 concept specific to heavy-ion fusion has been proposed.

1103

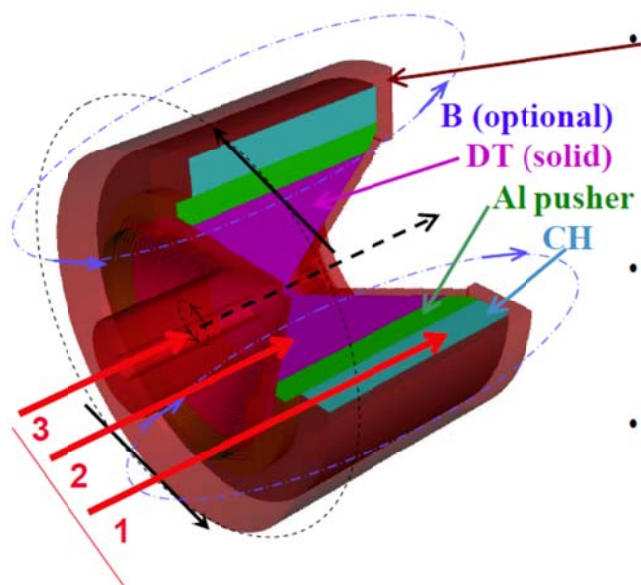
1104 **Target Fabrication**

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1106 Before the targets can be injected into the reaction chamber they must be fabricated to
1107 tight tolerances, which requires a well understood and reliable process that is suitable for mass
1108 production. The mass fabrication challenges posed for the different types of targets vary
1109 significantly, although there are technologies common to many of the targets that will benefit all
1110 concepts for fusion energy. In this section, the key challenges are outlined for the production of
1111 these targets for laser drivers, pulsed power drivers, and heavy-ion drivers.

1112 Targets proposed for each of the fusion energy concepts have equal mixtures of
1113 deuterium and tritium as the fuel. This fuel is confined in a spherical capsule for the laser-driven
1114 concepts and most of the heavy-ion concepts or in a conical “X-target” (see Figure 2-6) or
1115 cylindrical structure (see Figure 2-7) for direct-drive heavy-ion fusion and pulsed-power fusion,
1116 respectively. Fabrication of the conical and cylindrical structures appears to be straightforward,
1117 though the exact specifications are not yet well defined or tested. Fabrication of the spherical
1118 capsules is complicated—partially owing to the design and partially owing to the tight tolerances
1119 and stringent specifications. Researchers making these targets for the ICF and the HAPL
1120 programs produced targets with specifications that are acceptable for the laser-driven fusion
1121 concepts; however, it remains to be demonstrated that the fabrication process can be scaled to
1122 satisfy the requirements of an IFE program.

²⁰ B.G. Logan, Lawrence Berkeley National Laboratory, “Heavy-Ion Target Design” presentation to the panel on July 7, 2011.

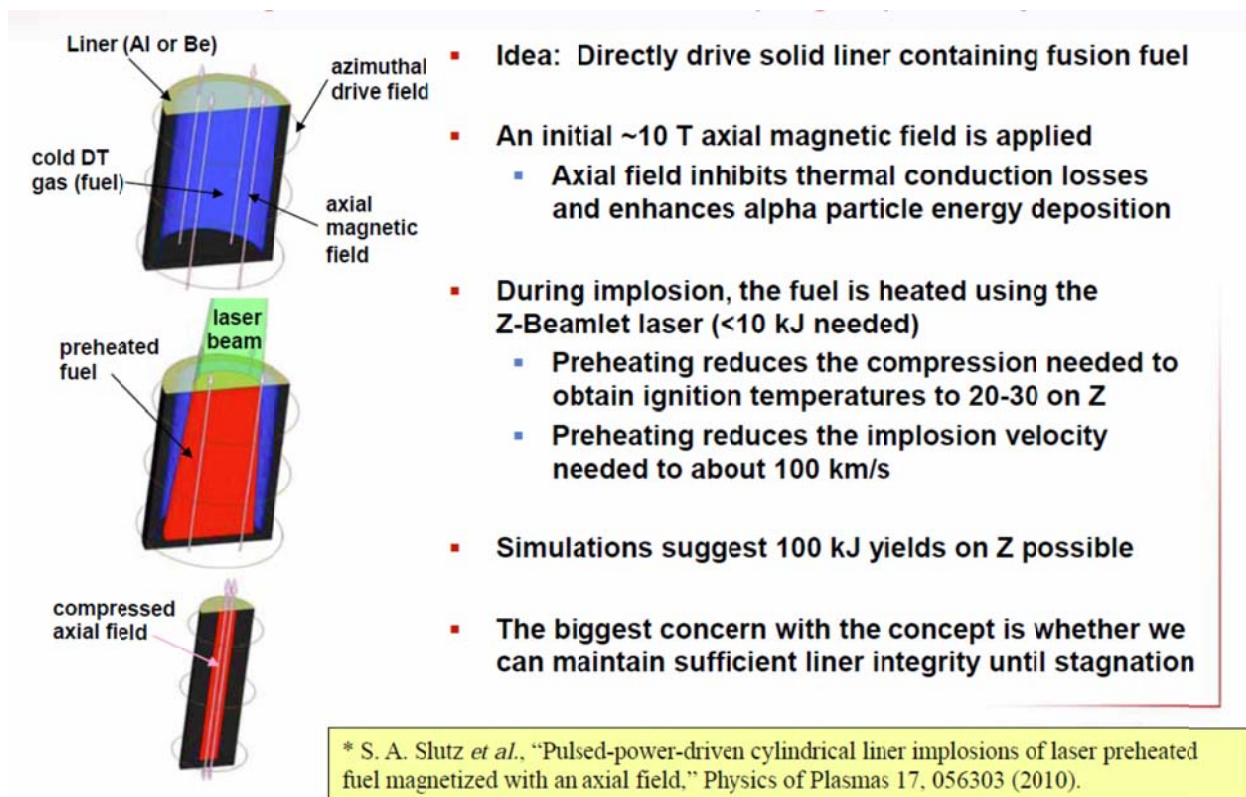


- The case can be assembled from 2 or 3 stamped metal pieces, about the mass and cost of 3 pennies. The thermal inertia of the case would protect the DT during 10 to 30 milliseconds exposure to hot chamber vapor during injection.
- At 4 deg K, the high strength of some metals for the case can withstand up to 300 MPa, allowing very high gyroscopic stability with 10^4 rps spin.
- With a high conductivity case over-coat such as aluminum, a target magnetized outside the chamber could have an L/R decay time > than the chamber transit time. The target dipole field then allows magnetic acceleration, guiding, and steering in the chamber.

Illuminated by 2-sequential annular heavy ion beam pulses for compression followed by 3rd pulse on axis for ignition—all beams from one side at high range (~ 2 g/cm² ; e.g., 90 GeV U)

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FIGURE 2-6 The heavy-ion-driven “X-target” concept. B, magnetic field; CH, plastic.
SOURCE: B. Grant Logan, LBNL, “Heavy-Ion Target Design,” presentation to the panel on July 7, 2011.



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1130

1131 **FIGURE 2-7** The cylindrical magnetized liner inertial fusion (MagLIF) target concept.
1132 SOURCE: S.A. Slutz, SNL, "Design and Simulation of Magnetized Liner Inertial Fusion
1133 Targets," presentation to the panel on May 10, 2011.

1134

1135 *Indirect-Drive Targets*

1136

1137 The indirect-drive targets proposed for laser-driven IFE (e.g., in the LIFE point design)
1138 are a modification of the target currently used at the NIF. The fundamental design is the same:
1139 DT fuel is contained inside a capsule that is supported inside a hohlraum. However, there are
1140 differences in both the capsule and the hohlraum. The capsule is a bilayered structure with an
1141 outer layer of high-density carbon (diamond) and an inner layer of low-density hydrocarbon
1142 foam. The hohlraum is elliptical (rather than cylindrical as is the NIF target) and made from lead
1143 rather than gold. Additionally, internal membranes ("shine shields") are introduced to prevent the
1144 capsule having a direct line of sight to the laser entrance holes in the hohlraum. The capsule is
1145 postulated to be manufacturable using a combination of microfluidic and vapor deposition
1146 techniques, and the DT fuel is added by drilling a hole 5 μ in diameter in the capsule and sealing
1147 it once the fuel is inserted. Cooling the target assembly liquifies the DT fuel, which is wicked
1148 into the foam layer to make a uniformly thick fuel layer. New technologies will be required to
1149 form the foam layer inside an existing capsule, and those technologies need to be consistent with
1150 a credible mass-production process.

1151

1152 *Direct-Drive Targets*

1153

1154 The direct-drive target proposed for fusion energy bears a close resemblance to the
1155 direct-drive target that is proposed for experiments at the NIF.²¹ The fusion energy target is a
1156 spherical foam capsule that is slightly larger than the NIF direct-drive target. The outer surface
1157 of the foam capsule has a fully dense plastic overcoat (to retain the fuel) and a thin reflective
1158 metallic coating to reduce the radiative heat load to the ice. Additional outer layers may be
1159 needed to provide greater protection to the target when it is injected into the reactor chamber.
1160 The DT fuel is diffused into the plastic shell and the target assembly is cooled to form the
1161 uniformly thick ice layer.

1162 The manufacturing processes for both laser-driven target designs are scalable for mass
1163 production. However, it remains to be demonstrated that these processes can achieve the
1164 production yield required for a fusion plant given the specifications that are required. At this
1165 point, such processes are near,²² but have not yet been proven for mass production. Any changes
1166 in the target design to improve the implosion physics (resulting from experiments at the NIF) are
1167 likely to be dimensional changes that can be easily accommodated by the existing manufacturing
1168 process instead of changes in configuration that would require new technologies.

1169 Two of the targets designs that are proposed for the heavy-ion driven fusion concept use
1170 indirect- and direct-drive implosion symmetries, so the manufacturing challenges are the same as
1171 for laser-driven fusion targets. A third more recently proposed target design is a single-sided
1172 direct-drive concept where liquid DT fills an X-shaped volume (two cones joined at the apex, see
1173 Figure 2-6). No production method has been proposed, nor are any tolerances proposed for the
1174 design, although it appears this target will have similar constraints and technical challenges as
1175 the other targets.

1176 The pulsed-power fusion energy targets are distinctly different from the other fusion
1177 energy targets. There are multiple designs; one is a cylinder made from beryllium and filled with
1178 cryogenic D-T gas. This target will be straightforward to manufacture and is considerably less
1179 complex than the other target designs. However, the additional components that are needed to
1180 inject this target into a pulsed-power fusion reactor must be better defined to fully evaluate the
1181 technological challenges to making the entire target assembly.²³

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Factors Most Likely to Determine the Cost of Targets

1186 It is important to appreciate that the technologies for making most of the components of
1187 the targets exist already; targets are being successfully manufactured for the existing ICF
1188 program, and with a few exceptions, any changes to the target to adapt it for energy applications
1189 appear to be technically feasible.

1190 Much of the cost of the ICF target today is due to the quality assurance process, in which
1191 each target must be thoroughly evaluated because the yield of acceptable targets is so low. Any
1192 future IFE technology program will need to evaluate whether current technologies can (1)
1193 produce a more consistent product and (2) maintain the high production yield when scaled to
1194 mass production.

²¹ P.B. Radha, University of Rochester, “Polar-Drive Target Design,” presentation to the panel on July, 7, 2011.

²² J. Sethian, NRL, “The HAPL Program to Develop the Science and Technologies for Direct-Drive Laser Fusion Energy,” presentation to the panel on September 20, 2011, and “M. Dunne, LLNL, “LIFE Target System Performance,” presentation to the panel on July 7, 2011.

²³ S.A. Slutz, SNL, “Design and Simulation of Magnetized Liner Inertial Fusion Targets,” presentation to the panel May 10, 2011.

1195 The material and production costs for manufacturing the targets appear to be acceptable
1196 and will benefit from the economies of large-scale production if a viable process is developed.
1197 The costs for developing the manufacturing process and constructing the manufacturing facilities
1198 are less predictable, with the latter depending strongly on the former. However, these are one-
1199 time costs that when amortized over the number of targets that are produced during the projected
1200 lifetime of the plant will likely be a small component in the cost of each target.

1201 A contributor to the cost of the target is the cost of the tritium fuel. Fusion energy has the
1202 appeal and requirement that tritium be bred in a reactor and be self-sustaining. Neutrons from the
1203 deuterium-tritium fusion process interact with a surrounding blanket of lithium/beryllium and
1204 produce proportional quantities of tritium. Once the plant is initially fueled with tritium, the cost
1205 of sustaining the fuel will be primarily the cost of extracting tritium from the by-products of the
1206 nuclear reaction and the cost of controlling the radiological hazards. (Deuterium, the other
1207 component of the fuel, is extracted from water.)
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1210 **Tritium Inventory Considerations**

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1212 A consideration for selecting a target production concept, and possibly even a fusion
1213 energy concept, is the amount of tritium that is required to maintain the power plant in constant
1214 operation. While tritium-breeding will allow a facility to be self-sustaining, the complexity of
1215 recovering tritium from the breeder and reactor-chamber effluent, and then refueling the targets,
1216 will scale with the complexity of the operation and amount of tritium in the facility.

1217 Minimizing the amount of tritium in a power plant was an important consideration in
1218 designing the indirect- and drive-direct targets.²⁴ More ambitious ideas were proposed for the
1219 indirect-drive concept that will require additional scientific and technical development to realize:
1220 drilling a hole in the target to add the fuel (and then resealing the hole) and achieving a
1221 uniformly thick fuel layer by suspending the fuel as a liquid within a foam layer. Combined,
1222 they would reduce the tritium inventory to less than 1 kg²⁵ by recycling tritium through the
1223 facility in less than 8 hours. The first approach adds steps to the manufacturing process and
1224 should be technically feasible; the latter approach is also technically feasible, but it is unclear
1225 whether the liquid fuel can be cooled below its freezing point and still remain a liquid, which is
1226 what has to be done to achieve the gas density required in the capsule. If this is not possible, then
1227 an alternative and lengthier process is needed to form the ice layer, which would increase the
1228 tritium inventory.

1229 Minimizing the tritium inventory was a less important consideration for developing the
1230 direct-drive target. In any case, target tritium inventory for the direct-drive targets is much higher
1231 than for the current indirect-drive configuration. About 10 times more tritium is present in this
1232 target than in the indirect-drive target. Additionally, tritium is diffused into the capsule instead of
1233 flowing through a hole, which takes 2 to 4 days because of the fragility of the target and the
1234 quantity of fuel that has to be added.²⁶ The process for forming the ice layer adds about 12 hours
1235 to the production cycle, which is the same process that the indirect-drive concept will use if it is
1236 not possible to subcool the liquid layer sufficiently to achieve the desired gas density.

²⁴ M. Dunne, LLNL, "LIFE Target System Performance," presentation to the panel on July 7, 2011.

²⁵ M. Dunne et al., LLNL, "Overview of the LIFE Power Plant," presentation to the panel on April 6, 2011.

²⁶ J. Sethian, Naval Research Laboratory, "The HAPL Program to Develop the Science and Technologies for Direct-Drive Laser Fusion Energy," presentation to the panel on September 20, 2011.

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- 1237 Two main contributors to the total tritium inventory of an IFE plant will be these:
1238 • The amount of tritium that is trapped inside the target during the target assembly
1239 phases and
1240 • The amount that is entrained in the tritium-breeding and recovery processes (from
1241 the gaseous effluent from the reaction chamber).

1242
1243 At this stage, there is insufficient information to know the optimum balance between
1244 these sources and whether the effort to minimize the amount of tritium in the target assembly
1245 process is worth the added manufacturing and technical complexities.

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Proliferation Risks Associated with Inertial Fusion Energy and with Specific Target Designs

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This chapter discusses the potential proliferation risks associated with inertial fusion energy (IFE). Many modern nuclear weapons rely on a fusion stage as well as a fission stage, and there has been discussion of the potential for nuclear proliferation—particularly vertical proliferation²⁷—in a country where an IFE power plant is sited.

We begin by providing some background on nuclear proliferation and inertial confinement fusion (ICF) and continue with discussions of several related topics: classification concerns, the relative proliferation risk associated with different target designs, weapons production in ICF facilities, knowledge transfer, other proliferation risks associated with ICF, and, finally, the importance of international engagement on this issue.

CONTEXT AND HISTORICAL PERSPECTIVE

The term “nuclear proliferation” refers to the spread of nuclear weapons knowledge, technology, and materials to countries or organizations that did not previously have this capability. Proliferation has been of increasing concern in recent years, particularly following the successful detonation of a North Korean nuclear weapon, and the signals that Iran may also be pursuing an illicit nuclear weapons program. With the breakup of the Soviet Union, special nuclear material (SNM) became available at lightly guarded facilities; it is unclear how much was lost to theft, but proliferation concerns remain. Another concern arises from the many nuclear weapons in Pakistan, and whether they are controlled adequately.

Proliferation could occur in several ways: (1) the spread of knowledge about how to build nuclear weapons to other countries, (2) knowledge of—and access to—the physical technology used to construct nuclear weapons, (3) access to the materials from which a nuclear weapon could be constructed (e.g., SNM), and (4) access to people who have been engaged in nuclear weapons technology in other nations.

Because the first nuclear weapons were built using technology that was later adapted for use in civilian nuclear power plants and the civilian nuclear fuel cycle, the role that fission power could play in proliferation has been considered for decades. An international safeguards regime to detect attempts at proliferation is currently in place and operated by the International Atomic Energy Agency (IAEA). This regime, which is based on the Treaty on the Non-Proliferation of Nuclear Weapons (NPT), involves cooperation in developing nuclear energy while ensuring that nuclear power plants and fuel cycle facilities are used only for peaceful purposes.

The risk of nuclear proliferation could also be associated with inertial confinement fusion (ICF) research facilities or, possibly in the future, inertial fusion energy (IFE) plants. For example, IFE plants and ICF research facilities provide an intense source of neutrons, which could, in principle, be used to generate ²³⁹Pu from ²³⁸U. In addition, information that could help

²⁷ Vertical proliferation refers to the enhancement of a country’s capability to move from simple weapons to more sophisticated weapons.

1290 countries develop more advanced boosted weapons or thermonuclear weapons could be gained
1291 from a thorough understanding of a fusion facility's operation.

1292 While the effect of a fission-only weapon can be devastating, the development of two-
1293 stage (both fission and fusion) thermonuclear weapons can provide much higher yield per
1294 weapon. By using an ICF facility to improve its understanding of the physics of fusion, a nation
1295 might glean information useful in transitioning its weapons program into a much more complex,
1296 modern, and threatening system. In fact, the U.S. research program in laboratory-based inertial
1297 confinement fusion has been largely funded by the nuclear weapons program, because valuable
1298 information can be learned from ICF that can otherwise be learned only from nuclear testing.²⁸

1299 Because IFE is still at an early stage as a potential energy source, international treaties
1300 related to nuclear weapons and proliferation do not clearly apply to IFE at this time. However,
1301 due to the value of IFE to the U.S. nuclear weapons program and the programs of other nations,
1302 the applicability of some treaties to ICF has been considered.

1303 The NPT does allow for laser fusion experiments, both in states that already have nuclear
1304 weapons and those that do not. As noted in 1998, this position is based on the unopposed, U.S.
1305 unilateral statement at the 1975 NPT Review Conference stating that “nuclear reactions initiated
1306 in millimeter-sized pellets of fissionable and or fusionable material by lasers or by energetic
1307 beams of particles, in which energy releases, while extremely rapid . . . are nondestructively
1308 contained within a suitable vessel . . . [do] not constitute a nuclear explosive device within the
1309 meaning of the NPT . . .” (U.S. DOE, 1995). Even so, the status of pulsed-power fusion
1310 experiments under the NPT remains unclear (Paine and Mckinzie, 1998).

1311 In the 1990s, there was discussion in the United States about whether the Comprehensive
1312 Nuclear Test Ban Treaty (CTBT) also banned the use of ICF.²⁹ Ultimately, the Clinton
1313 administration took the position that ICF is not a prohibited activity under the CTBT (Jones and
1314 von Hippel, 1998), and this position continues to be that of the Obama administration. However,
1315 some experts still debate the applicability of this treaty to ICF (Paine and McKinzie, 1998).

1316 ICF research has received a great deal of specifically directed funding in the United
1317 States in recent years, even though IFE per se has not. This research is funded primarily through
1318 the U.S. nuclear weapons program, which envisions using ICF experiments and modeling as a
1319 method of verifying codes and calculations related to the current U.S. nuclear weapons stockpile.
1320 Because many of the topics involved in ICF are related in some way to nuclear weapons, much
1321 of the work is classified. The next section provides a brief introduction to the history and current
1322 status of the classification and declassification of various ICF concepts.

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CLASSIFICATION: ICF AND IFE

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1327 The primary reason stated by the U.S. government for classifying information related to
1328 ICF is to protect information relevant to the design of thermonuclear weapons. The possibility of

²⁸ The moratorium on nuclear testing announced on October 2, 1992, by President George H.W. Bush and extended by the Clinton administration remains in effect. It was reinforced by the 1996 U.S. signing of the Comprehensive Nuclear Test Ban Treaty, which, however, has not been ratified by the United States Senate. The information gained by the nuclear weapons program is related to improving our understanding of weapons components built during the cold war, including the effects of aging on component performance.

²⁹It should be noted that the U.S. is not currently a party to the CTBT but as a signatory is bound not to act in violation of the fundamental restrictions of the CTBT.

1329 using lasers to ignite fuel was first considered by the Atomic Energy Commission (AEC) and the
1330 national weapons laboratories in the early 1960s. At that time, concerns about the potential for
1331 laser fusion weapons as well as close ties between ICF concepts and nuclear weapons design
1332 (particularly physics and simulation codes) led the AEC to classify research on ICF. The first
1333 classification guidance for inertial confinement fusion information was issued in 1964. Initially,
1334 all aspects of ICF were considered to be classified.

1335 Declassification of fusion concepts began slowly in the 1970s, and by August 1974,
1336 essentially all work with directly irradiated fusion targets was declassified. After a long pause,
1337 declassification began again in the late 1980s and continued through the early 1990s. Most
1338 notably, in late 1990, an Inertial Confinement Fusion Classification Review was requested by the
1339 Secretary of Energy with the intent of eliminating unnecessary restrictions on information
1340 relevant to the energy applications of inertial confinement fusion. The panel included
1341 representatives from the DOE national laboratories, the Department of State, the Arms Control
1342 and Disarmament Office, and other stakeholders, and the report was issued on March 19, 1991.
1343 The key panel recommendations included these: (1) “For laboratory capsules absorbing <10 MJ
1344 of energy and with maximum dimension <1 cm, all information should be declassified with some
1345 exemptions,” and (2) “Some Centurion-Halite declassification would be desirable to gain the
1346 scientific credibility needed to advance the energy mission of ICF.” (U.S. DOE, 2001). Later, on
1347 December 7, 1993, nearly all information on laboratory ICF experiments was declassified.³⁰
1348 At present, much of the information related to ICF targets has been declassified, with several
1349 notable exceptions. First, some aspects of computer codes and certain target designs remain
1350 classified, as well as the details of some historical experiments related to ICF (in particular, the
1351 Centurion-Halite program). Some aspects of classified targets are discussed in the classified
1352 Appendix F.

1353 Whether or not aspects of ICF are classified is highly relevant to the future of IFE. If
1354 essential parts of an IFE plant are classified, this could create significant complexities for
1355 commercialization. Although some commercial facilities rely on classified concepts (such as
1356 those involved in the enrichment or reprocessing of nuclear fuel), there are likely to be export
1357 controls or specific regulations involved in dealing with this situation.

1358 It is important to realize that classification or export controls could themselves indirectly
1359 cause proliferation risks if denial of information, technology, or materials causes some nations to
1360 mount covert programs or withdraw from the NPT.

1361 There are four possible scenarios for future classification of IFE concepts. The first
1362 possibility is simple—the target will be classified or other key aspects of the concept will be
1363 classified. The second possibility is that the target is unclassified, but the expertise needed to
1364 make or assess it will involve classified information or codes. A third possibility is that other
1365 parts of the plant (e.g., lasers) will be considered to be dual use and subject to export controls.
1366 Any of these three outcomes could be very troublesome at a commercial plant. On the other
1367 hand, a fourth possibility is that the target and expertise will be unclassified, and none of the key
1368 elements of the plant are subject to export controls. If this is feasible, it would be the simplest
1369 configuration and a highly desirable goal for the future commercialization of IFE.

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³⁰ Roy Johnson, LLNL, “The History of ICF Classification,” a document provided to the panel on February 24, 2011.

1372 **PROLIFERATION CONCERNS ASSOCIATED WITH DIFFERENT IFE TARGET**
 1373 **CONCEPTS**
 1374

1375 Any kind of ICF seeks to achieve thermonuclear ignition and burn. As noted previously,
 1376 this goal relates ICF to thermonuclear weapons, and for this reason ICF (whether in a research
 1377 facility or a power plant) is seen to pose some proliferation risk. However, this risk is mitigated
 1378 by the fact that (1) nuclear weapons are much larger than ICF targets, and (2) their operation
 1379 presents some different engineering challenges.

1380 Indirect-drive targets are associated with some proliferation concerns because the physics
 1381 involved is more closely related to the physics associated with thermonuclear weapons than is
 1382 the case with direct drive. In particular, the functioning of indirect-drive targets involves the use
 1383 of X-rays in the hohlraum to drive the capsule implosion. ICF using indirect drive was
 1384 declassified in 1991.

1385 In any case, the processes involved in heavy-ion deposition (for heavy-ion-driven fusion)
 1386 and the beam-plasma interactions that occur in direct-drive capsules are physically much more
 1387 remote from conditions in existing thermonuclear weapons. In addition, these processes do not
 1388 relate to any feasible design for a weapon that the panel is aware of. For these reasons, it is the
 1389 judgment of the panel that heavy-ion fusion and direct-drive fusion pose (arguably) fewer
 1390 proliferation concerns.

1391 The Z-pinch fusion concept is likewise remote from existing weapons. However, during
 1392 the cold war, the Soviet program in explosively driven magnetic implosion (MAGO) progressed
 1393 further than any other approach to pure fusion, though like all such approaches, it was still very
 1394 far from ignition (Garanin et al., 2006, Velikhov, 2008). Since the 1990s, LANL and the All
 1395 Russian Research Institute of Experimental Physics (VNIIEF) have carried out joint experiments
 1396 on MAGO (Lindemuth et al., 1995).

1397 In the future, as processing power for desktop and academic computers continues to
 1398 increase, and as knowledge of plasma physics continues to accumulate in the open literature,
 1399 many of these concerns may become less relevant, including the proliferation risk distinction
 1400 between indirect drive and other forms of ICF that might be used for IFE. Enough physics
 1401 knowledge may accumulate in the public arena that the use of indirect-drive IFE would not be
 1402 able to add much to publicly available knowledge. In such a world, codes would be classified
 1403 according to their direct use for (and calibration from) nuclear weapons, not according to the
 1404 physics that they model. However, if an IFE plant were to rely on classified codes for target
 1405 design or other operational aspects, and knowledge of these technologies could be used to gain
 1406 information about the codes' details, proliferation would be a concern.

1407
 1408 **CONCLUSION 3-1: At present, there are more proliferation concerns associated with**
 1409 **indirect-drive targets than with direct-drive targets.** However, the spread of technology
 1410 around the world may eventually render these concerns moot. Remaining concerns are likely to
 1411 focus on the use of classified codes for target design.
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1413
 1414 **WEAPONS MATERIAL PRODUCTION AT IFE PLANTS**
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1416 One of the key proliferation risks associated with any fusion plant (ICF or magnetic
 1417 confinement fusion) is that it is possible to use the plant to create materials that are essential for

1418 the construction of nuclear weapons. These materials fall into two primary categories: special
1419 nuclear materials and tritium. Both types of material can be produced without the use of fusion
1420 facilities, but commercial fusion plants may be a more convenient source for these materials for
1421 those who cannot acquire them easily in another way. The potential for the production of each
1422 type of material is discussed next.

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Special Nuclear Materials

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1427 As noted previously, it is technically possible to utilize the significant neutron flux
1428 emanating from a fusion reactor core to produce ^{239}Pu from ^{238}U . To accomplish this task
1429 covertly, it would be necessary to:

1430

- 1431 • Move quantities of uranium into the immediate vicinity of the fusion core and
- 1432 • Acquire technology for—and construct—the appropriate reprocessing facilities to
- 1433 separate the plutonium from the uranium and fission products.

1434

1435 The first task is likely to be operationally cumbersome. In addition, the transfer of large
1436 quantities of uranium into and out of a fusion power plant would likely be detectable, as such
1437 conveyance would not be a normal operation for such a plant. The development and construction
1438 of a reprocessing facility—assuming that it had not already been built and brought into
1439 operation—would also be necessary. The technology is not new, but it requires significant
1440 radiation-handling capability. The construction and operation of such a facility would probably
1441 be detectable by the current safeguards regime.

1442 Overall, the panel judges that the construction and diversion of an IFE plant in this
1443 fashion is not the simplest path for a host state to produce SNM. Research reactors and
1444 commercial nuclear plants capable of serving the same purpose (irradiation of uranium for
1445 plutonium production) exist in many nations. However, a previously built and operating fusion
1446 plant could serve as a path of opportunity for a nation interested in developing weapons. Such
1447 facilities may therefore have to be subject to inspection to assure that they would not be so used,
1448 and to IAEA safeguards in states that do not already have nuclear weapons.

1449 However, if terrorists were to seize an IFE plant, it could provide them with neutrons for
1450 the production of material to make a weapon of mass destruction. In this case, any facility
1451 capable of producing neutrons could be useful, but it is possible that no better solution would be
1452 available. Nonetheless, as noted above, an effective form of reprocessing would still be needed
1453 to isolate the plutonium.

1454 For these reasons, the panel believes that a fusion plant raises fewer proliferation
1455 concerns than a fission plant with respect to the production of nuclear materials. However, in a
1456 region free of nuclear facilities, siting of a fusion plant could increase the proliferation risk in
1457 that region if the fusion plant were totally exempt from inspection by the IAEA or other
1458 international body. A hybrid fusion-fission plant would have the proliferation disadvantages and
1459 the economic problems of both technologies.

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Tritium

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In order to fuel itself, a functioning IFE plant would likely be designed to continually breed a stream of tritium in vast amounts: about 60 kg per year for a plant of 1 GW (thermal) capacity. Tritium is not only an essential fuel for a fusion power plant, but it can also be used in part to fuel modern, boosted fission weapons or thermonuclear weapons.

The diversion of some portion of the substantial tritium stream would be relatively straightforward, but such diversion does not necessarily pose a significant proliferation threat per se. However, for a state already possessing nuclear weapons the diversion of only a few grams of tritium would be significant and would be difficult to detect. In addition, tritium can be produced in other ways if a state needs it. To date, tritium for nuclear weapons and other purposes has been produced using fission reactors.

With current technologies tritium alone, unlike SNM, cannot be used to build a nuclear weapon, and only a host state with relatively advanced capabilities would find such a stream of tritium to be useful. Indeed, for primitive nuclear weapons, tritium does not need to be used at all. However, if a significant diversion of tritium is observed, it could be a signal to the international community that the host state is considering increasing its nuclear capability to include more advanced weapons using boosting or thermonuclear burn.

KNOWLEDGE TRANSFER AT ICF FACILITIES

A second path for a potential proliferator might be the covert acquisition of key information about fusion, drawing on knowledge gained from operating a fusion facility. This path is discussed separately for research facilities and energy facilities in the following sections.

Inertial Confinement Fusion Research Facilities

Research facilities—such as the National Ignition Facility (NIF)—pose different proliferation concerns than a fully functioning inertial fusion power plant, and the concerns associated with a host country misusing a research facility are likely to be greater than those associated with a fusion power plant. A fusion research facility is designed for the purpose of increasing physics understanding on a range of topics, not for a specific function (i.e., energy production). A power plant, however, is likely to be highly specialized and not designed with the flexibility inherent in a research machine. In addition, research facility diagnostics by their nature will provide hints about the underlying physics that power plant diagnostics may not. If considered fully, the proliferation risk associated with a research facility can go beyond the physical presence of the facility in one nation or another. Research facilities may cater to a range of scientific interests beyond the needs of either the power generation community or the weapons community. For example, the NIF provides the plasma physics community with a highly effective experimental test and validation for a number of codes and theories that may indirectly or directly relate to the physics required for an understanding of thermonuclear weapons. Because the research community is intrinsically both open and international, such an improved understanding of plasma physics could provide a range of potentially useful information to a proliferator.

1510 This increase in understanding is unlikely to stop, regardless of U.S. decisions. In the
 1511 coming decades, both experiments and simulation in research facilities worldwide are likely to
 1512 surpass current U.S. capabilities. For example, continuing increases in computing speed and
 1513 understanding in the open research community could result in extremely capable physics codes.
 1514 However, it should be clear that information about physics is not the same as information about
 1515 weapons design. For a nation that has never successfully (or unsuccessfully) detonated a
 1516 thermonuclear weapon, no fusion research facility or power plant can adequately replace
 1517 experimental physics and engineering knowledge gained from nuclear testing.

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IFE Power Plants

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An IFE power plant, as noted above, is unlikely to be highly flexible, and a research facility is likely to provide more information to a potential proliferator. By the time a design is commercialized, the physics will likely have been well understood (or engineered around), and the designs of the individual components will have been optimized to the extent possible for power production. In addition, the diagnostics will be likely to be optimized for the needs of a power plant operator, not for the needs of a physicist attempting to learn useful weapons information.

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However, knowledge transfer remains a concern if an IFE power plant is deployed overseas in a country where proliferation is a concern, because local expertise will be needed to operate the plant. The plant may not yield useful information about the physics involved in the reaction, but could provide information about energies needed and other technological details that must be known to obtain ignition in a fuel pellet. Moreover, personnel would gain practical experience in handling tritium. Whether this knowledge would be greater than that obtainable in the open literature is unclear.

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CONCLUSION 3-2: The nuclear weapons proliferation risks associated with fusion power plants are real but are likely to be controllable. These risks fall into three categories:

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- Knowledge transfer,
- SNM production, and
- Tritium diversion.

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CONCLUSION 3-3: Research facilities are likely to be a greater proliferation concern than power plants. A working power plant is less flexible than a research facility, and it is likely to be more difficult to explore a range of physics problems with a power plant. However, domestic research facilities, which may have a mix of defense and scientific missions, are more complicated to put under international safeguards than commercial power plants. Furthermore, the issue of proliferation from research facilities will have to be dealt with long before proliferation from potential power plants becomes a concern.

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ICF FOR OTHER PURPOSES

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One proliferation concern associated with ICF is the potential for the development of a laser fusion weapon, as discussed briefly in the section on classification earlier in this chapter.

1556 However, owing to the size, complexity, and energy requirements of existing or planned driver
 1557 systems, the panel does not consider this to be a credible and immediate concern with respect to
 1558 current concepts for inertial fusion energy, such as laser-driven fusion energy. However, in the
 1559 distant future, advances in laser technology could change this picture.

1560 In a 1998 declassification decision, the Department of Energy (DOE) stated that “the U.S.
 1561 does not have and is not developing a pure fusion weapon and no credible design for a pure
 1562 fusion weapon resulted from the DOE investment.” (U.S. DOE, 1991). According to information
 1563 released after the cold war, the Soviet experience was similar. However, this concern might
 1564 someday materialize with currently unforeseen technology developments. For this reason and to
 1565 alleviate any current concerns, it will be important to address the possibility (or impossibility) of
 1566 pure fusion weapons in policy discussions and in the safeguards regime.

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1569 **THE IMPORTANCE OF INTERNATIONAL ENGAGEMENT**

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1571 As described in the previous sections, there are proliferation risks associated with the use
 1572 of ICF facilities around the world, and—should IFE concepts prove to be fruitful—with IFE
 1573 plants themselves.

1574 Managing proliferation, whether it is associated with fission concepts or fusion concepts,
 1575 is intrinsically an international problem. While one country may not allow the export of certain
 1576 technologies, other countries that do not consider the technology as sensitive may choose to
 1577 allow it. In addition, the result of proliferation—the successful construction of a nuclear weapon
 1578 by one more state—is international in its consequences.

1579 For this reason, preventing proliferation associated with fusion energy requires
 1580 international agreement on methods for managing the risks of the technologies involved,
 1581 including safeguards. The IAEA defines the purpose of its safeguards system as follows:

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1583 ...to provide credible assurance to the international community that nuclear material and other
 1584 specified items are not diverted from peaceful nuclear uses. Towards this end, the safeguards
 1585 system consists of several, interrelated elements: (i) the Agency’s statutory authority to establish
 1586 and administer safeguards; (ii) the rights and obligations assumed in safeguards agreements and
 1587 additional protocols; and (iii) the technical measures implemented pursuant to those agreements.
 1588 These, taken together, enable the Agency to independently verify the declarations made by States
 1589 about their nuclear material and activities.

1590

1591 This safeguards system has been in place for decades to verify compliance with the
 1592 Nuclear Nonproliferation Treaty (NPT) for fission plants and fuel cycle facilities around the
 1593 world. If new facilities that also pose a proliferation risk—such as fusion facilities—were to be
 1594 deployed around the world, it would be sensible to either include them in the current regime or to
 1595 design a similar safeguards regime for these facilities.

1596 Of course, these safeguards would need to take into account the design of a particular
 1597 fusion power plant. Although numerous design concepts have been advanced,³¹ the panel did not
 1598 see any credible, complete power plant designs. This has benefits, as it provides an opportunity
 1599 to consider “safeguardability” directly in the initial design of a fusion power plant.

³¹ See, for example, OSIRIS and SOMBRERO Inertial Fusion Power Plant Designs – DOE/ER-54100-1, March 1992, and “Inertial Fusion Energy Reactor Design Studies Prometheus-L and Prometheus-H,” DOE/ER-54101, March 1992.

1600 Early international discussions on this topic could be very helpful in reaching an
 1601 international consensus on the key proliferation concerns associated with the use of inertial
 1602 fusion power plants as well as how to manage these concerns (Goldston and Glaser, 2011).

1603

1604 **CONCLUSION 3-4: It will be important to consider international engagement regarding**
 1605 **the potential for proliferation associated with IFE power plants.**

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1608 **ADVANTAGES AND DISADVANTAGES OF FUSION PLANTS WITH RESPECT TO** 1609 **PROLIFERATION**

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1611 Proliferation is most tied to access to SNM, e.g., using enrichment processes. Richard
 1612 Meserve³² recently wrote that “There is no proliferation risk from the [fission] reactors.
 1613 Proliferation risks can arise from enrichment facilities because the technology could be used for
 1614 weapons purposes.” (Meserve, 2011) An advantage of fusion plants with respect to
 1615 nonproliferation is that SNM will not be used in the plants and SNM will not be accessible from
 1616 the waste products, as it is from fission plants. This lack of direct access to SNM is the major
 1617 nonproliferation advantage of a fusion plant.

1618 The disadvantage of inertial fusion power plants is that they allow access to knowledge
 1619 and experience with fusion, which will necessarily increase with the design and operation of
 1620 such plants. The latest nuclear weapons use fusion as a major source of the explosion energy.
 1621 These concerns were outlined in a presentation by an official (Massard, 2010):

1622

1623 As an EU [European Union] requirement, we keep a clear separation between IFE and
 1624 ‘sensitive’ weapons science (nonproliferation)

- 1625 • No use of weapons codes in the European programs
- 1626 • No benchmarking of physics code with weapons code
- 1627 • Not in favor of indirect drive capsule option in the European program for sensitivity
- 1628 issues

1629

1630 European countries have strong collaborations in ICF (for example, HiPER). The French
 1631 are building a laser fusion facility, LMJ, which is broadly similar to NIF and which will be the
 1632 most capable driver available in Europe. As a matter of policy, these programs will pursue
 1633 direct-drive ICF but do not intend to pursue indirect drive for IFE (Massard, 2010), because of
 1634 the perceived proliferation risk. The United Kingdom participates in LMJ and HiPER and also
 1635 actively participates at NIF in the United States, and in the latter context is pursuing indirect-
 1636 drive ICF.³³

1637 The Russian program in pure fusion evolved historically from the pre-1991 Soviet
 1638 nuclear weapons program (Velikhov, 2008). Its major emphasis is on magnetic confinement
 1639 fusion, which is not within the scope of this report. In ICF, two methods have received
 1640 continuing attention in Russia: laser fusion and magnetized target fusion (MTF). Although
 1641 research supporting ICF development is ongoing with smaller lasers (Kirillov et al., 2000;

³²Former Chair of the US Nuclear Regulatory Commission and chair of the IAEA safety advisory group.

³³ John Collier, UK Science and Technology Facilities Council, “Recent Activities and Plans in the EU and UK on Inertial Fusion Energy”, briefing to the NRC IFE Committee, June 15, 2011.

1642 Belkov et al., 2010), Russia currently has no laser facility comparable to NIF or LMJ,³⁴ and is
1643 unlikely to achieve laser-driven ignition in the near future. As for magnetized target fusion, the
1644 Russian MAGO concept has been widely advertised, and, as mentioned, joint work with LANL
1645 is ongoing. The proliferation risks of the MAGO MTF concept have been discussed in detail
1646 (Jones and von Hippel, 1998). Little concern about the potential for proliferation in MAGO is
1647 evident in Russian publications and policy. Indeed, in general, different countries have different
1648 classification policies.
1649

³⁴ A news report in Aug., 2011 suggests that plans for a NIF-class laser at VNIEFF are once again going forward, with commissioning expected in 2017; however the stated purpose is stockpile stewardship, not ICF (<http://english.ruvr.ru/2011/09/30/57370758.html>).

4

Evaluation of ICF Targets**SOLID-STATE LASER-DRIVEN, INDIRECT-DRIVE TARGETS****Current Status**

No laser fusion target has yet achieved ignition or breakeven,³⁵ but current understanding leaves open the possibility that given time, funding, and the existence of alternative design options with sufficient margin for ignition and a gain of one, ignition might eventually be achieved.

The current U.S. program aimed at achieving ignition, the National Ignition Campaign (NIC), lays out a path via laser indirect drive (ID), and significant progress has been made along that path, although not enough either to demonstrate success or to conclude that ignition cannot be achieved. It is the understanding of this panel that the current program plan anticipates a demonstration of ignition sometime after the beginning of FY2013, although the planning document scheduled that event for the end of FY2012. The closest Level 1 milestone as of this writing is to achieve, in FY2012, significant alpha-heating of a capsule's fuel. The expected signature of such an event is the production of at least 10^{16} D-T-equivalent neutrons. The significance of this milestone is that it would indicate that fusion bootstrapping of the ion temperature in the capsule fuel had occurred—a prerequisite to achieving fusion ignition and energy gain. The NIC Rev 5.0 target is designed to operate using indirect drive of a frequency-tripled (3ω) laser to reduce the negative effects of laser-plasma interactions (LPI) (see Box 4-1).

Recent and Upcoming Work

Recent work on indirect drive laser fusion has brought the NIC program to the point where it has transitioned from preparation for the actual ignition campaign to the campaign itself. The latter involves optimization of a set of parameterized characteristics of the target and laser system in order to achieve conditions under which ignition could be anticipated to occur; the development of these “tuning parameters” has itself been one of the areas of development, in part because most of the tuning campaigns will require the use of specially designed capsules to enable data acquisition of the type and accuracy needed for that specific campaign.

Box 4-1**Laser-Plasma Interactions**

In laser-driven ICF, the capsule implosion is driven by thermal pressure.^a Thus, the incident laser energy must be absorbed by matter and thermalized, either in the outer shell of the capsule (direct drive) or in the inner walls of the hohlraum (indirect drive), which become plasmas. The variety of LPI that take place when an intense laser pulse hits matter have been studied for more than 50 years; they have been a key limiting factor in laser ICF, and are still

³⁵ Breakeven occurs when fusion gain equals unity—that is, when the fusion energy released in a single explosion equals the energy applied to the target.

incompletely understood.

LPI that absorb and thermalize laser energy are desired. Undesirable, parasitic LPI include backscattering of laser light, which can result in loss of energy; cross-beam energy transfer among intersecting laser beams, which can lose energy or affect symmetry; acceleration of suprathermal “hot electrons,” which then can penetrate and preheat the capsule’s interior and limit later implosion; and filamentation, a self-focusing instability that can exacerbate other LPI. LPI are worse at longer laser wavelengths, so all modern drivers currently operate in the “blue” (3ω Nb:YAG at 353 nm) or ultraviolet (KrF at 248 nm). Moreover, lasers can be modulated so as to substantially ameliorate parasitic LPI by spectral broadening, spatially incoherent filtering, and/or polarization diversity, and great progress has been made over several decades on all the main kinds of laser drivers on such beam smoothing.^b Since LPI are threshold effects, target designers attempt to keep laser intensities below the threshold of major harm. However, neither fundamental understanding nor simulation are good enough to do so a priori; well diagnosed experiments remain essential for LPI control.^c

LPI are currently important in the NIC indirect-drive targets. Overall, backscattered light losses appear to be 10-15 percent of the incoming laser energy; however, the inner beams backscatter more because of their greater path length in the hohlraum plasma. Stimulated Raman scattering (SRS) of the inner beams appears to play a significant role in causing drive asymmetry and hohlraum temperature deficits.^d The asymmetry has been controlled by the use of cross-beam energy transfer mediated by Brillouin scattering, but fundamental understanding and simulation of this effect are incomplete, and its repeatability has not been established experimentally. Experiments so far are said to indicate that hot electrons are below the design threshold, but more diagnostics are needed, because hot electrons, if actually present, could explain the currently observed anomaly in capsule adiabat. Furthermore, other laser-produced sources of preheat, such as gold M-band emission, will require quantification in this new cross-beam environment.

Rapidly increasing computer performance has enabled LPI calculations that were unimaginable just 12 years ago, but full-scale National Ignition Facility (NIF) simulations remain beyond reach.^e The Lawrence Livermore National Laboratory (LLNL) typically performs single- or multiquad simulations using pF3D on the largest advanced simulation and computing (ASC) platforms. Improvements in hohlraum modeling have changed plasma conditions and the location of backscatter in LPI simulations, bringing them into better agreement with measurements. Recent simulations show that overlapping quads and spatial nonuniformities act to increase laser reflectivity. Simulations have suggested potential ways to mitigate the effect of overlap beam intensity on SRS, including changing the hohlraum aspect ratio and changing the pointing of inner cone quads. Substantial computational and experimental resources are being devoted to LPI issues within the NIC.

LPI for direct-drive targets is under experimental and theoretical study at LLE;^f the most important effect appears to be cross-beam energy transfer, which results in 20 percent energy losses in capsule experiments on OMEGA. The relatively short beam paths in coronal plasma suggest that other LPI, and hot electrons, may be controllable in the extrapolation to ignition targets for direct drive, though most of the key experiments remain to be done. However, the greater laser intensities needed for shock ignition may cause harmful LPI; this must be studied. OMEGA EP^g will be an important platform for studying direct-drive LPI issues at IFE-relevant plasma scale lengths. NRL is performing complementary LPI experiments at 248 nm on Nike.^h Two-plasmon decay experimental data seem to agree with thresholds calculated using simple

plane-wave-based threshold formulas, confirming the classical wavelength scaling. In direct drive, the initial target aspect ratio can be modified to limit the intensity and mitigate LPI risk at the penalty of greater sensitivity to Rayleigh-Taylor hydroinstabilities.

Increased LPI intensity thresholds and greater hydrodynamic efficiency for short wavelengths should combine to give better overall stability in direct-drive implosions. The Naval Research Laboratory (NRL) baseline shock ignition target is above the two-plasmon decay threshold during compression.ⁱ Extending the Nike laser to 20 kJ would provide a useful capability to study LPI and hydrodynamics at 248 nm in IFE-relevant scale-length plasmas and compare them with OMEGA EP and NIF data.

Plasma physics, including LPI, involves many degrees of freedom on a huge range of length scales; moreover, nonlocal propagation by electromagnetic fields and fast electrons are important. For these reasons, a priori simulation of a full-scale target will be impossible for the foreseeable future, although impressive simulations are now feasible for fundamental processes and small-scale regions. Future development of subgrid and mesoscale modeling on full-scale systems would help to understand the experiments and support better target design, but would require a large effort to create and perfect.

^a Radiation pressure of the laser light itself is too small by many orders of magnitude.

^b David Montgomery, LANL, “Overview of laser plasma instability physics and LANL understanding,” presentation to the panel on September 21, 2011.

^c Mordecai Rosen, LLNL, “Understanding of LPI and its impact on indirect drive,” presentation to the panel on September 21, 2011.

^d Ibid.

^e Denise Hinkel, LLNL, “State of the art for LPI simulation,” presentation to the panel on September 21, 2011.

^f Dustin Froula, LLE, “Laser-plasma interactions in direct-drive implosions,” presentation to the panel on September 21, 2011.

^g OMEGA EP (extended performance) is an addition to OMEGA and extends the performance and capabilities of the OMEGA laser system. It provides pulses having multikilojoule energies, picosecond pulse widths, petawatt powers, and ultrahigh intensities exceeding 10^{20} W/cm².

^h Andrew Schmitt, NRL, “Assessment of understanding of LPI for direct-drive (KrF),” presentation to the panel on September 21, 2011.

ⁱ Liu and Rosenbluth, 1976.

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Four key input variables are to be optimized in the NIC tuning campaigns:

- The implosion adiabat (usually designated α), which strongly affects the resistance of the capsule to implosion;
- The implosion velocity V ;
- The amount of capsule material involved in mixing across the single interface characteristic of this class of capsule designs, M ; and
- The overall shape of the implosion, which is characterized by a dimensionless parameter S .

These tuning campaigns are expected to use what are termed “keyhole” targets, backlit gas capsules, “symcap” capsules, and reemission capsules. Ignition is neither expected nor desired in these types of capsules, although tritium-hydrogen-deuterium (THD) capsules, which are intended for use in many of the preignition integrated experiments, utilize the ignition design but incorporate less DT thermonuclear fuel in favor of the less reactive HD. The use of THD capsules is expected to allow collection of data with which to confirm or calibrate calculations of

1700 the nuclear performance of the optimized implosion system (laser pulse + hohlraum + capsule
1701 design). Calibration of the nuclear diagnostics is planned using capsules of the so-called
1702 “exploding pusher” design.

1703 The work mentioned thus far has all been accomplished at the NIF facility at LLNL.
1704 Additional preparations for optimization and testing of ignition capsules have been carried out at
1705 other laser facilities, notably the OMEGA laser at the University of Rochester’s Laboratory for
1706 Laser Energetics (LLE). One aspect of this work has investigated some of the problematic
1707 aspects of LPI. Experiments at LLE have also facilitated the development and porting of
1708 diagnostics to the NIF and have provided data on the operation of noncylindrical, “rugby”
1709 hohlraums;³⁶ experiments are planned to provide similar data on the efficacy of “P2”³⁷ laser
1710 entrance hole (LEH) shields.

1711 If ignition can be achieved on NIF, target simulations presented to the panel suggest that
1712 optimization of the tuning parameters and increases in the driver energy could result in gains of
1713 between 50 and 100 at some future facility.

1714

1715 **Evaluation and Discussion of Remaining R&D Challenges**

1716

1717 It is too early in the experimental campaign to evaluate the performance of the NIC
1718 ignition target design. However, information already in hand does indicate some potential
1719 problem areas, which could become showstoppers. They are discussed individually below.

1720

1721 **Implosion Velocity**

1722

1723 Perhaps the most critical discrepancy is that the measured implosion velocity of
1724 nonoptimized capsules is ~10 percent lower than the calculated velocity, even early in the
1725 implosion. The fact that related quantities, such as capsule bang time, are likewise delayed
1726 compared to expectations confirms the interpretation of the velocity measurements. Possible
1727 explanations offered at the time the panel received its briefings are that the calibration of the
1728 hohlraum temperature measurement (Dante X-ray flux diagnostic) was incorrect, or that the
1729 opacity of the Ge dopant in the capsule wall (to reduce early-time heating of the interior portions
1730 of the capsule) was higher than expected.

1731 Plans are in place to explore these hypotheses by checking the calibration in question and testing
1732 capsules without that dopant for comparison.

1733 The principal means available to increase the implosion velocity is to increase the laser
1734 drive energy. Greater drive energy would, however, also increase the preheating from LPI,
1735 which, as discussed below, does not appear to be well understood. A path forward is thus not
1736 guaranteed.

1737

1738 **Implosion Symmetry**

1739

³⁶Rugby hohlraums are shaped not like a cylinder but like a rugby ball, with a wall having a tapered curve.

³⁷ ‘P2’ refers to the type of departure from sphericity that the shields are intended to reduce. A nearly spherical shape with azimuthal symmetry is often represented mathematically using Legendre polynomials, and “P2” is the standard means of referring to the second Legendre polynomial, which is needed to describe a shape that has been described as a “sausage.”

1740 The panel was told that there are some concerns about early-time imprinting of drive
1741 asymmetries based on observations of reemission targets. Furthermore, the overall implosion
1742 symmetry of baseline targets was routinely more prolate than predicted. Acceptable symmetry
1743 was obtained using interbeam energy transfer between outer and inner laser cones, but at present
1744 this process has not been successfully incorporated into the design simulations used to predict
1745 target performance. The consensus of the panel is that this situation may be a further indication
1746 of unknown LPI processes in the hohlraum or of other predictive inadequacies.

1747

1748 **Mix**

1749

1750 The prediction of mix across shocked interfaces and during convergent implosions has
1751 been a very active and controversial area of research in many technical communities for many
1752 years. Approximate simulations of mix are possible and are routinely included in some target
1753 simulations, but the calculated mix—and therefore its calculated effects—is recognized to be
1754 unreliable. Moreover, data to validate calculations of the consequences of mix is thus far
1755 unavailable. It is therefore planned to compensate for the effects of mix empirically—that is, it is
1756 planned to design and engineer for sufficient margin in ignition conditions and gain to
1757 compensate for whatever degradation the mix may cause.

1758 The lack of a definitive, quantitative understanding of the origins and evolution of mixing
1759 has raised concerns that isolated bumps and defects in the capsule shell could give rise to spikes
1760 of wall material that would penetrate into the central fuel region. The potential for such an
1761 occurrence clearly is related to the precision of target fabrication; some target fabrication
1762 technology issues are discussed below.

1763

1764 **Implosion Adiatat**

1765

1766 Measurements indicate the existence of disparities between the calculated and actual
1767 adiabat on which NIF capsules implode. Some workers have postulated that the disparities are
1768 due to inaccuracies in tabulated plastic ablator (CH) release isentropes, but there appears to be no
1769 technical evidence to support this hypothesis.

1770 LLNL briefings to the panel conveyed conviction that hot electron preheat from LPI in
1771 the NIF target has been adequately anticipated and that the implosion adiabat of the fuel can be
1772 managed by controlling shock heating. Nevertheless, the uncertainties concerning LPI processes
1773 within a target hohlraum (discussed below) and the strong sensitivity of a capsule's gain to
1774 preheat make the understanding and management of a capsule's implosion adiabat an area of
1775 concern to the panel.

1776

1777 **Laser-Plasma Interactions**

1778

1779 LPI diagnostics on an ID target assembly can only sample the small solid angle of light
1780 that is backscattered out of a hohlraum's laser entrance holes. The processes occurring inside the
1781 hohlraum, including those that can produce hot electrons, are difficult to observe. These
1782 circumstances significantly decrease the effectiveness of efforts to ascertain the adequacy of
1783 simulations of LPI.

1784 Initial experiments on the OMEGA laser have shown disparities between modeling for
 1785 both vacuum and gas-filled rugby hohlraums. Scattering of the inner beams entering a hohlraum
 1786 is reported to be greater than predicted, providing specific evidence of simulation inadequacies.
 1787 Current simulations approximate LPI using inverse Bremsstrahlung energy deposition models in
 1788 which the power balance of the beams is input by the user, although rad-hydro modeling has
 1789 apparently been improved through the use of nonlocal electron transport models and detailed
 1790 configuration analysis (DCA). Cross-beam transfer is estimated via analytic models. There is a
 1791 fluid model for LPI, called PF3D, which includes approximate models of kinetic effects; the use
 1792 of similar models might improve LPI simulations for laser fusion applications.

1793 It appears to the panel that the current state of understanding and simulation capability of
 1794 LPI presents a significant risk to both the NIC and the credibility of any indirect-drive IFE
 1795 design concept, such as the Laser Inertial Fusion Energy (LIFE) initiative. The effects of LPI
 1796 may be a central issue, contributing to observed disparities between measured and calculated
 1797 implosion entropy, velocity, and shape in the NIC.

1798

1799 **Capsule Fabrication**

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1801 There is extensive experience in fabrication of NIC-style targets, and there is a high
 1802 likelihood that the capsule and hohlraum system can be made to the desired specifications.

1803

1804 **CONCLUSION 4-1: The national program to achieve ignition using indirect laser drive**
 1805 **has several physics issues that must be resolved if it is to achieve ignition.** At the time of this
 1806 writing, the capsule/hohlraum performance in the experimental program, which is carried out at
 1807 the NIF, has not achieved the compressions and neutron yields expected based on computer
 1808 simulations. At present, these disparities are not well understood. While a number of hypotheses
 1809 concerning the origins of the disparities have been put forth, it is apparent to the panel that the
 1810 treatments of the detrimental effects of LPI in the target performance predictions are poorly
 1811 validated and may be very inadequate. A much better understanding of laser-plasma interactions
 1812 will be required of the ICF community.

1813

1814 **CONCLUSION 4-2: Based on its analysis of the gaps in current understanding of target**
 1815 **physics and the remaining disparities between simulations and experimental results, the**
 1816 **panel assesses that ignition using laser indirect drive is not likely in the next several years.**

1817 The NIC plan—as the panel understands it—suggests that ignition is planned after the
 1818 completion of a tuning program lasting 1-2 years that is presently under way and scheduled to
 1819 conclude at the end of FY2012. While this success-oriented schedule remains possible, resolving
 1820 present issues and addressing any new challenges that might arise are likely to push the timetable
 1821 for ignition to 2013-2014 or beyond.

1822

1823 **CONCLUSION 4-3: Ignition of a laser-driven, indirect-drive capsule will provide**
 1824 **opportunities for follow-up work to improve understanding of the potential for IFE.**

1825

- 1826 • If ignition is achieved with indirect drive at NIF, then an energy gain of 50-100
- 1827 should be possible at a future facility. How high the gain at NIF could be will be
- 1828 better understood by follow-on experiments once ignition is demonstrated. At this
- 1829 writing, there are too many unknowns to project a potential gain.

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- Achieving ignition will validate the assumptions underlying theoretical predictions and simulations. This may allow a better appreciation of the sensitivities to parameters important to ignition.

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USE OF LASER-DRIVEN, INDIRECT-DRIVE TARGETS IN A PROPOSED IFE SYSTEM

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The proposed—and de facto—baseline model for a laser ID power plant is the LIFE initiative of LLNL. The discussions in this section are therefore based on that design as presented to the panel.

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The current target design for LIFE was derived from the current baseline NIC design, with subtle but distinct differences. Modification was necessary to increase the calculated gain for IFE. Other modifications were to enable rapid, affordable fabrication in bulk, because the current plan for LIFE envisions firing approximately 1 million targets per day. The developers of LIFE plan to accommodate errors in the calculated target performance by adopting a design that is calculated to produce 125 percent of the gain for which LIFE was designed. The 25 percent surplus gain is viewed as a margin that would be eroded by the combined effects of inaccuracies in target design, fabrication, insertion, drive (shape, intensity, smoothing, and aiming), and LPI.

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As discussed above, in evaluating the current NIC target, issues relating to the target implosion velocity, implosion symmetry, mix, the implosion adiabat, and LPI must be addressed. In spite of the modifications to the NIC target design that adapt it for use in LIFE, sufficient similarities persist that the preceding issues apply fully, unless and until optimization and other research conducted under the NIC program lead to a favorable resolution of the underlying uncertainties. The differences between the NIC and LIFE targets also raise additional issues, as discussed below.

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Modifications to Increase Gain

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The design approach to increasing the gain of the IFE capsule stems from an approximate analytical expression in which capsule yield is proportional to $E_{\text{capsule}}^{5/3}$, where E_{capsule} is the energy absorbed by the capsule. The strategy is to increase the implosion energy primarily by increasing the drive temperature in the target hohlraum. The drive temperature is increased by increasing the laser driver energy and decreasing losses. The laser energy is to be increased from a maximum energy of 1.8 MJ at NIF to 2.2 MJ for LIFE.

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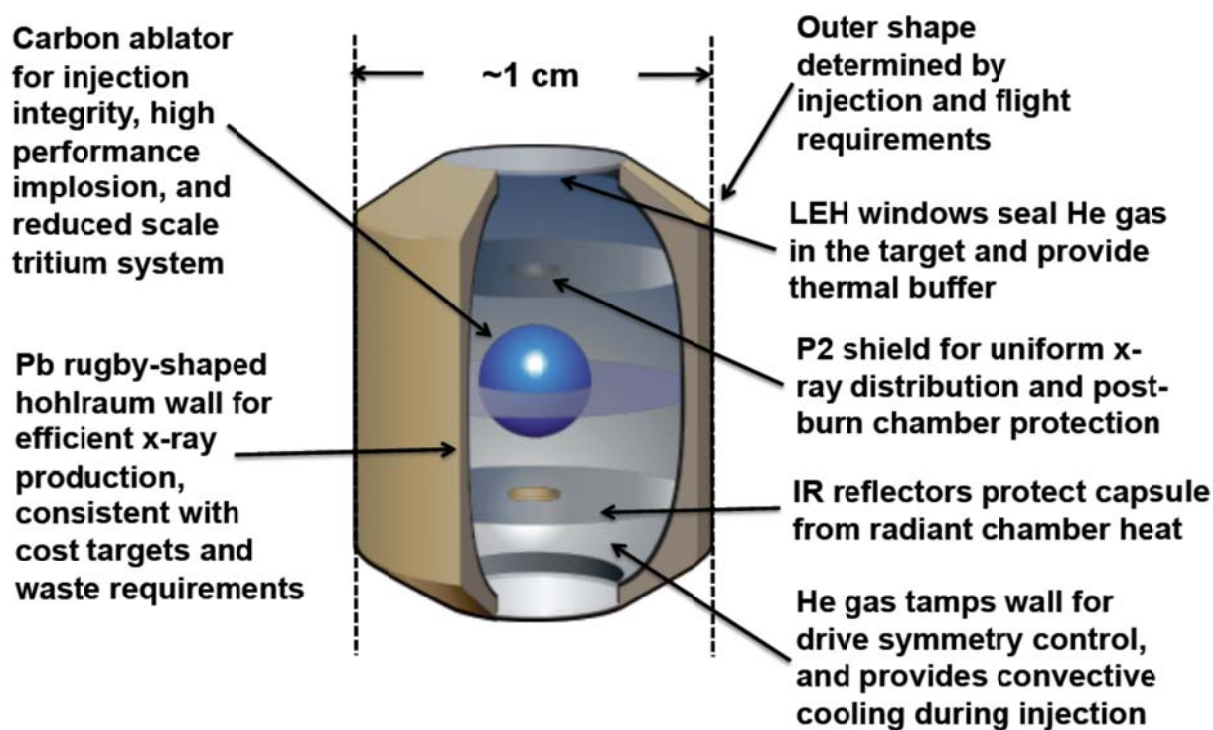
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A hohlraum shaped like a rugby ball has been designed to more efficiently partition the drive energy; the redesign includes reducing the case-to-capsule diameter ratio to 2.0-2.4. The energy lost by reradiation from the hohlraum is to be reduced by the use of P2 LEH shields, and the conversion of absorbed energy to implosion energy is to be increased by using a high-density carbon (HDC) shell to increase the ablation efficiency. An illustration of the LIFE target design is shown in Figure 4-1.



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1873

1874 FIGURE 4-1 The LIFE target design. Modifications from the NIC target design include the
1875 curved (“rugby”) inner wall of the hohlraum, the high-density carbon ablator, the LEH shields,
1876 and the P2 shine shields. SOURCE: Mike Dunne, LLNL, presentation to the panel on July 7,
1877 2011.

1878

1879

Modifications for Production Operation

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1881 The proposed manufacturing process of the LIFE target is a significant extension of the
1882 well-proven process for manufacturing targets for the NIC.

1883

Capsule Fabrication

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1885
1886 There is extensive experience in capsule fabrication, and it appears likely that the capsule
1887 can be made to the desired specifications. The technical challenges are (1) to demonstrate the
1888 formation a uniformly thick, low-density (20 mg/cc) foam wall inside the diamond shell using a
1889 technique that is suitable for mass production and (2) a cost-effective manufacturing process that
1890 can process more than 1 million targets per day through multiple steps where each target is
1891 individually handled. Proponents assert that automation can achieve the required throughput for
1892 an indeterminate capital and development cost; the bigger issue is whether the manufacturing can
1893 be done for the required per item cost (estimated to be in the range of 20-40 cents).³⁸

1894

1895 The method proposed for forming a uniformly thick fuel layer is a radical departure from
1896 the method used for making targets for the NIF. The reason for this new concept is to reduce the
1897 time required to form the fuel layer and thereby reduce the tritium inventory for the power plant.
The design is for the fuel layer to be maintained as a supercooled liquid at a temperature

³⁸ D.T. Goodin, General Atomics, presentation to the main IFE committee on January 29, 2011.

1898 sufficiently below the freezing point to achieve the required vapor pressure. The thickness
 1899 uniformity of the fuel layer is expected to be provided by the 20 mg/cc CH foam wall, the
 1900 interfacial liquid surface tension, and a controlled thermal profile along the surface of the
 1901 hohlraum. This process has to be demonstrated. A critical technical milestone is to demonstrate
 1902 that the DT liquid can be supercooled sufficiently to achieve the required vapor pressure, a
 1903 property that has not been observed in cryogenic fluids.³⁹ A second technical challenge will be to
 1904 preserve the uniformity of the liquid fuel when the capsule is accelerated to a velocity of 250 m/s
 1905 into the target chamber. The low mechanical stiffness of the low-density foam and the low
 1906 viscosity of the liquid will make the uniformity of the fuel layer thickness susceptible to the high
 1907 acceleration loads.

1908 Neither of the traditional methods of introducing fuel into the capsule—a capsule fill tube
 1909 or diffusion filling—is feasible for power plant targets. A method would have to be developed to
 1910 seal the capsules with a plug of some appropriate material after filling them with DT.

1911

1912 **Hohlraum**

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1914 The rapid capsule insertion necessary for a power plant will require structurally rigid
 1915 support for the capsule and the LEH shields. The hohlraum-capsule structure is a delicate and
 1916 intricate design with tight assembly tolerances on how precisely the capsule needs to be
 1917 positioned inside the hohlraum. In addition, there are two internal shine-shields that need to be
 1918 positioned precisely inside the hohlraum using a low-mass support structure so that neither the
 1919 thermal profile nor the x-ray radiation flux within the hohlraum is excessively perturbed. Further
 1920 work is required to define a construction that meets these requirements and will also survive the
 1921 high acceleration loads experienced when the assembly is injected into the target chamber.

1922 The hohlraum walls in the LIFE design are to be of a lead alloy that is optimized for high
 1923 opacity at the capsule drive temperature. Current hohlraums are constructed either entirely of
 1924 gold, or of gold-plated uranium. The latter are impractical for a high production rate. As an
 1925 example, a firing rate of 10 Hz translates to 8.6×10^5 capsules fired per day. With a hohlraum
 1926 mass of 3 g, 2.6 metric tons of lead must be collected and recycled per day. Using lead rather
 1927 than solid gold will reduce both the startup cost and the security requirements for the crucial
 1928 processes of hohlraum material recycling and target fabrication.

1929

1930

Evaluation

1931

1932 In evaluating the current NIC target, issues relating to the target implosion velocity,
 1933 implosion symmetry, mix, the implosion adiabat, and LPI were discussed above. The
 1934 modifications to the NIC target design that adapt it for use in LIFE leave it fully vulnerable to the
 1935 issues surrounding the performance of the NIC capsule, unless and until optimization and other
 1936 research conducted under the NIC program lead to a favorable resolution of the underlying
 1937 issues. The differences between the NIC and LIFE targets and drives also raise additional issues,
 1938 which are discussed below. This section concludes with an evaluation of the robustness of the
 1939 LIFE target design.

1940

³⁹ Different IFE target designs exist for different methods of achieving compression. Only one target design proposes supercooled DT liquid. If this step turns out to be physically impossible, alternative designs will be explored.

1941 **Modifications to Increase Gain**

1942

1943 The credibility of the effectiveness of the target design changes from NIC to LIFE is
 1944 directly related to obtaining and understanding the desired performance of the NIC Rev 5.0
 1945 design and understanding its operation. The seriousness of the issues discussed in this section
 1946 can be expected to become more apparent as the ignition campaign unfolds. Many of these
 1947 changes are scheduled for study on OMEGA, NIF, or both.

1948

1949 **Capsule Implosion**

1950

1951 The system modifications to increase the capsule drive are primarily intended to increase
 1952 the energy of the imploding capsule; the implosion velocity is one indicator of this energy.
 1953 The planned increase in the energy of the LIFE lasers should provide the most direct means of
 1954 increasing the energy of an imploded capsule. The outlook for carrying out this plan is clearly
 1955 independent of the target design, but any compromise in achieving this energy goal could
 1956 severely reduce the likelihood of achieving sufficient gain for a power plant to be feasible.
 1957 Calculations indicate that a redesign of the target hohlraum from the cylinder shape used thus far
 1958 at NIF to a rugby shape can increase the drive temperatures for the enclosed capsule. However,
 1959 initial experiments on the OMEGA laser using this hohlraum shape have shown disparities
 1960 between the expected and measured temperatures. This trend was observed for both evacuated
 1961 and gas-filled hohlraums. The disparities are not well understood and could be caused by
 1962 increased importance of missing models of laser-plasma interactions or by something as simple
 1963 as inadequate zone resolution. Although independent codes are used at the various laboratories,
 1964 they tend to have similar models. Until a better understanding of the disparities between
 1965 modeling and experiments on rugby hohlraums is achieved, there will be concerns that the
 1966 needed drive temperatures might not be obtained.

1967 Data appropriate for validating calculations of the temperature distribution and history in
 1968 a rugby hohlraum are not yet in hand. Aspects of the calculations needing validation include the
 1969 behavior of hohlraums with Pb walls, the radiation flow and hydrodynamic effects of P2 LEH
 1970 shields, and the radiation hydrodynamics of a target utilizing a 2.0-2.4 case diameter:capsule
 1971 diameter ratio. Such data must be acquired to attain confidence in predictions of target operation
 1972 for LIFE.

1973

1974 **Mix**

1975

1976 The HDC to be used in the LIFE outer shell is a more complex material than the CH it is
 1977 replacing; it exhibits a microcrystalline structure and is described by a complicated phase
 1978 diagram. Because three-dimensional, directional irregularities are intrinsic to a microcrystalline
 1979 structure, the potential for HDC to affect the hydrodynamic stability of the capsule requires
 1980 further study.

1981

1982 **LPI**

1983

1984 The modifications of the LEH and the addition of the P2 shields to the NIC hohlraum
 1985 create the potential for the LPI issues discussed above to be exacerbated by the use of a rugby
 1986 hohlraum. Some increased effect could also be expected from the approximately 20 percent

1987 increase in laser power. The introduction of LEH shields with the rugby hohlraum may increase
 1988 the mass of blown-off material in which LPI occur. Resulting changes in LPI phenomena may
 1989 also change the implosion adiabat for the capsule.

1990

1991 **Modifications for Production Operation**

1992

1993 Target Fabrication

1994

1995 A target of this design has not yet been made, and new technologies will be required to
 1996 make it. Only once the target is demonstrated to meet the specifications can the feasibility of
 1997 mass-producing these targets for the desired cost be accurately assessed.

1998 The plan to form the outer fuel layer of a LIFE target capsule by wicking liquid DT into a
 1999 layer of nanoporous foam is a radical departure from the method used for making targets for the
 2000 NIF. It will be necessary to demonstrate the formation of a uniformly thick, low-density (20
 2001 mg/cc) foam wall inside the HDC shell using a technique that is suitable for mass production.
 2002 The efficacy of the planned smoothing mechanisms, as well as the ability to create and maintain
 2003 the required thermal profile on the hohlraum through target insertion must also be demonstrated.

2004 Other specific issues of concern include the need to eliminate the polishing step for the
 2005 HDC shell and the significant length of time (approximately 2 days) involved for crucial
 2006 manufacturing steps (CVD deposition of the HDC and etching to remove the silicon mandrel)
 2007 (Biener et al., 2009). The hohlraum-capsule structure is a delicate and intricate design with tight
 2008 assembly tolerances on how precisely the capsule and two P2 LEH shields need to be positioned
 2009 inside the hohlraum using low-mass support structures so that neither the thermal profile nor the
 2010 X-ray radiation flux within the hohlraum is excessively perturbed. A construction method that
 2011 meets these requirements is not yet available.

2012 It would be important to the successful operation of the targets that the original
 2013 specifications for the composition and uniformity of the lead mixture used to make the hohlraum
 2014 walls be consistently maintained. The use of a “salted” Pb solution or alloy for the body of the
 2015 target hohlraum would probably complicate the recycling process for that material. When it exits
 2016 the reaction chamber, this material will have to be cycled through a full sequence of phases,
 2017 proceeding rapidly from a solid to a plasma and then somewhat more slowly to a gas and a
 2018 liquid. The composition of this liquid Pb mixture is unlikely to be uniform on the micron scale,
 2019 and some portion of the other target components would also be present.

2020 Whether fabrication to sufficiently tight specifications can be done for an acceptable per-
 2021 item cost is an important question. It should be apparent from the discussion above that there are
 2022 numerous technical challenges associated with developing an effective fabrication technology.
 2023 However, the fuel costs for an inertial fusion power plant are much larger than is typical for the
 2024 power industry,⁴⁰ so there is very little financial room for compromise. As currently envisioned,
 2025 a viable technology must be capable of producing approximately 1 million targets a day through
 2026 multiple steps in which each target is individually handled. Automation might achieve the
 2027 required throughput by eliminating individual handling, but the associated capital and
 2028 development costs are not known. The critical point from the standpoint of target design is that a
 2029 compromise on any target specification or other aspect of fabrication quality would be likely to
 2030 significantly reduce target gain.

⁴⁰ The LIFE point design puts fuel costs at nearly 28 percent of the cost of electricity, about the same as the laser costs. From Tom Anklam, LLNL, presentation to the main IFE committee on January 31, 2011.

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Additional Considerations

The combination of extreme conditions that exist in a power plant reaction chamber and the very tight specifications that must be maintained for an IFE power plant to function result in an unusually tight coupling between the target design and some of what would typically be considered the separable engineering aspects of a power plant design. For the LIFE concept, the target insertion mechanism and the protection of the reaction chamber's laser windows fall into this category.

Target Insertion

The target must be positioned precisely at the desired location and in the desired alignment at the specified instant in time to uniformly drive the implosion. Positioning tolerance within approximately 1 cm of the optimum position was demonstrated as part of the High Average Power Laser (HAPL) program (see Box 4-2) using a smaller target than the proposed LIFE target. However, the conditions of the HAPL demonstration did not include transport through hot Xe gas, which will be present in the LIFE chamber to help protect the walls. Turbulence in this gas due to the ~10 Hz firing rate is inevitable, and its effect on target positioning is currently unknown. The LIFE targets are to be inserted into the reaction chamber in a manner that is most reminiscent of a bullet, requiring an acceleration of 400-500 g to reach the required 250 m/s velocity. This acceleration places very great demands on the technology for target fabrication.

The nominally low-mass supports for the P2 LEH shields and for the capsule itself must survive target acceleration with a sufficiently predictable geometry that their position satisfies tight specifications. It is even more important that the geometry of the capsule layers be as designed at shot time. The low mechanical stiffness of the low-density foam and the low viscosity of the DT liquid wicked into it may make it difficult to ensure a uniform thickness at shot time. These capabilities have not yet been demonstrated.

Box 4-2

Highlights of the High Average Power Laser Program

The goal of the HAPL Program (FY1999-2009) was to pursue integrated development of science and technology for IFE that would be, to the extent possible, simple, durable, and affordable without sacrificing performance. The program featured parallel efforts on KrF and diode-pumped, solid-state lasers (DPSSLs). A high priority was placed on acquiring experimental data for both laser systems and technology concepts. The Sombrero Power Plant study^a was used as a starting point.^b

The HAPL program was based on laser-driven, direct-drive targets because of their potential for higher drive efficiency, simpler target fabrication, lower estimated cost, and smaller inventory for material recycling. Both conventional hot-spot ignition and shock ignition concepts were investigated. Predictions indicated that the drivers were equivalent for the conventional ignition and that the shorter-wavelength target produced higher gains for shock ignition. At the

program goal of no more than 25 percent recirculating power, a combined driver target gain (η_G) of 10 was needed, corresponding to a minimum target gain of 140 for a 7 percent efficient laser system (e.g., KrF). The HAPL program made significant progress in repetitive laser technologies for both diode-pumped Nd:glass and electron-beam-pumped KrF, demonstrating multihour runs at pulse rates from 5 to 10 Hz.

Research and development supported by the HAPL program included (1) calculations of neutron damage to optical ports and optics trains; (2) the development and successful testing of a new dielectric grazing incidence multilayer mirror for the first optical element of the laser system; (3) the development and demonstration of a method to mass-produce foam shells for target capsules; and (4) the development and demonstration of a cryogenic fluidized bed to make DT layers economically (the estimated cost of production was less than \$0.17 each).

Target injection by both light-gas gun and magnetic slingshot was developed and tested. A method to improve capsule illumination accuracy detected the reflection (“glint”) from the moving capsule, of the light of a small laser to determine the target’s trajectory. Real-time adjustment of the laser mirrors enabled illumination that was within 28 μ of the ideal to be demonstrated.

^a Sviatoslavsky et al., 1992.

^b An overview of the HAPL results is in Sethian et al., 2010.

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The HAPL program demonstrated active aiming of the drive laser that reduced its equivalent positioning error to 28 μ . The “glint” technique, in which the target capsule was illuminated during its trajectory through the essentially evacuated reaction chamber by a separate laser, utilized optical sensor location of the target by reflected laser light to determine the appropriate aim point. The firing rate in HAPL-sponsored tests was 5 Hz.

Successful translation of the glint technique to LIFE-style IFE would require that the target trajectory be sufficiently predictable to allow enough time to adjust the directions of the laser beam cones. Should perturbations of the target trajectory increase to problematic levels as it neared its aim point (the center of the turbulent region), very rapid detection and aiming adjustments would be needed to meet the 100 μ -equivalent error requirement for the LIFE design. Orientation of an ID target is also important, unlike the spherical HAPL target capsule. The target insertion technique includes inducing a spin along the LEH axis to stabilize its orientation. Successful irradiation would require that a target’s angular momentum sufficiently overwhelm the effects of its hydrodynamic interaction with vorticity in the Xe fill of the reaction chamber that its orientation remains within acceptable bounds. Any second-order effects from also adjusting the aim of the laser beams are assumed here to be negligible. The difficulty of the other half of the glint technique—the illumination and detection of the target entering the reaction chamber—will be increased by the Xe fill. An assessment of this effect has not been presented to the panel.

Some unspecified portion of the gain margin calculated for the LIFE target has been allocated to compensating for nonoptimum insertion, but turbulence or other irregularities in the Xe gas through which the targets must pass could lead to sufficient inaccuracy not only to overwhelm that margin, but also to preclude capsule ignition. A key issue here is the repeatability of any phenomena that significantly perturb the target’s trajectory.

The LEH shields are themselves inside LEH windows that are needed in the LIFE concept to separate the reaction chamber Xe from the He inside the hohlraum. The LEH windows also represent an interface between the cold interior of the target and the prevailing

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2089 conditions of the reaction chamber. Some fraction of any Pb plasma or vapor from previous
 2090 capsules through which a target travels might be expected to condense on the LEH windows
 2091 during insertion and could affect the irradiation of the hohlraum interior.

2092 Lastly, the accelerations must not cause any portion of the supercooled DT to change
 2093 phase. Significant solidification would break the HDC ablator shell, and isolated solidification
 2094 would create density nonuniformities that would spoil the implosion, either directly or by
 2095 seeding hydrodynamic instabilities.

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2097 **Target Robustness**

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2099 The Merriam-Webster online dictionary⁴¹ has several meanings for “robust,” one of
 2100 which is pertinent to the current discussion: “capable of performing without failure under a wide
 2101 range of conditions.” Robustness will be used in what follows to mean the quality of being
 2102 robust according to this definition, with the regrettable caveat that the current state of the art
 2103 limits an assessment’s tie to reality to relatively indirect data. A result of this limitation is that
 2104 degrees of robustness actually indicate the assessed likelihood that a system can be made robust
 2105 by actions and processes that are anticipated, proposed, or otherwise foreseeable, and, more
 2106 fundamentally, the assessed likelihood that a system can be made to work at all.

2107 Based on evaluations of the associated issues, the panel assesses the robustness of the
 2108 physics design for the LIFE target concept to be low. The main factors leading to this assessment
 2109 are the following:

2110

- 2111 • Ignition of a fusion target operating in the physics regime of laser-driven ICF has
 2112 never been observed, but a robust design would have to reliably produce a large gain
 2113 under much less controlled conditions than are normal in laboratory experiments.
 2114 Moreover, the parameter space over which simulations predict adequate gain for the
 2115 LIFE target capsule is relatively small, and the optimization of several parameters, an
 2116 integral part of NIC, can be expected to further narrow the parameter space over
 2117 which sufficient gain might be obtained;
- 2118 • Significant departures from predicted operation have been observed on implosion
 2119 experiments pertinent to the LIFE target design. These disparities, which were
 2120 observed at both the NIF and the OMEGA lasers, relate directly to important aspects
 2121 of target operation (e.g., implosion velocity), and the targets in which they were
 2122 observed are the closest available analogues to the LIFE target. The discrepant data
 2123 are important to the calibration or validation of the simulations on which predictions
 2124 of the operation of the LIFE target are based, but tentative explanations of the
 2125 disparities are at this time unsupported;
- 2126 • To achieve the gain required for the LIFE plan to be viable, its target design
 2127 incorporates modifications that are likely to further reduce the predictability of the
 2128 target performance; and
- 2129 • The outer, dense thermonuclear fuel region of the LIFE target is planned to be
 2130 constructed of liquid DT wicked into low-density foam, but obtaining the gas
 2131 pressure believed to be required for successful operation would require cooling the
 2132 target capsule below the thermodynamic triple point for DT. The ability to create a

⁴¹ Available at www.merriam-webster.com.

2133 LIFE target as currently designed therefore requires the existence of a physical
 2134 phenomenon—the stabilization of a supercooled DT liquid in a low-density foam for
 2135 an extended period of time—that has never been observed and for which there is no
 2136 theoretical prediction.⁴²

2137
 2138 **CONCLUSION 4-4: The target design for a proposed indirect-drive inertial fusion energy**
 2139 **system (the laser inertial fusion energy or LIFE program developed by LLNL)**
 2140 **incorporates plausible solutions to many technical problems, but the panel assesses that the**
 2141 **robustness of the physics design for the LIFE target concept is low.**

- 2142
- 2143 • The proposed LIFE target presented to the panel has several modifications relative to
 2144 the target currently used in the NIC (for example, rugby hohlraums, shine shields, and
 2145 HDC ablaters), and the effects of these modifications may not be trivial. For this
 2146 reason, R&D and validation steps would still be needed.
 - 2147 • There is no evidence to indicate that the margin in the calculated target gain ensures
 2148 either its ignition or sufficient gain for the LIFE target. If ignition is assumed, the
 2149 gain margin briefed to the panel, which ranged from 25 percent to almost 60 percent
 2150 when based on a calculation that used hohlraum and fuel materials characteristic of
 2151 the NIC rather than the LIFE target, is unlikely to compensate for the phenomena
 2152 relegated to it—for example, the effects of mix—under any but the most extremely
 2153 favorable eventuality. In addition, the tight coupling of LIFE to what can be tested on
 2154 the NIF constrains the potential design space for laser-driven, indirect-drive IFE.

2155 **SOLID-STATE LASER-DRIVEN, DIRECT-DRIVE FUSION**

2156 **Current Status**

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 2160 The leader in direct drive inertial confinement fusion with solid-state lasers is the
 2161 Laboratory for Laser Energetics (LLE) at the University of Rochester, which operates the
 2162 OMEGA Laser Facility (OMEGA and OMEGA EP) for the National Nuclear Security
 2163 Administration (NNSA). LLE is conducting research into direct-drive ICF targets that utilize
 2164 either the hot-spot ignition concept used by the NIC capsule or one of the more recent two-step
 2165 ignition concepts (fast or shock ignition). The 60-beam OMEGA laser system, which delivers
 2166 >30 kJ of 3ω light on target with 1-2 percent irradiation nonuniformity, has been operating since
 2167 1995, is fully instrumented, and is capable of up to 1500 shots/year. The OMEGA EP laser
 2168 system, which adds four NIF-like beamlines (6.5 kJ at 3ω), was completed in April 2008 and can
 2169 propagate to either the OMEGA or OMEGA EP target chamber. Two EP beams can be operated
 2170 as a high-energy petawatt (2.6 kJ in the infrared in 10 ps) system.

2171 The current ICF program is aimed at exploring, understanding, and quantifying the
 2172 physics issues of direct-drive laser targets at OMEGA drive energies and extrapolating the target
 2173 performance to ignition and high-yield regimes. LLE has been routinely fielding cryogenic
 2174 capsules since 2001 and has seen a steady improvement in implosion experiments as they have
 2175 improved the quality of the ice layer and the centering of the target in the chamber. The flexible

⁴² There are studies that suggest it is possible to supercool hydrogen isotopes and other fluids (See, for example, Beaudoin et al., 1996. It remains unclear whether this effect can be achieved in the nanoporous hydrocarbon foam material, and if the corresponding vapor pressure is the desired value.

2176 pulse-shaping capability of OMEGA enables the generation of multiple-picket pulse shapes that
2177 can drive ignition-scaled cryogenic DT implosions to ignition-relevant implosion velocities ($3 \times$
2178 10^7 cm/s) on a low adiabat ($\alpha \sim 2$ - 3)⁴³. The energies and relative timings of the three pickets and
2179 main pulse are adjusted to optimize the coalescence of four shocks to create a central hot spot,
2180 the same implosion strategy used at NIF. Areal densities (ρr) up to 300 mg/cm² have been
2181 measured using a magnetic recoil spectrometer in cryogenic DT implosions on OMEGA drive at
2182 $\sim 8 \times 10^{14}$ W/cm² (Goncharov et al., 2010). The measured areal density in these experiments is
2183 larger than 88 percent of the predicted 1-dimensional (1-D) value. The measured area mass
2184 density, ion temperature, and neutron yield can be combined with computed 1-D neutron yield to
2185 estimate the overall ignition parameter (χ)⁴⁴ for these experiments. These OMEGA cryogenic
2186 implosions have achieved an appreciable fraction (~ 3 percent) of the overall ignition parameter.
2187 The low inferred adiabats of these targets suggest that hot electron production from LPI and
2188 deposition into the fuel are within acceptable limits.

2189 LLE has developed a 1 MJ symmetric, direct-drive NIF ignition design using a triple-
2190 picket pulse scaled to NIF laser parameters⁴⁵ that has a 1-D gain of ~ 50 . Since direct drive has
2191 higher implosion efficiency than indirect drive, it is calculated to produce higher target gains,
2192 which should lead to lower laser cost.

2193 No existing solid-state laser system in a direct-drive configuration presently has sufficient
2194 energy to demonstrate ignition. A multilaboratory workshop was held in 2001 whose purpose
2195 was not to preclude direct drive on NIF (Meyerhofer, 2001). It was also agreed that the change
2196 board process would be used to ensure that future modifications did not preclude direct drive on
2197 NIF. However, it is not clear that the final assembly procedure strictly adhered to this principle.

2198 Reconfiguring NIF to symmetric direct drive geometry represents the lowest target
2199 physics risk but the highest facility cost, and it would disrupt weapons physics experiments using
2200 hohlraums. As an alternative, LLE has identified a so-called “polar drive” (PD) geometry that
2201 allows direct-drive target performance to be studied at lower facility cost and minimal disruption
2202 of other experiments but at the price of higher target physics risk. Calculations predict that by
2203 repointing the beams from the existing laser ports, a uniform target drive can be achieved with
2204 PD irradiation, assuming that the irradiation at the equator is compensated by increased laser
2205 intensity. The risk is that the oblique irradiation at the equator occurs at lower densities, which
2206 reduces laser absorption and hydroefficiency and requires lateral heat flow to the equator from
2207 nonradial beams (Skupsky et al., 2004). The NIF triple-picket PD design with expected
2208 nonuniformities and multiple phase-modulation frequencies (multi-FM) beam smoothing
2209 achieves a calculated 2-dimensional (2-D) gain of 32.

2210 LLE has identified five changes on the NIF that would implement a PD capability for an
2211 ignition demonstration. OMEGA EP can be used to test many of the modifications, including
2212 multi-FM 1-D SSD beam smoothing,⁴⁶ and to validate laser performance.

⁴³ α is a measure of the degree to which the actual adiabat of the implosion exceeds the ideal Fermi-degenerate adiabat (for which $\alpha = 1$).

⁴⁴ The ignition parameter is the energy that would have had to be absorbed by the target to produce ignition based on the other parameters achieved in the implosion—symmetry, density, and so on, as calculated in simulations.

⁴⁵ This involves targets whose dimensions are scaled down from the ignition design due to the reduced energy on OMEGA relative to NIF.

⁴⁶ One-dimensional SSD with multiple phase-modulation frequencies (multi-FM) requires pre-conditioning the laser pulse with three high frequency-modulators to increase the bandwidth and is followed by a dispersion grating to increase the temporal skew. Multi-FM 1D SSD has been optimized to provide the required beam smoothing to enable PD ignition. See Marozas et al., 2010.

2213 Advanced two-step ignition concepts such as shock ignition (SI) or fast ignition (FI) provide
 2214 alternatives to conventional hot-spot ignition. If successful, these ignition options will open the
 2215 path to high-gain ICF ($G \sim 150$) for ~ 1 MJ laser drivers (Perkins et al., 2009; Betti et al., 2006).

2216 Fast ignition requires a combination of long-pulse (implosion) and short-pulse (FI) lasers.
 2217 Aspects of fast ignition both by electrons⁴⁷ and protons⁴⁸ were briefed to the panel. Integrated FI
 2218 experiments have begun on OMEGA as part of the program of the DOE Office of Fusion Energy
 2219 Sciences, which is studying the fast-electron coupling into a compressed core. The inferred laser-
 2220 to-target heat coupling of ~ 3.5 percent needs to be increased significantly for FI to be a viable
 2221 concept. Integrated simulations of electron-driven fast ignition experiments are challenging and
 2222 do not presently suggest ways of improving the target coupling. In principle, FI can also be
 2223 achieved with protons accelerated by ultrashort-pulse lasers, which has the advantage of ballistic
 2224 ion transport and sharper energy deposition. However, proton FI is hindered by lower laser
 2225 conversion efficiency (~ 10 percent experimentally), a high intensity requirement ($\sim 10^{20}$ W/cm²),
 2226 and a high proton-dose requirement ($\sim 10^{16}$ protons) that complicates target fabrication. Further, a
 2227 more complicated capsule design is required if a reentrant cone is used to protect the proton-
 2228 generation foil. Although there is international interest in FI (e.g., the Fast Ignition Realization
 2229 Experiment (FIRE) project at ILE/Osaka and HiPER in the U.K.), funding is presently
 2230 insufficient for FI to challenge the mainline programs on NIF or the Laser Megajoule Facility
 2231 (LMJ), which is under construction in France. Furthermore, the recently proposed concept of SI
 2232 appears to be an easier and more attractive alternative to standard hot-spot ignition. SI utilizes a
 2233 standard long-pulse laser beam with a pulse shape that provides a high-intensity spike at the end
 2234 of the main drive pulse. The SI concept has been tested using CH shells on OMEGA. Higher
 2235 areal densities (30 percent) and significantly higher neutron yields ($\sim 4x$) were achieved with SI
 2236 pulse shapes (Theobald et al., 2008).

2237 Continued fundamental research into FI theory and experiments, the acceleration of
 2238 electrons and ions by ultrashort-pulse lasers, and related high-intensity laser science is justified.
 2239 However, issues related to low laser-target energy coupling, a complicated target design, and the
 2240 existence of more promising concepts (such as SI), led the panel to the next conclusion on the
 2241 relative priority of FI for fusion energy.

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 2243 **CONCLUSION 4-5: At this time, fast ignition appears to be a less promising approach for**
 2244 **IFE than other ignition concepts.**

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Recent and Upcoming Work

2249 The in-flight shell adiabat has been tuned using shock-velocity measurements using a
 2250 variant of the NIF “key-hole target” (Boehly et al., 2011). Cross-beam energy transfer (CBET)
 2251 has been identified as an issue that may be reducing laser energy absorption on OMEGA by 20
 2252 percent. Near-term experiments are planned to study mitigation strategies using modified phase-
 2253 plate designs. Initial shock ignition designs for the NIF have 1-D gains of 70 at 680 kJ, with
 2254 about half of that total energy in the shock generation pulse. PD diagnostic commissioning
 2255 targets using existing ID phase plates are being imploded on the NIF (Cok et al., 2008).

⁴⁷ David Meyerhofer, LLE, “Fast and Shock Ignition Research,” presentation to the panel on July 6, 2011.

⁴⁸ Juan Fernandez, LANL, “Inertial Confinement Fusion (ICF) Targets at Los Alamos National Laboratory,” presentation to the panel on May 10, 2011.

2256 LLE continues to demonstrate hydroequivalent scaling experiments on OMEGA to
 2257 validate design codes that are then used for PD ignition calculations for NIF.
 2258 Upcoming experiments using targets with improved quality and reduced offset from the target
 2259 chamber center are predicted to increase the χ from 3 percent of ignition to 5-6 percent,
 2260 achieving the maximum credible performance for a 30-kJ driver.
 2261 LLE is developing a project execution plan (PEP) to demonstrate PD ignition on the NIF
 2262 in 2017.

2263 **Evaluation and Discussion of Remaining R&D Challenges**

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 2266 Direct-drive, capsule-implosion data exist only at the 30 kJ level. The predicted
 2267 hydroequivalent scaling requires validation at the MJ energy level, including issues of LPI,
 2268 shock ignition at MJ energies, and symmetry. The modifications of NIF for PD need to be
 2269 developed and tested on OMEGA and deployed on NIF. There are target physics risks for polar
 2270 drive that need to be studied. Further, there are target fabrication, injection, and survival issues
 2271 that are specific to the direct drive approach. Specific issues are discussed individually below.
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2273 **LPI**

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 2275 The larger energies for ignition targets are achieved through longer laser pulses, which
 2276 result in long-scale-length plasmas that are more susceptible to LPI. There is a need to study and
 2277 demonstrate acceptable laser energy deposition and hot electron production for ignition-scale
 2278 plasmas. Relevant experiments can be done on OMEGA EP, which has NIF long-pulse beam
 2279 lines. In particular, planar two-plasmon decay (TPD) experiments can quantify the hot electron
 2280 production by collecting all electrons.

2281 There are critical uncertainties in extrapolating TPD physics in planar geometry to the
 2282 oblique irradiation geometry of the equatorial beams for NIF PD. Integrated TPD experiments on
 2283 OMEGA will be very important in quantifying the production and deposition of hot electron
 2284 energy.

2285 The plasma physics community requires a better understanding of cross-beam energy
 2286 transfer, including better theory and modeling, additional measurements, and tests of potential
 2287 mitigation techniques.

2288 The ability to model underdense plasma conditions is important for understanding LPI,
 2289 since most LPI depend exponentially on electron density and temperature. Continued
 2290 development of these models—including the effects of nonlocal transport—is important,
 2291 especially for PD beam geometries.
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2293 **Shock Ignition**

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 2295 Fully integrated 2-D point designs for NIF PD shock ignition targets are required in order
 2296 to plan for experimental campaigns on NIF. Experiments need to continue on OMEGA to
 2297 identify whether there are any LPI issues that are unique to the SI approach, especially in PD
 2298 geometries. Experiments need to be done on OMEGA and later on NIF to determine whether the
 2299 hot-electron production by the high-intensity spike is acceptable for high-gain target
 2300 performance. Calculations and experiments need to be performed to study the implementation of

2301 shock ignition pulses, including the trade-offs among laser beam parameters, illumination
2302 symmetry, and SI performance.

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2304 **Symmetry**

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2306 It remains to be seen whether sufficiently smooth laser beams can be created on the NIF
2307 to allow direct drive experiments, particularly in the PD geometry. Pointing errors and nonradial
2308 deposition geometries could lead to low-mode symmetry errors. Insufficient beam smoothing
2309 could lead to high-mode asymmetries. Symmetry issues related to providing both normal and
2310 high-intensity beams to illuminate SI targets need to be investigated, including calculations and
2311 experiments in PD geometry.

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2313 **Reconfiguring NIF for Polar Drive**

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2315 The following steps need to be taken to enable polar drive experiments on NIF:

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- 2317 • Demonstrate new multi-FM 1-D SSD beam smoothing technique and validate on
- 2318 OMEGA EP.
- 2319 • Design and demonstrate tailored phase plates to increase equatorial beam coupling.
- 2320 • Design and demonstrate polarization smoothing for OMEGA EP to reduce focal-spot
- 2321 irradiance modulation. Design and demonstrate distributed polarization rotators
- 2322 (DPRs) that are sufficient to achieve polar-drive ignition on NIF.
- 2323 • Demonstrate integrated NIF PD beam smoothing on OMEGA EP.
- 2324 • Complete development of a NIF fill-tube target that meets polar-drive ice layer
- 2325 specifications.
- 2326 • Complete development of concepts for a PD ignition target insertion cryostat.

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2328 **Polar Drive Physics**

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2330 Understanding of the following areas of polar drive target physics need to be improved:

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- 2332 • Deposition in low-density plasma by oblique beams at equator, including 3-
- 2333 dimensional (3-D) laser ray trace algorithms that are compatible with PD geometry.
- 2334 • Ability of laser to deliver increased intensity to equatorial beams.
- 2335 • Nonlocal transport and heat conduction for nonradial beams; this may require
- 2336 extensions to existing theory and algorithms.
- 2337 • Possible LPI issues unique to PD illumination geometry; e.g., CBET between
- 2338 overlapping beams.

2339

2340 **CONCLUSION 4-6: The prospects for ignition using laser direct drive have improved**
2341 **enough that it is now a plausible alternative to laser indirect drive for achieving ignition**
2342 **and for generating energy.**

2343

- 2344 • The main concern with laser direct drive has been the difficulty of achieving the
- 2345 symmetry required to drive such targets. Advances in beam-smoothing and pulse-
- 2346 shaping appear to have lessened the risks of asymmetries. This assessment is

2347 supported by data from capsule implosions (performed at the University of
 2348 Rochester's OMEGA laser), but it is limited by the relatively low drive energy of the
 2349 implosion experiments that have thus far been possible. Because of this, the panel's
 2350 assessment of laser-driven, direct drive targets is not qualitatively equivalent to that
 2351 of laser-driven, indirect-drive targets.

- 2352 • Further evaluation of the potential of laser direct-drive targets for IFE will require
 2353 experiments at drive energies much closer to the ignition scale.
- 2354 • Capsule implosions on OMEGA have established an initial scaling point that
 2355 indicates the potential of direct-drive laser targets for ignition and high yield.
- 2356 • Polar direct-drive targets will require testing on the NIF.
- 2357 • Demonstration of polar-drive ignition on the NIF will be an important step toward an
 2358 IFE program.
- 2359 • If a program existed to reconfigure NIF for polar drive, direct-drive experiments that
 2360 address the ignition scale could be performed as early as 2017.

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Potential for Use in an IFE System

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2365 If ignition and high yield can be demonstrated for DD targets, the higher target gain
 2366 translates into greater system efficiency and lower laser energy (size). The even higher predicted
 2367 gains of shock ignition targets make this DD concept very attractive. Shock ignition is not an
 2368 option for ID targets due to the inherent integrating nature of the hohlraum, which limits the
 2369 ability to spike the temperature drive.

2370 Demonstrating PD ignition on the NIF is an important step toward an IFE program. This
 2371 should include experiments to explore the performance of shock ignition targets on NIF.

2372 To date, the LLE ICF program has been focused on the development of laser beam
 2373 smoothing technologies and single-shot ICF target physics experiments, which is the appropriate
 2374 scope of the NNSA program. With the exception of some work in developing mass-production
 2375 techniques for fabricating cryogenic DD targets and studying their survival in IFE-relevant
 2376 thermal environments, LLE has not conducted research into either repetitive solid-state laser
 2377 technologies or the host of issues associated with an IFE power plant. Through the HAPL
 2378 program, LLNL has been the lead laboratory in developing repetitive solid-state lasers (DPSSL
 2379 technology). Similarly, through the HAPL program, the Naval Research Laboratory (NRL) has
 2380 supported the study of many of the technology and material issues related to the operation of a
 2381 DD power plant. This suggests that there are opportunities for teaming among LLE, LLNL, and
 2382 NRL if an IFE program is established to explore the potential of a DD fusion power plant with
 2383 solid-state lasers. Further, LLE has much to contribute in target physics and target fabrication if
 2384 KrF lasers prove more attractive as the laser driver in a DD power plant.

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Additional Considerations

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Target Injection

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A key issue here is the repeatability of any phenomena that significantly perturb the
 target's trajectory.

2393 Survival of Cryogenic Target

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2395 LLE has been studying the survival of cryogenic DD targets via complete Monte Carlo
 2396 and computational fluid dynamics modeling of heat load to the target and its effect on the ice
 2397 during injection into the chamber. These calculations will be supplemented by experiments in a
 2398 surrogate IFE chamber. This issue was also addressed in the HAPL program, but more study is
 2399 needed.

2400

2401 Reactor Chamber Issues

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2403 Most direct-drive IFE schemes are predicated on a dry-wall concept and an evacuated
 2404 chamber. There are a host of structural and material issues that need to be addressed. The HAPL
 2405 program supported initial research in most of these areas, but much more work will be required
 2406 before a power plant design can be completed. The HAPL final optic train was designed to meet
 2407 the requirements for illumination uniformity, adequate tritium breeding, the threshold for
 2408 damage to the grazing incidence metal mirror, and neutron damage to the conventional DD
 2409 target. This design was applicable to both DPPSLs at 351 nm and KrF at 248 nm (Sethian et al.,
 2410 2010).

2411

2412 **CONCLUSION 4-7: In general, the science and engineering of manufacturing fusion**
 2413 **targets for laser-based ICF are well advanced and meet the needs of those experiments,**
 2414 **although additional technologies may be needed for IFE.** Extrapolating this status to predict
 2415 the success of manufacturing IFE targets is reasonable if the target is only slightly larger than the
 2416 ICF target and the process is scalable. However, subtle additions to the design of the ICF target
 2417 to improve its performance (greater yield) and survivability in an IFE power plant may
 2418 significantly affect the manufacturing paradigm.

2419

2420 **CONCLUSION 4-8: There are important differences between the direct-drive and**
 2421 **indirect-drive based targets. The direct-drive target is simpler to build than is the indirect-**
 2422 **drive target, and it is more vulnerable to the environment when it is injected into the target**
 2423 **chamber.** Understanding these nuances and demonstrating a viable manufacturing process
 2424 would likely be an important early priority for an IFE program because the quality and
 2425 variability in the target's specifications can strongly affect the target's gain.

2426

2427 **CONCLUSION 4-9: One major area where the IFE laser-driven target differs from the**
 2428 **ICF target is the method of delivering the target to the target chamber at a high frequency.**
 2429 The high-velocity projectile techniques proposed for laser-based fusion show promise, but there
 2430 has been little quantification of the degree to which the target will be compromised during the
 2431 process and what effect any degradation may have on the target's gain. Also, changes that need
 2432 to be made to the ICF target to improve its survivability in the IFE target chamber environment
 2433 have been identified, but the consequence of these changes for the manufacturing process is not
 2434 known. These are issues that need to be thoroughly addressed early in any future IFE program.

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KRYPTON FLUORIDE LASER-DRIVEN, DIRECT-DRIVE FUSION

The leader in DD inertial confinement fusion with krypton fluoride (KrF) lasers is the Naval Research Laboratory (NRL) in Washington, D.C., which operates the Nike and Electra lasers. Nike is the world's largest KrF laser. Its amplifier with 60-cm aperture delivers a pulse of between 3 and 5 kJ at 248 nm to planar geometry targets using a smoothing technology called "induced spatial incoherence" (ISI). Nike has demonstrated "focal zooming," which allows the laser to more efficiently deliver late-time energy to the imploding spherical ICF pellet.

Electra is a repetitive KrF laser that was developed as part of the HAPL program to study the technology issues of repetition rate, durability, efficiency, and cost for inertial fusion energy. The HAPL program is discussed in Box 4-2. NRL has also developed the FAST (Gardner et al., 1998, and Zalesak et al., 2005) radiation-hydrocode, which has several unique features that make it complementary to the ICF codes used at other laboratories.

The current ICF program on the Nike laser is focused on studying the hydrodynamic performance of planar targets accelerated by very smooth laser beams at 248 nm. LPI theories predict higher intensity thresholds for shorter wavelength lasers, proportional to the square of the wavelength. Further, shorter wavelengths enable higher absorption efficiency, larger drive pressure, and higher hydrodynamic efficiency. Experiments to quantify the growth of Richtmyer-Meshkov and Rayleigh-Taylor instabilities in planar cryogenic (deuterium wicked into foam) targets with thicknesses close to that of a high-gain target have been published (Pawley et al., 1999) and found to be in good agreement with theoretical predictions by the FAST3D code. The use of a thin high-Z layer to mitigate the imprinting of nonuniformities in the low-intensity laser foot was proposed and validated on Nike (Obenschain et al., 2002).

Further collaborative validation experiments on OMEGA demonstrated "significant and absolute (2X) improvements in neutron yield when the shells are coated with a very thin layer (~200–400 angstrom) of high-Z material such as palladium" (Mostovych et al., 2008). Thus, this imprint mitigation technique has been shown to work in both planar and spherical geometries at 248 and 351 nm. The utility of the high uniformity and higher ablation pressure generated by the Nike KrF laser was recently demonstrated in experiments on hypervelocity acceleration of planar targets in collaboration with researchers at the Institute of Laser Engineering at Osaka University in Japan. Whereas the Gekko XII/HIPER glass laser (351 nm) achieved a 700 km/s velocity, the KrF laser was able to achieve a 1,000 km/sec foil velocity (Karasik et al., 2010). Extrapolating this performance to spherical DD implosions, ISI and zooming with a KrF laser offer the potential to use targets having lower aspect ratios and to reduce hydroinstability growth, thereby achieving higher target gain for less laser energy.

In 2008, Nike was upgraded to enable high-intensity LPI target experiments. The $2\omega_{pe}$ instability at quarter-critical density is of greatest concern in DD targets, where measurement of $\omega_0/2$, $3\omega_0/2$, and hard X-ray (>20 keV) emissions indicate the onset of the instability. The quarter-critical critical instability thresholds observed in Nike experiments with ISI-smoothed beams are in approximate agreement with planar beam $2\omega_{pe}$ theory, which does not account for the effects of beam smoothing, beam overlap, or saturated levels. This agreement includes an attempt to study the scaling with plasma scale length by varying the laser pulse length. OMEGA experiments with beams smoothed by SSD show similar agreement, and the predicted wavelength scaling appears to have been obtained. The OMEGA experiments have been modeled using the FAST and LILAC codes, both of which are in agreement with respect to the onset of LPI (Seka et al., 2009). However, DD ignition targets will likely need to operate above

2485 this theoretical threshold, and further research to understand, model, and measure LPI is
 2486 required. This includes utilizing the NIF-equivalent OMEGA EP beam parameters to study LPI
 2487 at plasma scale lengths that are relevant to ignition high-yield DD IFE targets.

2488 A series of DD IFE target designs have been studied with the goal of maximizing target
 2489 gain while minimizing laser energy. A conventional DD design provided IFE-relevant 1-D gains
 2490 ($G \sim 100$) at laser energies of ~ 1.3 MJ (Bodner et al., 2002). Later designs gave 1-D gains of
 2491 order 50 with 500 kJ of KrF laser light by going to higher implosion velocities and using early
 2492 time spikes in the pulse shape to tailor the implosion adiabat and diminish Rayleigh-Taylor
 2493 instability growth (Colombant et al., 2007).

2494 The shock ignition concept proposed by Betti (Betti et al., 2007), and discussed in more
 2495 detail in the preceding section, is now the baseline for KrF designs because of the higher
 2496 predicted gains. An initial step in validating these designs was obtaining the agreement of FAST
 2497 simulations of neutron yields with LLE simulations and experiments (Theobald et al., 2008). At
 2498 IFE energies, FAST simulations of ISI-smoothed KrF beams using focal zooming give shock-
 2499 ignition 1-D gains that are roughly twice as high as the best conventional designs (Schmitt et al.,
 2500 2009). High-resolution, 2-D FAST simulations (for Legendre modes $l = 1-256$), which include
 2501 the effects of inner and outer surface finishes and laser imprint, predict that these targets are
 2502 robust to such perturbations.

2503 The KrF research program would benefit from further 3-D implosion studies, improved
 2504 LPI simulations, and experimental validation from LPI and implosion experiments on both
 2505 OMEGA and NIF in PD configuration. However, in PD geometry, the oblique irradiation near
 2506 the equator occurs at lower densities, which reduces absorption and hydroefficiency and
 2507 introduces nonradial beam illumination geometries and lateral heat flow. These are the remaining
 2508 R&D challenges.

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2511

Recent and Upcoming Work

2512

2513 Having adequate numerical models for nonlocal thermal and hot-electron transport has
 2514 been a challenge for several decades. Of special concern for DD, electron thermal transport in a
 2515 laser-produced plasma cannot be described with a local approximation in many regions because
 2516 the electron mean free path is longer than the temperature gradient scale length. NRL researchers
 2517 have found that a Krook model provides reasonable descriptions of both preheat and flux
 2518 limitation and have developed a computationally tractable algorithm; they are now verifying the
 2519 accuracy of the model. This improved model will soon be available to apply to the analysis and
 2520 design of ongoing experiments, as well as to the design of PD experiments on NIF. These models
 2521 are also relevant to the uncertainties in NIF hohlraum modeling.⁴⁹

2522 NRL has recently begun to simulate polar, DD implosions on NIF using the FAST code.
 2523 This will complement ongoing work by LLE in defining DD experiments for a polar-drive
 2524 platform on NIF. The growing collaboration will allow development of conventional and shock
 2525 ignition designs for NIF and will enable use of the new Krook model to study the effect of
 2526 nonlocal transport in the PD geometry.

2527

2528

Evaluation and Discussion of Remaining R&D Challenges

⁴⁹ M. Rosen, LLNL, “Understanding of LPI and its impact on indirect drive,” presentation to the panel on September 21, 2011.

2529
2530 NRL presented a path forward to IFE DD target physics that included implosion
2531 experiments on OMEGA, LPI experiments on both Nike and OMEGA EP, and polar DD
2532 experiments on NIF. The theory and simulation efforts included the development of better
2533 physics models for the FAST code, improved two- and three-dimensional hydroimplosion
2534 simulations, and improved ability to perform LPI simulations. NRL also proposed the
2535 development of one KrF IFE beam line that was capable of delivering ~20 kJ on target to study
2536 target interaction and LPI physics at IFE-relevant intensity and plasma scale lengths. The goal of
2537 this program, to be carried out in collaboration with LLE, would be to validate the fundamental
2538 physics of DD, to determine whether sufficient gains were feasible for IFE, and to validate the
2539 physics models for comparing DD target performance at 248 nm and at 351 nm.

2540 The fundamental issues for DD capsules are the same at these two wavelengths, and the
2541 plans discussed in the solid-state laser DD section are all relevant and necessary. The importance
2542 of extending the OMEGA target performance database to NIF energies cannot be
2543 overemphasized. Specific issues relevant to the NRL program are discussed individually below.
2544

2545 **Direct-Drive Theory and Physics Models**

2546

2547 There is a continued need to develop improved physics models for DD in FAST,
2548 especially for potential megajoule-class experiments on NIF, but in a nonradial, PD geometry.
2549 This includes continued development of nonlocal thermal and hot electron transport models,
2550 improved non-local thermodynamic equilibrium (non-LTE) radiation modeling (particularly for
2551 thin, high-Z layers) and improved laser ray tracking for NIF PD geometries. There is also a need
2552 for improved LPI modeling, perhaps by teaming with other groups that have developed this
2553 capability and applying it at KrF wavelengths.
2554

2555 **Laser-Plasma Interactions**

2556

2557 As part of an increased effort toward understanding LPI, data on thresholds at KrF
2558 wavelengths will be useful. If a 20-kJ KrF laser was developed, it would provide the capability
2559 to study LPI at 248 nm in relevant scale-length plasmas and compare the results with OMEGA
2560 EP data. LLE is currently studying the role of CBET in DD experiments on OMEGA. KrF IFE
2561 designs may need to account for this physics, including the trade-off between CBET and
2562 illumination symmetry.
2563

2564 **Polar Drive Physics, Symmetry, and Shock Ignition**

2565

2566 All of the issues listed under the solid-state DD section are relevant to the KrF DD
2567 program. Research into the physics issues of PD geometries, illumination symmetry in all DD
2568 geometries, and exploration of the potential of shock ignition as a high-gain target concept might
2569 best be pursued as a collaborative ICF/IFE program with both OMEGA and NIF.
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2571 **Capsule Fabrication, Injection, and Survival**

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2573 These issues are similar to those already described for solid-state laser-driven targets.
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Potential for Use in an IFE System

As noted in the preceding section, if ignition and high yield can be demonstrated for DD targets, the higher target gain translates into higher system efficiency and lower laser energy (size) at either 351 nm or 248 nm. If high-gain shock ignition proves feasible, the theoretical increase in gain for KrF with focal zooming as compared to frequency-tripled glass (with or without zooming) appears significant enough to merit serious consideration in IFE power plant economics. Further, from a driver perspective, the simplicity and effectiveness of ISI beam smoothing and focal zooming, the self-repairing nature of a gaseous gain medium, and the promising performance of the Electra laser system make KrF an IFE laser technology worth exploring. The final decision between 351 and 248 nm should be based on a total system performance analysis, including laser efficiency, durability, power plant integration issues, and overall target gain and performance. At this point, it would seem that an overall collaboration in direct drive target physics and a competition between driver technologies at the beamline level would be a prudent technology maturation path.

CONCLUSION 4-10: Experiments on Nike in recent years give technical credence to using the deep-ultraviolet KrF wavelength to improve hydrodynamic coupling and increase LPI thresholds for direct-drive targets.

- Implosion experiments at 351 nm on OMEGA have made DD an attractive option for IFE. Planar experiments at 248 nm on Nike using ISI-smoothed beams have demonstrated the expected favorable scaling, with shorter wavelengths for laser absorption, increased drive pressure, and higher hydrodynamic efficiency, as well as higher LPI thresholds.
- The DD community would benefit from conventional and shock ignition experiments in PD geometry on OMEGA and NIF, which might best be pursued as a national collaborative effort.
- Extending the Nike laser to 20 kJ would provide a valuable capability to study LPI and hydrodynamics at 248 nm in IFE-relevant scale-length plasmas and compare the results with OMEGA EP and NIF data.
- An overall collaboration in DD target physics and a competition between driver technologies at the beamline level would appear to be a prudent technology maturation path. The ultimate choice of laser wavelength and associated technology for DD IFE will be based on a total system analysis.

CONCLUSION 4-11: The lack of understanding surrounding LPI remains a substantial but as yet unquantified consideration in ICF and IFE target design.

RECOMMENDATION 4-1: DOE should foster collaboration among different research groups on the modeling and simulation of laser-plasma interactions.

HEAVY-ION-DRIVEN TARGETS

Current Status

2621 The U.S. Heavy-Ion Fusion Science Virtual National Laboratory is a collaboration
 2622 between LBNL, LLNL, and the Princeton Plasma Physics Laboratory (PPPL). The research is
 2623 headquartered at LBNL. The Fusion Energy Sciences (FES) program within the Department of
 2624 Energy manages the heavy-ion fusion program. Historically, the mainline heavy-ion fusion
 2625 (HIF) target design was developed to leverage the NIF experiments to demonstrate hot-spot
 2626 ignition of an indirect drive target. Correspondingly, the most mature HIF target designs are for
 2627 hohlraums with two-sided illumination (like NIF) that indirectly drive a scale-up of the NIF
 2628 capsule using repetitive accelerator technologies to provide the driver energy. ID hohlraums
 2629 with NIF-like hot-spot ignition implosion physics are a well-documented approach (Callahan et
 2630 al., 2002). For example, the 2002, two-dimensional Lasnex (Zimmerman et al., 1978) design
 2631 called for a 7 MJ heavy-ion driver delivering 3 and 4 GeV Bi⁺¹ ions to the hohlraum, giving a
 2632 fusion gain of 68.

2633 ID, and DD with hot-spot ignition or shock ignition using heavy-ion beams, are based on
 2634 laser concepts but exploit the classical physics of ion-plasma energy deposition.⁵⁰ The briefing
 2635 the panel received on heavy ion target design at the July 2011 meeting⁵¹ focused on the much
 2636 newer X-target. The X-target is a HIF-motivated design that uses single-sided illumination by
 2637 three sequential beam pulses and has features that offer new opportunities in accelerator driver
 2638 technology, chamber technology, and driver-chamber interface.

2639 Two preliminary target designs were presented to the panel at the Rochester meeting: 1) a
 2640 1-D Lasnex design of a DD target requiring 3 MJ of 3 GeV Hg⁺¹ ions, giving a gain of ~150, and
 2641 2) a single-sided direct-drive X-target also utilizing 3 MJ of ions with a calculated 2-D gain of
 2642 between 50 and 400 (see Figure 2-6). There are plans to extend the DD target design to 2-D
 2643 design to incorporate a PD illumination geometry as well as a tamper and shock ignition assist.

2644 Uranium beams of 80 GeV are already focused to <300 μm (full-width at half maximum)
 2645 at GSI in Germany (transverse emittance sufficiently low), but beam current and space charge
 2646 effects are small, and the bunch pulse durations are too long for fast ignition (>100 ns).
 2647 Experiments at LBNL (NTX and NDCX-I) have shown that intense beam space charge can be
 2648 neutralized with pre-formed target chamber plasma >>> beam density. However, plasma
 2649 neutralization cannot prevent the spread of the focal spot size due to chromatic aberrations
 2650 (random momentum spread in the beam).

2651 The sole LBNL target designer is continuing to evolve the X-target calculations in 2-D
 2652 using the LLNL HYDRA code.

2653

2654 Evaluation and Discussion of Remaining R&D Challenges

2655

2656 The limitation of present accelerators in energy and focal intensity means that there are
 2657 only a few data on ion-stopping powers in warm dense matter and no ICF target data. The PD
 2658 and X-target performance estimates are purely based on rad-hydro code simulations that need to
 2659 be greatly increased in sophistication and resolution to deal with all of the issues in a
 2660 computational sense. The entry-level price of a heavy-ion target physics facility is sufficiently
 2661 high that it is unlikely to be constructed by the DOE/NNSA program in the near or medium term.

2662

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⁵⁰ L.J. Perkins, LLNL, "Targets for Heavy Ion Fusion Energy," a presentation to the panel on February 16, 2011.

⁵¹ B.G. Logan, LBNL, "Heavy-Ion Target Design," presentation to the panel on July 7, 2011.

2665 Integrated 3-D target design

2666

2667 The 3-D nature of the HIF targets and highly sheared flows will require increasingly
2668 sophisticated simulations at very high resolutions (massively parallel).

2669

2670 Mix

2671

2672 The sheared flows in the X-target with high-Z slide surfaces make mix with the DT fuel a
2673 serious concern.

2674

2675 Acceleration Compression Physics

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2677 It will be very challenging to reach the 200 ps /200 μm radius goals of the accelerator
2678 physics program. Ultimately, the limits of focusing and compression are determined by
2679 Liouville's theorem. The NDCX-II experiments will explore more intense beam compression
2680 and focusing physics related to subnanosecond heavy-ion shock ignition and fast ignition.

2681

2682 Neutralized Ballistic Focusing

2683

2684 The conceptual X-target designs assumed neutralized ballistic focusing of heavy ions
2685 through a background chamber plasma as simulated by the IBEAM systems code (Meier et al.,
2686 2002). Some panel members question the maturity of the models for dynamic charge state; the
2687 degree of neutralization in the reactor chamber environment; and the potential impact of beam
2688 space charge on the final focus. This is a transport issue that is unique to heavy-ion fusion and
2689 will require further research through detailed simulations and validation by experimental data
2690 (Sharp et al., 2004).

2691

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2693 Potential for Use in an IFE System

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2695 All three heavy-ion target physics options are intended to use multiple-beam linac drivers
2696 with thick liquid-protected chambers to mitigate material neutron damage risks. The liquid-
2697 protected chamber technology is synergistic with some aspects of the pulsed-power approach to
2698 IFE.

2699

2700 In principle, the injection of targets into the reactor chamber for heavy ions has the same
2701 features as laser fusion. Light-gas-gun or magnetic-slingshot systems developed for laser fusion
2702 should be applicable. If the heavy-ion chamber uses a liquid lithium protection for the first wall,
2703 there may be some differences in injection system implementation and the specifics of cryogenic
2704 layer survivability in the reactor environment, which would be accounted for in a detailed system
2705 study.

2706

2707 All of the DD heavy-ion fusion target concepts are at a very early stage. Similarly, the
2708 proposed novel accelerator techniques for compressing heavy-ion beams to 200 ps with focusing
2709 to 200 μm radius are challenging and at an early stage of research. While heavy ions may
2709 represent a promising long-term option for efficient, reliable, repetitive fusion power plants, they
2710 probably represent a second- or third-generation capability.

2710

2711 **CONCLUSION 4-12: The U.S. heavy-ion-driven fusion program is considering direct-drive**
 2712 **and indirect-drive target concepts. There is also significant current work on advanced**
 2713 **target designs.⁵² This work is at a very early stage, but if it is successful, it may provide**
 2714 **very high gain.**

- 2715 • The work in the HIF program involves solid and promising science.
- 2716 • Work on heavy-ion drivers is complementary to the laser approaches to IFE and
- 2717 offers a long-term driver option for beam-driven targets.
- 2718 • The HIF program relating to advanced target designs is in a very early stage and is
- 2719 unlikely to be ready for technical assessment in the near term.
- 2720 • The development of driver technology will take several years, and the cost to build a
- 2721 significant accelerator driver facility for any target is likely to be very high.

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Z-PINCH TARGETS

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Description of Current U.S. Efforts

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2728 The main research in Z-pinch-driven ICF is performed at Sandia National Laboratories in
 2729 Albuquerque, New Mexico. After the conversion of the PBFA-II accelerator to “Z” in 1997 to
 2730 increase the radiated power from its wire-array Z-pinches, Sandia transitioned its ICF research
 2731 from light-ion beam drivers to Z-pinches. The initial ICF concepts utilized thermal radiation
 2732 from Z-pinches to indirectly drive ICF capsules. For example, the double-ended hohlraum
 2733 concept drew heavily from ID ICF design experience at NIF. Initial experiments on this concept
 2734 demonstrated control of radiation symmetry via backlit capsule implosions; however,
 2735 calculations showed that significant fusion experiments required much higher currents than
 2736 achievable on Z (60 MA for high yield versus 20 MA Z capability). After completion of the Z
 2737 Refurbishment Project in October 2007 (26 MA peak current), NNSA issued guidance that the
 2738 primary mission of Z should be to support the Science Campaigns within its Stockpile
 2739 Stewardship program, especially in the areas of dynamic materials and nuclear weapons effects.
 2740 Presently, the limited portion of the Z experimental program that is devoted to ICF research is
 2741 focused on concepts utilizing the DD of high magnetic field pressure to implode DT fuel to
 2742 fusion conditions, citing an estimated 25-fold increase in theoretical efficiency for direct
 2743 magnetic drive versus indirect X-ray drive. The Magnetized Liner Inertial Fusion (MagLIF)
 2744 concept (see Figure 2-7) has been theoretically developed, and initial experiments to study the
 2745 stability of the shell during magnetic implosion have been completed. Future experiments will
 2746 add laser preheat to the magnetic implosions, with the eventual goal of $G = 1$ laboratory
 2747 breakeven (DT fusion yield equals energy delivered to fuel). Quantitatively, this translates to
 2748 ~100 kJ DT yields, although D₂ experiments will initially be performed for simplicity. High-
 2749 yield (GJ-class), high-gain (>500) target designs are under development. Much of the relevant
 2750 physics can be tested on Z.

2751

R&D Challenges and Requirements

2752

2753 Some Z-pinch IFE system concepts were developed several years ago during a brief
 2754 period when limited funding for IFE technology was provided within the NNSA ICF program.

⁵² Advanced designs include DD, conical X-target configurations (see Chapter 2).

2755 The concept of a recyclable transmission line (RTL) was explored as part of this technology
 2756 project, although it was intended for use with the ID target designs that were being studied at that
 2757 time. Extrapolated calculations of Z-pinch target designs typically require around 60 MA of
 2758 current to be delivered from the pulsed-power driver to the implosion system to achieve high
 2759 fusion yields. In contrast to laser and heavy-ion targets, which receive their energy from beams
 2760 that are transported either in a vacuum or through small amounts of gas within the reactor
 2761 chamber, the RTL directly connects the driver to the Z-pinch fusion target. This energy delivery
 2762 strategy leads to a unique set of challenges and requirements for achieving the Z-pinch fusion
 2763 system performance. The economics of this system design favor a low repetition rate and a high
 2764 fusion target yield.

2765 Technical and program managers at Sandia indicated to the panel that they perceive that
 2766 ICF target research is not considered a high priority given the extensive funding necessary for
 2767 the NIC and DOE's current prioritization of high-energy-density-physics experiments on Z (e.g.,
 2768 the plutonium equation of state). Nevertheless, the existing program recently accommodated a
 2769 modest amount of scientific work that shows significant promise for IFE. However, magnetically
 2770 driven ICF ultimately needs to achieve robust fusion burn conditions, just as laser or heavy-ion
 2771 ICF do. It has unique features that appear to the panel to provide an alternative risk-mitigating
 2772 path to fusion energy. The Sandia Z100 program has been developed to address some of the key
 2773 target physics issues in pulsed-power ICF. The pulsed-power technology program within the
 2774 NNSA Science Campaigns is developing some of the next-generation technologies that would
 2775 advance the pulsed-power driver issues of a fusion energy technology program. The following
 2776 summarizes the overall program status:

- 2777 • Single-shot, magnetically driven fusion target designs, funded by the NNSA, are
 2778 being investigated on the Z accelerator.
- 2779 • The MagLIF concept has been developed to exploit the favorable ignition
 2780 requirements that, in theory, apply to target designs with magnetized and preheated
 2781 fuel. The MagLIF design is to be investigated in near-term validation experiments
 2782 and simulations.
- 2783 • Benchmark experiments on Z have shown excellent agreement between magneto-
 2784 Rayleigh-Taylor simulations and observations.
- 2785 • Development of an overall system for pulsed-power IFE was supported from 2004 to
 2786 2006 by modest (~\$10 million) internal research funding. Sandia has indicated that
 2787 internally funded research (\$700,000) is now under way to continue the development
 2788 of the RTLs.

2789 Numerous issues surrounding target physics, driver technology, and fusion power system
 2790 parameters stand between the current state of technology and magnetic IFE. These issues include
 2791 the following:

- 2792 • Liner dynamics
 - 2793 —Obtain requisite velocities with suitable shell integrity.
 - 2794 —Demonstrate sufficient control over the fuel adiabat during the implosion (e.g.,
 2795 pulse shaping).
 - 2796 —Demonstrate tolerable levels of mixing at stagnation.
 - 2797 —Demonstrate required level of axial asymmetry.
 - 2798 —Demonstrate required level of azimuthal asymmetry.
- 2799 • Fuel assembly
 - 2800 —Demonstrate the required stagnation pressure.

- 2801 —Demonstrate required confinement time.
- 2802 —Compress sufficient current to a small radius to create extreme conditions.
- 2803 —Compress magnetic flux in the stagnating plasma.
- 2804 • Driver scaling
- 2805 —Determine the driver parameters required for ignition and/or high yield.
- 2806 —Demonstrate scientific breakeven and support target approach with validated
- 2807 simulations.
- 2808 —Develop robust, high-yield targets designs in state-of-the-art 2-D and 3-D
- 2809 simulations.
- 2810 —Demonstrate a repetitive coupling with an RTL system.
- 2811 —Design a system for reliably creating, handling, and utilizing repetitive, high
- 2812 fusion yield with high availability.
- 2813 Some additional specific technical issues still need to be explored:
- 2814 • The MagLIF target design benefits from short implosion times; that is, the final
- 2815 density of the imploded fuel varies as (100 ns)/implosion time. However, the cost
- 2816 and the complexity of the pulsed power driver have the opposite scaling. It was
- 2817 also stated that some target designs might be able to operate at longer implosion
- 2818 times. This would obviously be a huge lever arm on the total system that requires
- 2819 further investigation.
- 2820 • The MagLIF performance scaling simulations have been primarily performed in
- 2821 1-D, with limited exploration of 2-D Rayleigh-Taylor instability issues. However,
- 2822 the physics of thermal conduction and transport in magnetized plasmas is fully 3-
- 2823 D in nature and requires exploration in greater detail. 1-D simulations provide
- 2824 ideal energy scaling; 2-D begins to bring in Rayleigh-Taylor instabilities.
- 2825 Magnetized performance, however, will require 3-D studies.
- 2826 • As stated by Sandia, “batch burn” (volume ignition) will result in a low yield, and
- 2827 a “levitated fuel” layer should give better performance. This will require
- 2828 additional calculations, target fabrication techniques, and experimental
- 2829 implementation. While providing improved performance, it also makes the
- 2830 fabrication and fielding logistics in a fusion power plant more complicated.
- 2831 • Traditional magnetized target fusion concepts have not been shown to scale to
- 2832 high yield and gain. Sandia states that it has recently calculated high-yield
- 2833 performance with MagLIF targets. However, the additional cost of the magnets
- 2834 and optics that would be destroyed on each shot and the complexity of
- 2835 transporting the heater laser through the thick-liquid-wall chamber environment
- 2836 must both be accounted for in the system economics and design.
- 2837 • References from the 2005 Sandia IFE program discuss potential issues of
- 2838 operating RTLs if the final radius and gap become too small. At that time the
- 2839 baseline power flow was relatively large wire-array Z-pinches. It will be
- 2840 important to study the compatibility of the RTL concept with the smaller diameter
- 2841 of direct magnetic-drive targets.
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2847 **Potential for Use in an IFE System**

2848
2849 Concepts for IFE systems using Z-pinch targets were presented to the panel,⁵³ but
2850 sufficient uncertainties remain that it would be premature to attempt an evaluation at this time.
2851 As presently envisioned, each 3-GJ fusion energy pulse would require the insertion, connection,
2852 and energizing of an RTL and fusion target assembly at a 0.1 Hz repetition rate. The assembly is
2853 comprised of an evacuated RTL system that contains the cryogenically cooled Z-pinch target at
2854 its center. The details of this concept are complex and will require extensive research and
2855 development if Z-pinches are pursued as an IFE technology. It is too early in both the target
2856 physics and fusion technology research programs to evaluate the target fabrication and economic
2857 issues quantitatively, but the material and fabrication costs of the expended portions of the
2858 system will certainly be a factor in Z-pinch power plant economics. Because of the limited ICF
2859 target physics database, incomplete validation of the design tools and methodologies, and related
2860 lack of an integrated, high-yield target design, a consistent set of requirements and solutions for
2861 the pulsed power driver, RTL, and ICF target cannot be articulated at this time. Therefore, the
2862 overall credibility of the energy delivery system and the ICF target performance cannot be
2863 quantitatively evaluated.

2864
2865 **CONCLUSION 4-13: Sandia National Laboratory is leading a research effort on a Z-pinch**
2866 **scheme that has the potential to produce high gain with good energy efficiency, but**
2867 **concepts for an energy delivery system based on this driver are too immature to be**
2868 **evaluated at this time.**

2869
2870 The Z-pinch scheme is completely different from the NIF and HIF approaches and
2871 therefore serves as risk mitigation for the ICF and IFE programs. It is not yet clear that the work
2872 at SNL will ultimately result in the high gain predicted by computer simulations, but initial
2873 results are promising and it is the panel's opinion that significant progress in the physics may be
2874 made in a year's time. The pulsed power approach is unique in that its goal is to deliver a large
2875 amount of energy (~10 MJ) to targets with good efficiency (≥ 10 percent) and to generate large
2876 fusion yields at low repetition rates.

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2878 **CONCLUSION 4-14: The target manufacturing and delivery processes that are proposed**
2879 **for direct-drive heavy-ion and pulsed-power fusion energy are less developed conceptually**
2880 **and technically than the targets for laser-based fusion energy.** This is primarily because the
2881 priority has been to emphasize the implosion physics and driver issues (pulsed-power and linear
2882 accelerators). The pulsed-power target appears to be straightforward to manufacture, difficult to
2883 field, and challenging to reprocess after the thermonuclear event. In contrast, the heavy-ion
2884 targets possess many synergies with the laser-based target, but because a final target design is far
2885 from being defined, potential manufacturing complexities cannot be accurately assessed. The
2886 target delivery method for pulsed-power fusion is more conceptual than for laser- or heavy-ion
2887 based fusion and presents very different problems—for example, a very much larger mass
2888 (~1000 times larger), a slower replacement frequency (~100 times slower), and potentially a
2889 greater radioactive waste disposal problem.

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⁵³ M. Cuneo et al., Sandia National Laboratories, "The Potential for a Z-pinch Fusion System for IFE," presentation to the panel on May 10, 2011.

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OUTPUT SPECTRUM FROM VARIOUS IFE TARGETS

The fusion reaction of each type of IFE target produces a spectrum of threats (X-rays, ions, neutrons, and debris) to the first wall of the reaction chamber. The HAPL program studied the spectrum of threats to the first wall posed by direct-drive targets and developed candidate mitigation strategies and materials. It should be noted that while 14 MeV neutrons and 3.5 MeV α -particles are the universal products of the DT fusion reaction, the different target material and configurations for direct drive and indirect drive produce different threat spectra at the reactor chamber first wall. An IFE engineering test facility could be an intermediate step, before full-scale electrical power production, wherein fusion material issues could be studied.

Indirect Drive

The high-Z hohlraum materials used in ID absorb most of the α -particles and radiate more energy as X-rays. The actual threat spectrum is dependent on the details of the hohlraum design. For an ID, heavy-ion target, calculations show that 69 percent of the energy is in neutrons, 25 percent is in X-rays (500 eV peak), and 6 percent is in ions.⁵⁴ For the LIFE target, the X-ray fraction is about 12 percent, the ion fraction about 10 percent, and the remainder in neutrons.⁵⁵ X-rays are the dominant threat to the first wall for ID targets. The Osiris heavy-ion target chamber uses walls wetted by liquid lithium to mitigate the X-ray threat, while LIFE uses Xe gas to protect a dry solid wall.

Direct Drive

DD targets for both KrF and DPPSL systems produce the same threat spectrum, where approximately 1.3 percent of the energy is released in X-rays (4 keV peak) that produce surface deposition in less than the first 1 μm ; 24 percent is in ions that have subsurface deposition in less than 5 μm , and the remainder is in neutrons that have volumetric deposition. Ions produce the greatest first wall heating for direct drive, and the implantation of α -particles presents a helium retention challenge. The HAPL program studied both of these challenges, combining modeling with experiments using lasers, ions, and plasma arc lamps to test thermomechanical cyclic stresses. The helium retention issue was similarly modeled, and experiments were performed on both the Van de Graff and the Inertial Electrostatic Confinement fusion devices at the University of Wisconsin. A nanoengineered tungsten wall material showed an encouraging ability to mitigate helium retention. Experiments showed that cyclic heating in the IFE chamber mitigates helium retention.

Z-Pinch

The spectrum output issues associated with the RTL/Z-pinch system are unique to this approach. The mass of material in this assembly is much greater than in any other concept,

⁵⁴ L.J. Perkins, LLNL, "Targets for Heavy Ion Fusion Energy," a presentation to the panel on February 16, 2011.

⁵⁵ M. Dunne, LLNL, "LIFE Target System Performance," presentation to the panel on July 7, 2011.

2935 leading to greater recycling requirements. Further, the interaction of the fusion output with the
 2936 RTL structure could lead to unique problems with the formation of shrapnel and debris. These
 2937 problems are not presently understood but appear to require a thick liquid-wall chamber.

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TARGET FABRICATION

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2941 The primary concern of this panel with regard to ICF target fabrication relates to the
 2942 technical feasibility of various proposed fabrication methods and the remaining technical risks
 2943 and uncertainties associated with these methods. The question of whether the targets can be
 2944 made cost-efficiently for a power plant is beyond the purview of this panel and is addressed by
 2945 the NRC's IFE committee. Some promising approaches are discussed below.

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Microfluidic Methodologies for Manufacturing Targets

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2950 The polymer shell that contains the DT fuel for DD laser and heavy-ion-beam fusion is
 2951 proposed to be manufactured using a microfluidic droplet formation method.⁵⁶ This is an
 2952 established technology that is used to make ICF capsules for current DD and ID experiments.
 2953 The principle is to flow three immiscible fluids coaxially through two nozzles where the
 2954 Rayleigh-Plateau instability that occurs in the region where they intersect produces individual
 2955 droplets. Each droplet is an emulsion consisting of a thin shell of water surrounding a spherical
 2956 oil droplet; these droplets are collectively immersed in oil. The thin shell of water contains the
 2957 polymer precursors that form the plastic capsule. The final phase of the production process is to
 2958 remove the fluids using supercritical drying.

2959 This process has a very high production rate that is needed for a fusion energy program.
 2960 However, the repeatability and precision of the process must be improved if the process is to be a
 2961 viable option for an energy program. (The repeatability of the current process does not ensure
 2962 that each capsule meets the required specifications, so each capsule is individually measured to
 2963 determine its suitability; this raises the cost of the targets, which is acceptable for ICF
 2964 experiments but not for an IFE program.) In all other aspects, this production process offers a
 2965 potentially viable method for producing targets cost-effectively.

2966 One modification to the current microfluidic method that may improve the reliability is to
 2967 introduce electromechanical control into the process (Cho et al., 2003). This process, referred to
 2968 as "lab-on-a-chip," has demonstrated the feasibility and benefits of using electric fields and
 2969 electronics to control important steps in the target production process (Bei et al., 2010, Wang et
 2970 al., 2011). This concept can potentially reduce the production time and physical size of a target
 2971 production facility and address the precision and reliability concerns with the existing process.
 2972 Further development of the process is needed.

2973 The lab-on-a-chip concept is being evaluated as a method to accomplish the cryogenic
 2974 operation of loading the DT fuel into the capsule.⁵⁷ Preliminary proof-of-concept experiments
 2975 show that it is possible to form individual droplets of liquid deuterium of the correct size and
 2976 wick them into a foam capsule in a short period of time. This would have the benefit of

⁵⁶ A. Nikroo, General Atomics, "Technical Feasibility of Target Manufacturing," presentation to the panel on July 8, 2011; see also Utada et al., 2007.

⁵⁷ R. McCrory, LLE, "Target Fabrication for IFE Reactors: A Lab-on-a-chip Methodology Suited for Mass-Production," submission to the panel on July 6, 2011.

2977 simplifying the target fueling process and shorten the process time, which would reduce the
 2978 tritium inventory that is required by an IFE plant. Additional work is required to further develop
 2979 this concept—specifically, to demonstrate that the process works with tritium and that it is
 2980 practical to apply a condensed gas (argon, neon, or xenon) seal-coat onto the capsule once the
 2981 fuel is loaded.

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OVERARCHING CONCLUSIONS AND RECOMMENDATION

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2985 Based on the discussion in this chapter, the panel reaches the following overarching
 2986 conclusions and makes a recommendation:

2987 **OVERARCHING CONCLUSION 1: NIF has the potential to support the development and**
 2988 **further validation of physics and engineering models relevant to several IFE concepts, from**
 2989 **indirect-drive hohlraum designs to polar direct-drive ICF and shock ignition.**

- 2990 • **In the near to intermediate term, NIF is the only platform that can provide**
 2991 **information relevant to a wide range of IFE concepts at ignition scale. Insofar as**
 2992 **target physics is concerned, it is a modest step from NIF scale to IFE scale.**

- 2993 • **Targets for all laser-driven IFE concepts (both direct-drive and indirect-drive)**
 2994 **can be tested on NIF. In particular, reliable target performance would need to**
 2995 **be demonstrated before investments could confidently be made in development**
 2996 **of laser-driven IFE target designs.**

2997 NIF will also be helpful in evaluating indirectly driven, heavy-ion targets. It will be less
 2998 helpful in gathering information relevant to current Z-pinch, heavy-ion direct-drive, and heavy-
 2999 ion advanced target concepts.

3000 **OVERARCHING CONCLUSION 2: It would be advantageous to continue research on a**
 3001 **range of IFE concepts, for two reasons:**

- 3002 • **The challenges involved in the current laser indirect-drive approach in the**
 3003 **single-pulse NNSA program at the NIF have not yet been resolved and**
- 3004 • **The alternatives to laser indirect drive have technical promise to produce high**
 3005 **gain.**

3006 In particular, the panel concludes that laser direct drive is a viable concept to be pursued
 3007 on the NIF. SNL's work on Z-pinch can serve to mitigate risk should the NIF not operate as
 3008 expected. This work is at a very early stage but is highly complementary to the NIF approach,
 3009 because none of the work being done at SNL relies on successful ignition at the NIF and key
 3010 aspects of the target physics can be investigated on the existing Z-machine. Finally, emerging
 3011 heavy-ion designs could be fruitful in the long term.

3012 **OVERARCHING RECOMMENDATION: The panel recommends against pursuing a**
 3013 **down-select decision for IFE at this time, either for a specific concept such as LIFE or for a**
 3014 **specific target type/driver combination.**

3015 Further R&D will be needed on indirect drive and other ICF concepts, even following
 3016 successful ignition at the NIF, to determine the best path for IFE in the coming decades.

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Appendix A

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Biographical Sketches of Panel Members

John F. Ahearne (NAE) *Chair*, is the executive director emeritus of Sigma Xi, The Scientific Research Society, an adjunct professor of engineering at Duke University, and an adjunct scholar at Resources for the Future. He has extensive expertise in nuclear and radiation engineering and risk assessment. His professional interests are in reactor safety, energy issues, resource allocation, and public policy management. Dr. Ahearne served in the U.S. Air Force from 1959 to 1970, resigning as a major. He has also served as deputy and principal deputy assistant secretary of defense (1972-1977), in the White House Energy Office (1977), as deputy assistant secretary of energy (1977-1978), and as commissioner and chairman of the U.S. Nuclear Regulatory Commission (chairman, 1979-1981). He is a fellow of the American Physical Society, the Society for Risk Analysis, the American Association for the Advancement of Science, the American Academy of Arts and Sciences, and a member of the National Academy of Engineering, Sigma Xi, and the American Nuclear Society. He has previously chaired or served as a member on committees for over 30 other NRC studies. Dr. Ahearne received a Ph.D. in physics from Princeton University.

Douglas Eardley, *Vice Chair*, is professor of physics at the Kavli Institute for Theoretical Physics at the University of California, Santa Barbara. Dr. Eardley's research interests include general relativity: black holes, gravity waves, and quantum gravity; theoretical astrophysics: X-ray sources, quasars, active galactic nuclei, and cosmology; mathematical physics: nonlinear partial differential equations and geometry; physics and society: national security, nuclear weapons, and arms control. Dr. Eardley has been a member of several National Research Council study committees, including the Working Group on Related Areas of Science of the Astronomy Survey Committee ("Field Committee") in 1979-1980; the Committee on the Atmospheric Effects of Nuclear Explosions in 1983-1984; and the Science Panel of the Astronomy Survey Committee in 1989-1990. He was chair of the External Advisory Board of the Institute for Fundamental Theory of the University of Florida at Gainesville from 1990 to 1994; a member of the Physics Advisory Committee of Lawrence Livermore National Laboratory from 1991 to 1996; the plenary speaker at the Texas Symposium on Relativistic Astrophysics in 1992; a member of the Openness Advisory Panel of the Secretary of Energy Advisory Board for DOE from 1996 to 2002; and co-coordinator of the Institute for Theoretical Physics' Program in Black Hole Astrophysics from 1999 to 2002. Professor Eardley has been a member of the JASON Study Group since 1981; a member of the National Security Panel of the University of California's President's Council on the National Laboratories from 2000 to 2007; chair of the External Review Panel for the Radiation Effects Sciences Program for Sandia National Laboratories since 2000; and a member of the Joint Mission Committee for Los Alamos National Laboratory and Lawrence Livermore National Laboratory since 2007. He received a B.S. in physics from the California Institute of Technology and M.S. and Ph.D. degrees in physics from the University of California, Berkeley.

3214
3215 **Robert C. Dynes (NAS)** is professor emeritus of physics at the University of California, San
3216 Diego. He served as the 18th president of the University of California (UC) from 2003 to 2007,
3217 and as chancellor of UC San Diego from 1996 to 2003. His position as chancellor followed 6
3218 years in the physics department, where he founded an interdisciplinary laboratory in which
3219 chemists, electrical engineers, and private industry researchers investigated the properties of
3220 metals, semiconductors, and superconductors. Prior to joining the UC faculty, he had a 22-year
3221 career at AT&T Bell Laboratories, where he served as department head of semiconductor and
3222 material physics research and director of chemical physics research. Dr. Dynes received the 1990
3223 Fritz London Award in Low Temperature Physics, was elected to the National Academy of
3224 Sciences in 1989, and is a fellow of the American Physical Society, the Canadian Institute for
3225 Advanced Research, and the American Academy of Arts and Sciences. He serves on the
3226 executive committee of the U.S. Council on Competitiveness. A native of London, Ontario,
3227 Canada, and a naturalized U.S. citizen, Dr. Dynes holds a bachelors degree in mathematics and
3228 physics and an honorary doctor of laws degree from the University of Western Ontario, and
3229 masters and doctoral degrees in physics and an honorary doctor of science degree from
3230 McMaster University. He also holds an honorary doctorate from Université de Montréal.

3231
3232 **David Harding** is a senior scientist at the University of Rochester's Laboratory for Laser
3233 Energetics and a professor in the Department of Chemical Engineering. His research interests
3234 include the science and engineering associated with the making of fuel capsules for fusion
3235 experiments performed at the University of Rochester's Laboratory for Laser Energetics. He has
3236 worked at the University of Rochester for 15 years; prior to that he was a senior research
3237 engineer in the Materials and Structures Division at the NASA Lewis Research Center. He has
3238 participated as a panel member on two review committees: the National Ignition Facility Target
3239 Fabrication Review (2008) at Lawrence Livermore National Laboratory and a DOE review of its
3240 Solar Thermal Program (1992). Dr. Harding received a Ph.D. from Cambridge University.

3241
3242 **Thomas Mehlhorn** is superintendent of the Naval Research Laboratory (NRL) Plasma Physics
3243 Division, and a member of the Department of the Navy Senior Executive Service with
3244 responsibility for a broad spectrum of research programs in plasma physics, laboratory discharge
3245 and space plasmas, intense electron and ion beams and photon sources, atomic physics, pulsed
3246 power sources, radiation hydrodynamics, high-power microwaves, laser physics, advanced
3247 spectral diagnostics, and nonlinear systems. He began his career at Sandia National Laboratories
3248 in 1978 and worked on a variety of projects related to the generation, focusing, and interaction of
3249 intense beams of electrons and ions with plasmas. From 1989 to 1998 he was a manager in the
3250 Sandia Light Ion ICF Program and from 1998 to 2006 he managed Sandia's High Energy
3251 Density Physics and ICF Target Design Department in the Pulsed Power Fusion Program. From
3252 2006 to 2009 he was a senior manager with accountability for dynamic materials and shock
3253 physics, high energy density physics theory and modeling, and advanced radiographic source
3254 development and applications. Dr. Mehlhorn joined NRL in 2009. He is a recipient of two
3255 NNSA Defense Programs Award of Excellence (2007 and 2008), a Lockheed Martin NOVA
3256 award (2004), and an Alan Berman Research Publication Award from NRL (1983). Dr.
3257 Mehlhorn is a fellow of the American Association for the Advancement of Science (AAAS) in
3258 Physics (2006). He serves on the Advisory Board for Plasma and Atomic Physics at GSI,
3259 Darmstadt, Germany (2004-present, chair in 2006). He is a member of the Nuclear Engineering

3260 and Radiological Sciences Department Advisory Board at the University of Michigan (1996-
3261 1999, and 2004-present), as well as of the University of Michigan College of Engineering
3262 Alumni Society board of governors (2009-present). In 2010 Dr. Mehlhorn served on the
3263 Department of the Navy Space Experiments Review Board as well as the University of
3264 Missouri's Research and Development Advisory Board. Dr. Mehlhorn received B.S, M.S., and
3265 Ph.D. degrees in nuclear engineering from the University of Michigan.
3266

3267 **Merri Wood-Schultz** is a part-time consultant for SAIC and serves as a laboratory associate at
3268 LANL for Improvised and Foreign Devices. Dr. Wood-Schultz's early career focused on the
3269 physics design of secondaries of thermonuclear weapons. She was responsible for the conceptual
3270 and physics design of numerous nuclear tests and add-on experiments; the areas of focus of these
3271 tests included stockpile systems, weapons physics, and advanced development. Dr. Wood-
3272 Schultz played an active role in the development of nuclear weapons-related laboratory
3273 experiments (AGEX), serving as the lead designer for a series of experiments on the Sandia
3274 National Laboratories' SATURN pulsed-power machine and as a member of the inaugural
3275 LANCE (neutron scattering facility) Users Group. Later phases of Dr. Wood-Schultz's career
3276 included involvement in developing concepts and methods for certification without nuclear
3277 testing, notably the quantification of margins and uncertainty (QMU), and an increase in her
3278 work in nuclear intelligence. The latter led to a 6-month, change-of-station assignment to a DOE
3279 intelligence organization. Dr. Wood-Schultz is currently a member of the Nuclear Forensics
3280 Science Panel for the Department of Homeland Security and engages in continuing technical
3281 collaborations on nuclear weapons design, yield certification using QMU, and nuclear
3282 intelligence. Dr. Wood-Schultz became a fellow of Los Alamos National Laboratory in 2001,
3283 received the Department of Energy Award of Excellence in 1988, 1999, and 2004, the
3284 STRATCOM Medal of Excellence in 1997, and the Los Alamos National Laboratory
3285 Distinguished Performance Award in 1996. Dr. Schultz received B.S., M.S., and Ph.D. degrees
3286 in physics from the Georgia Institute of Technology.
3287

3288 **George Zimmerman** is a part-time consultant on computations and modeling for LLNL and on
3289 nuclear reactor modeling for TerraPower, LLC. He joined LLNL in 1970 as a staff member in
3290 the A Division, where he developed the LASNEX computer program to design laser fusion
3291 targets and analyze experiments. In 1980 he was appointed associate division leader in the X
3292 Division, where he led a group of physicists responsible for developing numerical methods to
3293 accurately perform integrated simulations involving laser absorption, magnetohydrodynamics,
3294 atomic physics, and the transport of photons, neutrons, and charged particles. From 1984 to
3295 1987 he was leader of the Computational Physics Division. He then led the inertial confinement
3296 fusion code development project in the AX Division until his retirement. Mr. Zimmerman
3297 received the Department of Energy's 1983 E.O. Lawrence Award for contributions to national
3298 security and the 1997 Edward Teller Award for developing the LASNEX inertial confinement
3299 fusion code. He also received the Defense Programs Award of Excellence for significant
3300 contributions to the Stockpile Stewardship Program in 2002 and 2005. He retired from
3301 Lawrence Livermore National Laboratory (LLNL) in 2007 and is currently a fellow of the
3302 American Physical Society. Mr. Zimmerman received a B.S. in physics from Harvey Mudd
3303 College and an M.A. in astronomy from the University of California, Berkeley.
3304

Appendix B

Panel Meeting Agendas and Presenters

WASHINGTON, D.C.
FEBRUARY 16-17, 2011

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3315 Call to order and welcome
3316 *John Ahearne, Chair*
3317
3318 Overview of the study task and origins and the National Academies' study process
3319 *Sarah Case, Study Director; John Ahearne, Chair*
3320
3321 IFE committee briefing to the panel on expectations
3322 *Gerald Kulcinski, Inertial Fusion Energy Committee Co-Chair*
3323
3324 Review of charge to the panel, the U.S. Department of Energy's interests in the committee and
3325 panel reports, and nuclear weapons proliferation risks for an inertial fusion energy program
3326 *David Crandall, Office of the Under Secretary for Science, U.S. Department of Energy System*
3327
3328 Indirect drive target physics at the National Ignition Facility (NIF)
3329 *John Lindl, Lawrence Livermore National Laboratory*
3330
3331 Direct drive target physics at the Naval Research Laboratory (NRL)
3332 *Andrew Schmitt, NRL*
3333
3334 Direct drive target physics at NIF
3335 *David Meyerhofer, Laboratory for Laser Energetics*
3336
3337 Heavy ion target physics
3338 *John Perkins, Lawrence Livermore National Laboratory*
3339
3340 Z pinch target physics
3341 *Mark Herrmann, Sandia National Laboratory*
3342
3343 Non-proliferation considerations associated with inertial fusion energy
3344 *Raymond Jeanloz, University of California, Berkeley*
3345
3346
3347 PLEASANTON, CALIFORNIA
3348 APRIL 6-7, 2011
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3350

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- 3351 Welcome and call to order
 3352 *John Ahearne, Chair*
 3353
 3354 System considerations for IFE
 3355 *Tom Anklam, Lawrence Livermore National Laboratories (LLNL)*
 3356
 3357 Overview of laser inertial fusion energy system and key considerations for IFE targets
 3358 *Michael Dunne, LLNL*
 3359

ALBUQUERQUE, NEW MEXICO
 May 10-11, 2011

- 3360
 3361
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 3363
 3364
 3365 Welcome and call to order
 3366 *John Ahearne, Chair*
 3367
 3368 Inertial confinement fusion (ICF) targets at Los Alamos National Laboratory (LANL)
 3369 *Juan Fernandez, LANL*
 3370
 3371 Design and simulation of magnetized liner inertial fusion targets
 3372 *Steve Slutz, Sandia National Laboratories (SNL)*
 3373

ROCHESTER, NY
 July 6-8, 2012

- 3374
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 3377
 3378
 3379 Welcome and call to order
 3380 *John Ahearne, Chair*
 3381
 3382 Welcome and overview of Laboratory for Laser Energetics (LLE) ICF program
 3383 *Robert McCrory, LLE*
 3384
 3385 Direct-drive progress on OMEGA
 3386 *Craig Sangster, LLE*
 3387
 3388 Polar drive target design
 3389 *Radha Bahukutumbi, LLE*
 3390
 3391 Facilitating NIF for polar drive
 3392 *David Meyerhofer, LLE*
 3393
 3394 Fast and shock ignition research
 3395 *David Meyerhofer, LLE*
 3396

- 3397 LPI issues for direct drive
3398 *Dustin Froula and Jason Myatt, LLE*
3399
3400 Heavy ion target design
3401 *B. Grant Logan, Lawrence Berkeley National Laboratory*
3402
3403 Discussion of LIFE targets and program
3404 *Michael Dunne, Lawrence Livermore National Laboratories*
3405
3406 Technical feasibility of target manufacturing
3407 *Abbas Nikroo, General Atomics*
3408
3409
3410
3411 WASHINGTON, D.C.
3412 September 20-22, 2012
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3414
3415 Welcome and call to order
3416 *John Ahearne, Chair*
3417
3418 Development of the technologies for laser fusion direct drive
3419 *John Sethian, NRL*
3420
3421 Overview of current NRL program for ICF/IFE
3422 *Steve Obenschain and Andrew Schmitt, NRL, and Frank Hegeler, Commonwealth Technology at*
3423 *NRL*
3424
3425 Overview of LPI Physics and LANL understanding
3426 *David Montgomery, LANL*
3427
3428 Understanding of LPI and its impact on indirect drive
3429 *Mordechai Rosen, LLNL*
3430
3431 Assessment of understanding of LPI for direct drive (solid-state)
3432 *Dustin Froula, LLE*
3433
3434 Assessment of understanding of LPI for direct drive (KrF)
3435 *Andy Schmitt, NRL*
3436
3437 State of the art for LPI simulation
3438 *Denise Hinckel, LLNL*
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Appendix C

Acronyms

3443		
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3446		
3447	AEC	Atomic Energy Commission
3448	ASC	Advanced Simulation and Computing
3449	CBET	cross-beam energy transfer
3450	CEA	Atomic Energy Commission (France)
3451	CH	carbon-hydrogen (plastic used as an ablator)
3452	CTBT	Comprehensive Nuclear Test Ban Treaty
3453	CVD	chemical vapor deposition
3454	DCA	Detailed Configuration Analysis
3455	DD	direct drive
3456	DOE	U.S. Department of Energy
3457	DPSSL	diode-pumped, solid-state laser
3458	DT	deuterium-tritium
3459	EP	extended performance
3460	EU	European Union
3461	FCT	flux-corrected transport
3462	FI	fast ignition
3463	FIRE	Fast Ignition Realization Experiment
3464	HAPL	High Average Power Laser (Program)
3465	HDC	high-density carbon
3466	HIF	heavy-ion fusion
3467	HiPER	High Power Laser Energy Research (facility)
3468	IAEA	International Atomic Energy Agency
3469	ICF	inertial confinement fusion
3470	ID	indirect drive
3471	IFE	inertial fusion energy
3472	ILE	Institute of Laser Engineering (Japan)
3473	IR	infrared
3474	ISI	induced spatial incoherence
3475	LANL	Los Alamos National Laboratory
3476	LBNL	Lawrence Berkeley National Laboratory
3477	LEH	laser entrance hole
3478	LIFE	Laser Inertial Fusion Energy
3479	LLE	Laboratory for Laser Energetics
3480	LLNL	Lawrence Livermore National Laboratory
3481	LMJ	Laser Megajoule Facility (France)
3482	LPI	laser-plasma interactions
3483	LTE	local thermal equilibrium
3484	MagLIF	magnetized liner inertial fusion
3485	MAGO	Explosively Driven Magnetic Implosion (Russia)
3486	MCF	magnetic confinement fusion
3487	MRT	magneto-Rayleigh-Taylor
3488	MTF	magnetized target fusion

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3489	NAS	National Academy of Sciences
3490	NIC	National Ignition Campaign
3491	NIF	National Ignition Facility
3492	NNSA	National Nuclear Security Administration
3493	NPT	Nonproliferation Treaty
3494	NRC	National Research Council
3495	NRL	Naval Research Laboratory
3496	FES	Fusion Energy Sciences
3497	ORNL	Oak Ridge National Laboratory
3498	PD	Polar Drive
3499	PEP	project execution plan
3500	RT	Rayleigh-Taylor
3501	RTL	recyclable transmission line
3502	SI	shock ignition
3503	SNL	Sandia National Laboratory
3504	SNM	special nuclear materials
3505	SRS	stimulated Raman scattering
3506	SSD	smoothing by spectral dispersion
3507	THD	tritium-hydrogen-deuterium
3508	TPD	two-plasmon decay
3509	VNIIEF	All-Russian Research Institute of Experimental Physics
3510	YOC	yield over clean
3511		
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